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HAL Id: hal-01427475
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Submitted on 30 Jul 2017

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Energy balance of droplets impinging onto a wall heated above the Leidenfrost temperature

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Keywords:
Leidenfrost effect
Spray cooling
Film boiling
Laser-Induced Fluorescence
Infrared thermometry

Abstract

This work is an experimental study aiming at characterizing the heat transfers induced by the impingement of water droplets (diameter 80–180 μm) on a thin nickel plate heated by electromagnetic induction. The temperature of the rear face of the nickel sample is measured by means of an infrared camera and the heat removed from the wall due to the presence of the droplets is estimated using a semi-analytical inverse heat conduction model. In parallel, the temperature of the droplets is measured using the two-color Laser-Induced Fluorescence thermometry (2cLIF) which has been extended to imagery for the purpose of these experiments. The measurements of the variation in the droplet temperature occurring during an impact allow determining the sensible heat removed by the liquid. Measurements are performed at wall conditions well above the Leidenfrost temperature. Different values of the Weber numbers corresponding to the bouncing and splashing regimes are tested. Comparisons between the heat flux removed from the wall and the sensible heat gained by the liquid allows estimating the heat flux related to liquid evaporation. Results reveal that the respective level of the droplet sensible heat and the heat lost due to liquid vaporization can vary significantly with the droplet sizes and the Weber number.

1. Introduction

Liquid cooling is widely used in applications, which requires a high heat dissipation rate. Cooling techniques such as pool boiling or jet impingement can provide high heat dissipation rates, but they generally fail to insure a uniform cooling. For example, studies on liquid jet impingement show that heat transfers are very high at the stagnation zone of the jet, whereas the spreading region is characterized by moderate heat transfer coefficients (Webb et al., 1995). Comparatively, spray cooling technology is of increasing interest since it is characterized by high heat transfer rates, uniformity of heat removal and small fluid inventory. In electronic systems and power electronics, spray cooling is required to maintain lower operating temperature of the component (Kim, 2007). In the steel industry, environmental and economic constraints have imposed greater demands for a reduction of the water and energy consumptions.

Spray quenching is very efficient compared to other cooling techniques. The main reason is that vapor can easily escape even if the temperature of the wall is well above the Leidenfrost temperature. However an optimization of the industrial processes and an increase of the predictive capabilities in that field need a complete understanding of the complex fluid flow and heat transfer characteristics when sprays interact with hot surfaces. The investigation of the impingement of droplets on solid surfaces has received a considerable attention throughout the decades. Up to now, mechanisms are poorly understood. This is mainly due to the dependence of these phenomena on many parameters, which cannot be easily varied independently. Many studies have been focused on impacts at high or low droplet velocity, with deep or shallow liquid film, on hot or cold solid surfaces, sometimes on micro- or nanostructured surfaces (Bhushan and Chae Jung, 2008; Lembach et al., 2010). Depending on these conditions, different behaviors can occur: the drops can spread over the solid surface, can splash by creating a crown or can rebound (Yao and Cai, 1988). Extensive experimental investigations were carried out to determine the parameters influencing the behavior of a single drop impact in order to characterize their respective influence. Some of these parameters describe the geometry and the dynamic of the drops (Rioboo et al., 2001; Yao and Cai, 1988), some refer to the physical properties of the liquid (Rioboo et al., 2001; Sikalo et al., 2002) or the solid surface (Cossali et al., 2006). Correlations based on dimensionless numbers characterizing the relative magnitude of the forces acting on the impinging droplet and estimated with the physical properties of the liquid before impact, i.e. Reynolds, Weber, Ohnesorge and Mundo numbers have been found (García Rosa et al., 2006).

In this study, the focus is placed on non-wetting conditions; the wall temperature is above the Leidenfrost temperature, which

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corresponds to the film boiling regime. A thin vapor layer forms quasi-instantaneously between the droplet and the wall and prevents the droplet to stick the wall. The splashing and the rebound of the droplets are thus the only behavior that can occur. The rebound regime is observed for low Weber number while an increase of the droplets are thus the only behavior that can occur. The rebound and splashing regimes, velocity of the outcoming droplets has been widely investigated (Mundo et al., 1995; Schmehl et al., 1999; Wachters and Westerling, 1966). In the same way, post-impact droplets size distribution has been widely investigated in the literature, for temperatures greater than the Leidenfrost limit (Dewitte et al., 2005) or above (Schmehl et al., 1999). To the best of our knowledge, there is no data in the literature related to the post-impact droplet temperature except the recent works of Castanet et al. (2009) and Dunand et al. (2012).

Heat transfers at the wall were generally characterized using thermocouples embedded in the sample thickness. It has been possible to monitor the history of the surface temperature at the location of the droplet impingement (Baumeister and Simon, 1973). The estimation of the heat flux extracted from the wall was also made possible when the experimental set-up was designed to ensure a one-dimensional heat flux along the instrumented section containing a set of thermocouples. Nevertheless, if measurements related to the wall provide valuable data to quantify the efficiency of the cooling, they have only a limited interest when focusing on the heat transfer occurring within the liquid phase. In particular, questions remain concerning the respective level of the droplet sensible heat variation and the heat removed due to liquid vaporization. When splashing occurs, very few correlations for the mass loss during the impact can be found in the literature, despite of its importance for practical applications. In the bouncing regime, the ratio between outgoing and impinging liquid mass was evaluated (Le Clercq (2000), using Phase Doppler measurements (PDA) and digital image processing. However, the outcoming droplets may be strongly deformed after their impingement and measuring small variations in their volume with a direct optical method such...
as PDA is a real challenge. An experimental correlation is proposed for the relative loss in mass of the droplets. It is based on the Mun
do number, the liquid boiling temperature, Leidenfrost tempera-
ture and wall temperature.

Presently, a quite different approach has been undertaken. The basic idea is to use combined measurement techniques for an indi-
rect estimate of the mass of liquid evaporated during the droplet/ wall interaction. Heat transfers are characterized within the liquid droplets and at the wall. The temperature variation of the droplet is measured using the two-color planar Laser-Induced Fluorescence (2cPLIF) thermometry. In addition, an infrared camera provides the temperature field at the rear face of a thin nickel target, heated by electromagnetic induction. A semi-analytical inverse heat con-
duction model allows estimating the heat flux on the front face of the plate where the droplets impinge. Finally, the heat flux re-
moved from the wall by the droplets is compared to the sensible heat stored in the outgoing droplets. Energy conservation is finally
invoked to estimate the heat flux associated to evaporation. The respective contributions of the liquid sensible heat, the heat of
evaporation, and the heat removed from the wall are analyzed in
terms of incident droplet size and normal Weber number.

2. Droplet generation and experimental set-up

In order to study droplet/wall interactions, an experimental set-
up was specifically designed. A sketch of the experimental set-up is
shown in Fig. 1. A linear monodisperse droplet stream is generated by
the disintegration of a cylindrical liquid jet. The breakup is driven by a Rayleigh-type instability that can be triggered by mecha-
nical vibrations using a piezoceramic. For some specific frequencies of the vibrations, the liquid jet split into equally spaced and mono-
disperse droplets. The size of the injector orifice and the inlet pressure
are fixed by the disintegration of a cylindrical liquid jet. The breakup is dri-

The droplet generator can be rotated to any prescribed angle
finj of incidence. The temperature of the injector body is regulated
by a Rayleigh-type instability that can be triggered by mechan-
ical vibrations using a piezoceramic. For some specific frequencies
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disperse droplets. The size of the injector orifice and the inlet pressure
are fixed by the disintegration of a cylindrical liquid jet. The breakup is dri-

Fig. 1. Experimental set-up.

3. Measurement techniques

3.1. Measurements of the droplets size and velocity by shadowgraphy

A high-speed (HS) camera is used to visualize droplets imping-
ing onto the heated wall. The HS camera is a Phantom v710 equipped with a 12-bits CMOS sensor that can provide up to
7500 fps at full resolution (1280 × 800 pixels). It is used with a re-
duced resolution to perform the image acquisition at a much higher frame rate, typically in the order of 100,000 fps. This acquisition rate is sufficient to resolve in time the droplet/wall interactions in the experimental conditions encountered in this study. The dro-

3.2. Two-color planar Laser-Induced Fluorescence thermometry

The two-color planar Laser-Induced Fluorescence (2cPLIF) was used to measure the variation in the droplet temperature during their interaction with the wall. This technique already demonstrated its ability to characterize the temperature of droplets in various situations including droplet evaporation in either inert or reactive flows (Castanet et al., 2003; Depréduand et al., 2010). It was also used to determine the droplet change in temperature during

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wise measurements, which imply a cumbersome point-by-point scanning to reconstruct the temperature distribution in the liquid phase of the flow. More recently, the technique was extended to planar laser induced fluorescence (PLIF) in order to obtain the tem-

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wise measurements, which imply a cumbersome point-by-point scanning to reconstruct the temperature distribution in the liquid phase of the flow. More recently, the technique was extended to planar laser induced fluorescence (PLIF) in order to obtain the tem-

The 2cLIF thermometry is based on the measurement of the fluorescence intensity of a single dye tracer. In liquids, the fluorescence quantum yield is strongly influenced by the quenching, which depends on the temperature. When a laser beam induced the fluorescence of a dye dissolved into a liquid, the intensity of the fluorescence signal detected on a given spectral band \( i \) can be expressed as (Castanet et al., 2003):

\[
I_{f,i} \approx K_{opt,i}K_{spec,i}I_0cV \exp \left( \frac{\beta_i}{T} \right), \quad i = 1 \text{ or } 2
\]

where \( K_{opt,i} \) is a parameter depending on the optical properties of the detection system (e.g. the solid angle of the detection, the spectral sensitivity of the detectors, the spectral band of detection), \( K_{spec,i} \) is a parameter depending on the spectroscopic properties of the tracer in its solvent on the designated spectral band. The parameter \( c \) is the concentration in dye molecules and the product \( cV \) corresponds to the number of molecules that are illuminated by the laser beam in the field of view of the detector. \( I_0 \) is the intensity of the laser beam before crossing the absorbing medium. In Eq. (1), it is implicitly assumed that the absorption of the laser beam and the fluorescence can be neglected along the ray path in the liquid medium. The parameter \( \beta_i \) corresponds to the temperature sensitivity of the fluorescence signal detected on the spectral band \( i \). They are specific to a given combination of dye, solvent, excitation wavelength, and spectral band of detection. In contrast, \( K_{opt,i} \) depends on the exact configuration of the experimental system and can change from one measurement configuration to another. For this reason, it must be determined by a reference measurement. In this study, rhodamine 640 (C\(_{32}\)H\(_{31}\)N\(_2\)O\(_3\)Cl\(_6\), also called rhodamine 101) was selected as a fluorescent tracer. The ratio of the fluorescence intensity measured on two bands, for which the temperature sensitivity is highly different (Lavieille et al., 2001), allows eliminating the effects of parameters that are unknown or difficult to control such as the laser intensity, the tracer concentration, the measurement volume \( V \) (the triple intersection between the laser excitation volume, the droplet and field of view of the collection optics), which varies continuously during the droplet transit in the probe volume. When the technique is applied in imagery, only the coefficients \( \beta_i \) do not depend on the pixel position in the image. All other variables can change from one pixel to the other, especially the parameter \( K_{spec,i} \). Even under isothermal conditions, the fluorescence ratio is not necessarily uniform, due mainly to the non-uniformity of the CCD detection matrix. To eliminate the influence of the detection system, a reference image at a known temperature \( T_0 \) (with the same optical configuration as for the measurement) is recorded. According to Eq. (1), denoting \( R_0 \) the fluorescence ratio obtained in the reference measurement, the temperature can be derived from:

\[
\ln \left( \frac{R_i}{R_0} \right) = (\beta_1 - \beta_2) \left( \frac{1}{T} - \frac{1}{T_0} \right).
\]

where \( R_i = I_{f,i}/I_0 \) and \( R_0 = I_{f,0}/I_0 \). Once the difference \((\beta_1 - \beta_2)\) is known, Eq. (2) can be used to determine the liquid temperature. Parameters \( \beta_1 \) and \( \beta_2 \) are obtained by a calibration in a temperature-controlled cell using the measurement system with the appropriate optical filters. The bands of detection correspond to the ranges [555–565 nm] and [635–685 nm]. They are selected with regards to their intensity level as well as their sensitivity to the temperature. The variation of the fluorescence ratio as a function of the temperature, measured in a temperature regulated cell is depicted in Fig. 2. The variation of the fluorescence ratio \( R_i \) is about 1.4%/K which is enough in practice to measure the droplet temperature with an accuracy of about ±2°C. A more precise estimate of the measurement uncertainties is given in Appendix A.

The measurement system is illustrated in Fig. 3. The excitation of Rh640 is achieved by means of a CW:YAG laser (Laser Quantum Finees, 6W @532 nm). An arrangement of spherical and cylindrical lenses provides a laser sheet with a thickness of 220 μm and a height of 16 mm in the measurement zone. This latter is observed by a Questar QM-1 long distance microscope, which is positioned at right angle at a working distance of about 84 cm. The microscope field of view is then about 3.5×3.5 mm\(^2\). A holographic filter (Notch Plus, Kayser Optical) is used to block the Mie scattering of the laser light at 532 nm. A neutral beamsplitter (R/T 45/55%) allows splitting the fluorescence signal for its acquisitions by the cameras. Interference filters are mounted in front of the cameras and allows selecting the aforementioned spectral bands. For the detection of the fluorescence images, two electron-multiplying CCD cameras (Hamamatsu 14 bits EM-CCD camera C9100-02) with a spatial resolution of 1000 × 1000 pixels are used. To improve the detection statistics, a 4 × 4 binning of the pixels is applied, even if the spatial resolution is reduced. Two images (one for each camera) are acquired simultaneously by using a common external trigger source. During the experiments, the cameras exposure time is set to a few tens of milliseconds to ensure a sufficient signal-to-noise ratio. Meanwhile, the droplets injected at about 10 kHz cover several millimeters and their contribution to the fluorescence signal is cumulated. The images correspond thus to a time averaged fluorescence field. Measurements were also performed with moderate gains to improve the signal/noise ratio.
One of the main difficulties related to this experimental setup is the pixel-by-pixel correspondence of the camera images. The approach described in Dunand et al. (2012) is applied here. The least mean squares are used to find a combination of rotation and translation that minimizes the distance between the images of the fluorescence field taken by the two cameras. The solution obtained for one pair of images is generally optimal, i.e. it can be applied to another pair of images with a maximum error less than one pixel.

Finally, the measurements are performed with a concentration in Rh640 equal to $5 \times 10^{-5}$ mol/L. This concentration is relatively high and the re-absorption of the fluorescence within the droplet cannot be ignored. However, this high concentration is required to limit the effect of the droplet size and shape on the fluorescence spectrum which has been described by Labergue et al. (2010). The re-absorption of the fluorescence is likely to modify differently the fluorescence ratios of incoming and outgoing droplets only in the case of a splashing. In the case of a rebound, the droplet does not change significantly in diameter. Re-absorption of the fluorescence can be accounted for in the case of a splashing. The extinction coefficients of each spectral band being known, the fluorescence ratio can be corrected by:

$$R_{\text{corr}} = R_f \cdot \mu = \frac{\int_{x_{\text{min}}}^{x_{\text{max}}} l_{\text{ex}}(x) \exp(-e_1 c x) dx}{\int_{x_{\text{min}}}^{x_{\text{max}}} l_{\text{em}}(x) \exp(-e_2 c x) dx},$$

In this expression, $\mu$ is a correction factor taken into account the effect of absorption, $x$ is a given position in the droplet, $e_1$ and $e_2$ the coefficients of absorption related to the spectral bands of detection and $V_d$ is the droplet volume. The evolution of $\mu$ is plotted in Fig. 4 as a function of the droplet diameter for a concentration in Rh640 equal to $5 \times 10^{-5}$ mol/L. In the case of a splashing, only the ratio corresponding to droplet before impact is corrected. The sizes of the secondary droplets are generally too small to be significantly influenced by the re-absorption of the fluorescence. For a 200 µm, the applied correction is on the order of 2.8 °C and it decreases to about 1.5 °C when the diameter is 100 µm.

Fig. 5 shows an example of measurement in the case of a splashing and Fig. 6 in the case of a rebound. A significant heating of the droplets resulting from their impingement is observed. The fluorescence field is not uniform in an image; this is mainly related to the time averaged liquid concentration, which varies strongly in space. The liquid concentration is the more important near the impact region where the droplet are strongly squeezed. From these images of the temperature field, the average temperature of primary and secondary droplets can be calculated. The average is weighted by the fluorescence intensity since this latter is roughly proportional to the liquid mass flow rate crossing the region of interest (ROI) during the integration time of the cameras:

$$T_m = \frac{\int_{\text{ROI}} T(x,y) \, dA \cdot l_{\text{ex}}(x,y) \, dy}{\int_{\text{ROI}} l_{\text{em}}(x,y) \, dy}$$

ROI are defined for the incident droplets and the secondary droplets. The difference in temperature $\Delta T$ between these regions is finally computed and allows evaluating the gain of sensible heat of the liquid.

3.3. Infrared thermography and inverse conduction model for the wall heat flux estimation

In all the experiments, the nickel sample is first heated up to 700–750 °C. Then heating is stopped and cooling by the water droplets stream occurs. As the slab is impacted by the droplets...
on one of its face (referred as “front face”), the temperature field is measured on the other face (referred as “rear face”). The IR camera is built around a focal plane array of photonic detectors working in the [3–5 μm] spectral range. It is equipped with a narrow [3.97–4.01 μm] filter. Acquisitions are performed at sampling frequencies of about 60 Hz and a maximum resolution of 320 × 240 pixels. A specific inverse heat conduction algorithm was developed to recover the heat flux removed from the front face corresponding to the droplets. This requires the analytical resolution of the heat transfer equation using integral transforms.

3.3.1. Solution of the direct heat transfer problem

The internal transient conduction within a disc whose radius is \( R = 12.5 \) mm and thickness \( \varepsilon = 500 \) μm is considered. As the resident time (duration of the droplet interaction with the wall) is very low compared to the time associated with the sampling frequency of the IR camera, an averaged value of the heat flux over the resolution time of the camera (or a multiple of it) will be inferred from the measurements. In order to decrease the number of unknowns associated with the flux distribution, the 3D modelling is reduced to a 2D modelling through angular averaging in a cylindrical coordinate system. Then, assuming constant thermophysical properties, the following set of equations is obtained:

\[
\begin{align*}
\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial z^2} &= \frac{1}{a} \frac{\partial T}{\partial t}, \\
T(r, z, t = 0) &= T_{\text{in}}(r), \\
-\lambda \frac{\partial T}{\partial r}_{r=0} &= 0 \quad \text{and} \quad -\lambda \frac{\partial T}{\partial r}_{r=R} = 0, \\
\lambda \frac{\partial T}{\partial z}_{z=0} &= q_f \quad \text{and} \quad -\lambda \frac{\partial T}{\partial z}_{z=\varepsilon} = q_R,
\end{align*}
\]

with : \( q_f = h_{EQ}^f (T_F - T_m) + q_d (r, t) \) and \( q_R = h_{EQ}^R (T_R - T_m) + K_{\text{cond}} \delta (r - R) (T_R - T_{\text{stum}}) \).

\( h_{EQ}^f \) corresponds to the heat loss coefficient which is the sum of both convective and radiative losses over the rear \((\beta = R)\) and front \((\beta = F)\) faces of the disc. \( q_d \) is the heat flux removed from the front face by the droplet stream. It is averaged angularly over a circle at radius \( r \). \( T_m \) is the air temperature and \( T_F \) and \( T_R \) are the front and rear face temperatures, \( T_{\text{in}} \) is the initial temperature field and \( T_{\text{stum}} \) the temperature of the support (three insulating beads placed at 120° and a radius \( R \), in between the Nickel disc and a hollow cylindrical support in stumatite (a ceramic). \( K_{\text{cond}} \) is a contact conductance between sample and support. Convective contributions to \( h_{EQ}^f \), have been calculated using natural convection correlations over horizontal surfaces. Linearized radiative contributions to \( h_{EQ}^f \), the radiative transfer coefficients, differ since the front face radiative environment is the ambient while the rear face is coupled with both the ambient and the stumatite support. All these coefficients have been estimated thanks to a relaxation experiment in the absence of any droplet stream. The resolution of the direct problem is detailed in Appendix B. It uses different integral transforms in time and space.

3.3.2. Inverse heat transfer problem

Solving the Inverse Heat Conduction Problem (IHCP) consists in using discrete temperature measurements inside a solid or at one of its external boundary in order to recover a time and/or space boundary condition (in the present case the distribution of the cooling flux at the droplet impact). This inverse estimation is a not well-posed problem, which means that low magnitude perturbations in the temperature measurement (noisy temperature) can generate large deviations in the estimated wall heat flux. This is caused by the discrete features of measurements while the information that is looked for is a continuous function of time and/or space. The model has been reduced (from 3D to 2D) and a regularized least square estimator has been used to overcome this effect and to stabilize efficiently the inversion algorithm. A detailed description of the inversion method used to estimate the heat flux removed by the droplets is provided in Appendix C.

3.3.3. Uncertainty for heat flux estimation

The uncertainty resulting from the use of an inverse heat transfer algorithm cannot be directly assessed, even if the error in every
data (thermophysical properties and temperature measurement) is known. The estimation of the bias in the heat flux must be done using simulated temperature obtained after a direct simulation of system (5)(Gradec et al., 2009). The obtained temperature field knowing all inputs (boundary conditions, thermophysical properties) is then used as input data of the inverse model in order to compare the outcomes of the inversion procedure with the primary input data (i.e. boundary conditions (8)). From these tests, the uncertainty on $Q_w(t)$ was estimated to about $+0.02 \text{ mJ}$.

4. Results and Discussion

Measurements were carried out for three different sizes of droplets (80 µm, 133 µm and 180 µm). Droplet size is modified by changing the diameter of the injector orifice and the frequency of the droplet injection. In this study, the frequency ranges between 9 kHz and 12 kHz. Variations in frequency is not expected to play a important role in this range, since the characteristic time for thermal diffusion in the thickness of the nickel slab ($e^2/\alpha$) is much longer the droplet period (10 ms compared to 0.1 ms). The impact angle (angle between the droplet stream and the horizontal wall) was modified step by step from 10° to 80° while the injection velocity was fixed at about 10 m/s. The case of normal impacts was not considered because of the need to discriminate between primary and secondary droplets that can coalesce in this configuration. The temperature of the liquid in the injector body was set at 24 °C throughout the experiments.

4.1. Liquid phase heating

Fig. 7 shows the liquid temperature increase as a function of the Weber number at fixed wall temperature $T_w = 500 \degree C$. The Weber number, $Wen = \rho V^2 D_d/\gamma$, based on the normal component of the droplet velocity is changed with the angle of incidence of the droplet stream. When $Wen < 60$, the impact regime is a rebound, whereas the regime observed beyond this value of $Wen$ corresponds to the splashing. It can be noticed that the heating of the droplets is all the more important than the Weber number is high. The droplet variation in temperature increases progressively with the Weber number, without any discontinuity between the bouncing and the splashing regimes. However, it appears clearly that the Weber number alone is not sufficient to describe the heat transfer from the wall to droplets having different size and normal impact velocity.

For the lowest values of $Wen$, the droplet heating rises rapidly with the $Wen$, and reaches a plateau in the beginning of the splashing region. The same evolution was already observed in Castanet et al. (2009). A certain scattering of the measurements can be observed, mainly due to the measurement errors, which have been estimated (error bars in Fig. 7). However, this scattering is not only due to measurement inaccuracies: the stability of the measurement conditions is generally not perfect, in particular in the bouncing regime, since the wall surface can be altered by a thermal fatigue during the measurement.

The influence of wall temperature on the droplet heating was also investigated (Fig. 8). Experiments were performed for wall temperatures ranging from 430 °C to 680 °C, and three different Weber numbers (modified by changing the angle of incidence of the droplets) stream. Again, it is found that the liquid heating increases with the Weber number. In contrast, the wall temperature seems to have a very limited effect on the liquid increase in temperature, which was already noticed in Castanet et al. (2009). This

![Fig. 7](image1.png)

**Fig. 7.** Increase in the droplet temperature during an impact onto the heated wall as a function of the normal Weber number $Wen$ for different droplets diameters ($T_w = 500 \degree C$, $f$ ranges from 9000 to 12,000 Hz).

![Fig. 8](image2.png)

**Fig. 8.** Influence of the wall temperature on the droplet heating for three different Weber numbers, $f = 12,500 \: \text{Hz}, D_d = 139 \: \mu \text{m}$.

![Fig. 9](image3.png)

**Fig. 9.** Heat removed per droplet as a function of the wall temperature for different values of the normal Weber number ($D = 180 \: \mu \text{m}, f = 12 \: \text{kHz}$).
may be explained by the fact that the experiments were carried out beyond the Leidenfrost temperature, where there is no direct contact between the droplet and the wall.

4.2. Wall cooling

The temperature at the rear face of the nickel sample was recorded during the cooling once the heating by the inductor is switched off. The temperature was recorded two times by the IR camera, one time during the cooling of the sample in the presence of droplets and a second time during the same cooling in the absence of droplets. We have applied the inversion procedure presented in Section 3.3, in order to recover the rate of heat flux presented in Section 3.3, in order to recover the rate of heat flux and thus the efficiency of the heat transfer increase.

4.3. Energy balance of the droplet/wall interaction

The energy balance of the droplet/wall interaction can be written as the sum of three contributions, the sensible heat gained by the liquid, the phase change heat and the heat drained by the vapor flow:

\[ Q_w = Q_l + \Delta m \cdot (L_v + C_p \Delta T_v). \]  

This result has been already observed in several studies including [Bernardin et al., 1997] where the cooling of a heated target by monodisperse droplet streams was also investigated. As the Weber number increases, the spreading diameter of the droplets and hence the heat transfer are enhanced. Furthermore, the cooling efficiency is all the more important than the droplet size is small. When the droplets are getting smaller, the ratio surface/volume and thus the efficiency of the heat transfer increase.

\[ \Delta m = m - m_{0} \]  

The parameters \( T_s \) and \( \Delta T_v \) are not measured but in a first approach, \( T_s \) can be replaced by the boiling temperature of the liquid and \( \Delta T_v \) can be estimated by assuming \( \Delta T_v = T_f - T_b \). A model for the thickness of the vapor film squeezed between the wall and the droplet and for the flow field within this vapor film would be required to estimate more accurately the sensible heat \( C_p \Delta T_v \) gained by the vapor and \( T_s \), but this is not in the scope of the present paper. Introducing the Jakob number \( J_a = C_p \Delta T_v / L_v \):

\[ Q_w = Q_l + L_v \Delta m (1 + J_a) \]  

Finally, the mass of the droplet evaporated during its interaction with the wall can be expressed by:

\[ \Delta m / m \]  

![Fig. 10](image1.png)  

**Fig. 10.** Evolution of the cooling efficiency as a function of the normal Weber number, for different droplet sizes.

![Fig. 11](image2.png)  

**Fig. 11.** Evolution of the relative loss of mass \( \Delta m/m \) as a function of the normal Weber number for different droplet sizes (\( f \) ranges from 9500 Hz to 12,000 Hz).
\[ \Delta m = \frac{Q_w - mC_p\Delta T_i}{L_v(1 + ja) + C_p(T_b - T_a)} \]  
\( \text{(14)} \)

where \( T_a \) is the temperature of the liquid after the impingement.

Fig. 11 shows the evolution of \( \Delta m/m \) as function of the Weber number, for different droplet sizes. It can be seen that this parameter varies strongly with the droplet size. For the smallest droplets (\( D_d = 80 \mu m \)), the relative variation in mass is very significant as it can reach about 25%, while it is negligible for the largest droplets.

The uncertainty on the droplet temperature \( \Delta T \) can be evaluated by analyzing the different sources of uncertainties. From Eq. (2),

\[ \Delta T = \frac{\partial T}{\partial \beta} \beta + \frac{\partial T}{\partial T_0} \Delta T_0 + \frac{\partial T}{\partial R_0} \Delta R_0 + \frac{\partial T}{\partial R_f} \Delta R_f \]  
\( \text{(17)} \)

In this expression, \( T_0 \) is the reference temperature associated with \( R_0 \). The derivatives in Eq. (17) are calculated at the mean temperature encountered in this study (\( T = 22^\circ C \)). The contribution of the different sources of uncertainties is given in Table 1.

### 5. Conclusions

The implementation of innovative non-intrusive diagnostics allows investigating the different contributions to the energy balance in the Leidenfrost effect.

Infrared thermography combined with an inverse heat conduction model allowed estimating the heat flow rate removed from the wall by the impact of monodisperse droplet streams and consequently the heat removed per droplet.

Furthermore, 2cPILF thermometry was used to measure the increase in the temperature of the droplets during an impact and thus to determine the sensible heat gained by the liquid. The contribution of evaporation to the wall cooling was obtained from the closure of the energy balance. The main interest of this approach is that it is almost impossible to quantify directly the mass of liquid after the impact, since the droplets can be strongly deformed after impinging the wall. It was clearly observed that the main contribution to the cooling is the gain of sensible heat by the liquid in the case of the large droplets. When the droplet size decreases, the heat removed by evaporation becomes dominant. In all the cases, heat transfers increase with the normal Weber number. A better assessment of the evaporated mass would require quantifying the enthalpy of the vapor trapped between the droplet and the wall.

### Acknowledgements

This work has been supported by the French National Agency (ANR) in the frame of the research program IDHEAS (ANR-NT09 432160).

### Appendix A

The uncertainty on the droplet temperature \( \Delta T \) can be evaluated by analyzing the different sources of uncertainties. From Eq. (2),

\[ \Delta T = \frac{\partial T}{\partial \beta} \beta + \frac{\partial T}{\partial T_0} \Delta T_0 + \frac{\partial T}{\partial R_0} \Delta R_0 + \frac{\partial T}{\partial R_f} \Delta R_f \]  
\( \text{(17)} \)

In this expression, \( T_0 \) is the reference temperature associated with \( R_0 \). The derivatives in Eq. (17) are calculated at the mean temperature encountered in this study (\( T = 22^\circ C \)). The contribution of the different sources of uncertainties is given in Table 1.

### Appendix B

The Laplace (~) and Hankel (\( \sim \)) transforms are used so that Eq. (6) becomes:

\[ \frac{\partial^2 \theta}{\partial z^2} - \left( \varpi^2 + \frac{p^2}{a^2} \right) \theta = 0 \]  
\( \text{(18)} \)

where \( \theta = T - T_{\text{ref}} \), \( p \) is the Laplace parameter, \( x_n = u_n/R_0 \), \( u_n \) solutions of \( f_1(u_n) = 0 \) and:

\[ \bar{\theta}_n(z, t) = \int_0^\infty \theta(r, t, z) f_{0n}(x_n r) dr ; \quad \bar{\theta}_n(p, z) = \int_0^\infty \bar{\theta}_n(z, t) e^{-r^2 dt} \]  
\( \text{(19)} \)

Finally, the quadrupole method (Mailet et al., 2000) yields a linear relationship between the rear face temperature \( z = e \) and the cooling heat flux \( z = 0 \):

\[ \bar{\theta}^0_n(p) = -Z_n \left[ (p + ax_n^2) \bar{\theta}_{n,d}(p) + \overline{W}_{n}^0 \right] (p + ax_n^2) \theta_{n,\text{ref}} - \overline{W}_{n}^{\text{sum}}(p + ax_n^2) \theta_{n,\text{sum}} \]  
\( \text{(20)} \)

<table>
<thead>
<tr>
<th>( \Delta T_w ) (K)</th>
<th>( \Delta T_0 ) (K)</th>
<th>( \Delta R_0/R_0 )</th>
<th>( \Delta R_0/R_f )</th>
</tr>
</thead>
<tbody>
<tr>
<td>10</td>
<td>1</td>
<td>4%</td>
<td>3%</td>
</tr>
</tbody>
</table>

This work has been supported by the French National Agency (ANR) in the frame of the research program IDHEAS (ANR-NT09 432160).
Three transfer functions, in the Hankel–Laplace domain, explain the response of the rear face temperature:

- one impedance $\sim_{h}\tilde{Z}_{h}(p)$, which accounts for the effect of the droplet stream on the flux at the front face
- and two transmittances:

\[ \tilde{W}_{h}^{\text{num}}(p) \]

which stems from the initial thermal imbalance between the slab and ambient.

$\tilde{W}_{h}^{\text{num}}(p)$ that is associated to the conduction heat losses through the stumatite sample holder.

These transfer functions are:

\[ \tilde{Z}_{h}(p) = \left[ (h_{R}^{p} + h_{Q}^{p})\cosh(ke) + \left( \frac{h_{R}^{p}h_{Q}^{p}}{2k} + \lambda k \right) \sinh(ke) \right]^{-1}, \]  

(21)

\[ \tilde{W}_{h}^{\text{num}}(p) = \left( \cosh(ke) + \frac{h_{R}^{p}}{2k} \sinh(ke) \right) \tilde{h}_{h}(p), \]  

(22)

\[ \tilde{W}_{h}^{\text{num}}(p) = K_{\text{num}}R_{h} \tilde{J}_{h}(x_{R}R_{h}) \tilde{Z}_{h}(p). \]  

(23)

Numerical Laplace inversion of the preceding expressions allows calculating the Hankel transform of the rear face temperature field at any time:

\[ \tilde{\phi}_{h}(t) = \int_{0}^{t} \exp(-x_{h}(t-t')) \tilde{Z}_{h}(t-t') \tilde{\phi}_{h}(d't) \, dt' \]

\[ \tilde{\phi}_{h}^{\text{num}}(t) = \tilde{\phi}_{h}(t) + \tilde{W}_{h}^{\text{num}}(t) \tilde{\phi}_{h}(t) \]

(24)

The real temperature in the time–space domain can be finally obtained through the Hankel inversion of the previous equations:

\[ \tilde{\phi}_{h}^{\text{exp}}(r, t_{k}) = \frac{2}{R^{2}} \sum_{n=d}^{n_{h}} J_{n}(x_{h}R_{h}) \tilde{\phi}_{h}^{\text{num}}(t_{k}). \]  

(25)

Of course, because of its last term, Eq. (24) is only implicit and iterations starting from a case with the absence of conduction losses are possible.

Appendix C

Temperature at the rear of the wall $\tilde{\tau}(x, y, t)$ is measured by the infrared camera in a Cartesian reference frame while the input data $\tilde{\tau}(r, t)$ of the model must be known in a cylindrical reference frame. Thus, it requires the conversion and the averaging of the initial Cartesian temperature field in order to obtain the input data. As a consequence, only an averaged heat flux $\bar{q}_{d}(r, t)$ can be obtained, but not the local heat flux $q_{d}(x, y, t)$.

The first step of the inversion procedure consists in obtaining the input temperature by angular averaging the initial field:

\[ \hat{\phi}_{h}(r_{m}, t_{k}) = \frac{2}{n_{m}} \sum_{n=d}^{n_{h}} J_{n}(x_{h}R_{m}) \phi_{h}^{\text{num}}(t_{k}) \rightarrow \hat{\phi}_{h}(t_{k}) = \hat{\tau}_{h}(t_{k}), \]  

(26)

where $\hat{\phi}_{h}(t_{k})$ is the column vector of the rear face temperature differences for all observable pixels, of size $n_{m} n_{n}$ and $\hat{\tau}(t_{k})$ the column vector composed of its Hankel harmonics of orders 0 to $n_{h}$.

A Gauss Markov inversion of the preceding model, using the experimental temperature field, where $n_{m}$ pixels share the same radius $r_{m}$ allows estimating the vector of harmonics $\hat{\theta}$ at a given time:

\[ \hat{\theta}(t_{k}) = (X^{\Omega}X)^{-1}X^{\Omega}\theta^{\text{exp}}(t_{k}) \quad \text{with} \quad \Omega_{mn} = \delta_{mn}n_{m}. \]  

(27)

More details about the previous modelling can be found in Graeck et al. (2009) and Maillet et al. (2010).

Eq. (24), expressed in Hankel domain ($\sim$), is the second step of the inverse heat conduction problem (IHCP); the integral form can be expressed using a quadrature:

\[ \tilde{\alpha}_{h}^{\text{exp}}(t_{k}) = \tilde{\alpha}_{h}(t_{k}) - \exp(-x_{h}^{2}t_{k}) \tilde{W}_{h}^{\text{num}}(t_{k}) \]

\[ = \sum_{j=1}^{k} \exp(-x_{h}^{2}t_{j-1}) \tilde{Z}_{h}(t_{k}) \tilde{q}_{h,d}(t_{k}) \Delta t \]

\[ = \sum_{j=1}^{k} \tilde{S}_{h} \tilde{q}_{h,d}(t_{j}) \Delta t. \]  

(28)

where $\Delta t$ is the time step of the infrared camera (1/60 s here). It is the time regularization hyperparameter that has been chosen not too high, in order to get unbiased estimates of temperature and flux and not too low to prevent an explosion of the inversion because of the presence of noise in the temperature measurements.

Calling $\tilde{S}_{h}^{\text{exp}}$ the vector of the $n$th harmonics of the corrected experimental rear face temperature, calculated according to (27) and (28), a least square inversion results from model (28):

\[ \tilde{q}_{h,d} = (S^{\Omega}S)^{-1}S^{\Omega} \tilde{\alpha}_{h}^{\text{exp}}. \]  

(29)

Estimation of the wall flux $q_{d}$ is then made using a truncation of its spectrum to a maximum of $n_{h} + 1$ harmonics:

\[ \tilde{q}_{d}(r, t_{k}) = \frac{2}{R^{2}} \sum_{n=d}^{n_{h}} J_{n}(x_{h}R_{h}) \tilde{q}_{h,d}(t_{k}). \]  

(30)

In practice, only two harmonics ($n = 0$ and 1) are used, because of the weakly local effect of the droplet stream on the rear face temperature field (the disc is thin and highly diffusive). This number is the hyperparameter for space regularization. The total rate of heat flow $\tilde{Q}_{d}(t)$ is estimated by:

\[ \tilde{Q}_{d}(t) = 2\pi \int_{0}^{R} \tilde{q}_{d}(r, t) \, dr = 4\pi \sum_{n=d}^{n_{h}} J_{1}(x_{h}R_{h}) \tilde{q}_{h,d}(t_{k}). \]  

(31)

Finally, the energy removed from the wall at each droplet impact is computed from (31) knowing frequency $f_{\text{inj}}$ of the injector of the droplet stream:

\[ \tilde{Q}_{w}(t) = \frac{Q_{d}(t)}{f_{\text{inj}}}. \]  

(32)

References


