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Effects of structural defects and polarization charges in InGaN-based double-junction solar cell

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Abstract

The performance of a double heterojunction solar cell based on Indium Gallium Nitride (InGaN) including a tunnel junction was simulated. The most challenging aspects of InGaN solar cells development being the crystal polarization and structural defects detrimental effects, their impact on the solar cell performances has been investigated in detail. The solar cell simulation was performed using physical models and InGaN parameters extracted from experimental measurements. The optimum efficiency of the heterojunction solar cell was obtained using a multivariate optimization method which allows to simultaneously optimize eleven parameters. The optimum defect free efficiency obtained is 24.4% with a short circuit current $J_{SC} = 12.92 \text{ mA/cm}^2$, an open circuit voltage $V_{OC} = 2.29 \text{ V}$ and a fill factor $FF = 82.55\%$. The performances evolution as functions of the polarization and the defects types and parameters was studied from their maximum down to as low as a 2% efficiency.

Keywords: Simulation, Optimization, Solar cell, InGaN, Heterojunction, Tunnel, Polarization, Defects

1. Introduction

Among the materials currently being studied for high efficiency solar cells lies the Indium Gallium Nitride (InGaN) alloy. Indeed, InGaN possesses unique properties, such as its wide band-gap that is tunable from 0.7 eV to 3.4 eV [1]. It also presents strong interatomic bonds and high thermal conductivity [2], which makes it an ideal material for photovoltaic solar panels resistant to high energy radiation.

Several *p-n* and *p-i-n* homojunction and heterojunction structures were studied. The structures with low Indium composition have conversion efficiencies below 5% [3, 4]. Indeed, in practice, the InGaN layers with a high concentration of Indium are very difficult to achieve today. For relatively high-Indium-content structures, Doolittle *et al.* [5] theoretically showed that a *p-i-n* structure having an Indium composition of 60% exhibits

a maximum efficiency of about 19%. Our previous studies [6, 7] have shown that single junction structures have conversion efficiencies not exceeding 20%. The *p-n* structure presents an optimal efficiency of 17.8% while that of a *p-i-n* structure is 19.0%. These previous works, using empirical parameters and models, demonstrated that the theoretical maximum efficiency of a single junction is limited to about 20%.

In this paper, we simulate and optimize an InGaN-based dual-junction solar cell including a tunnel junction. This allows the conversion of different parts of the solar spectrum, thus enhancing the solar light absorption efficiency and the photocarriers transport. Photovoltaic panels using such solar cells can potentially achieve efficiencies exceeding those of systems that use conventional single junction solar cells.

In the past, only two papers on InGaN-based double-junction solar cells have been published. A first study published in 2008 [8] reports a dual-junction InGaN solar cell with a conversion efficiency of 35.1%. Another similar study published in 2016 [9] confirms these results with a roughly similar efficiency of

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34.9%. In these works the two subcells of the solar cell are studied separately. The tunnel junction between the two subcells is thus not included in the simulations. The effects related to the junction, such as the electrical resistance, the discontinuity of the bands, are thus ignored. Furthermore, the physical parameters of the material used, as well as the optical properties, notably the refractive index, are not reported. Furthermore, the authors did not take into account the structural defects in the InGaN layers and their polarization.

Moreover, the optimum efficiency obtained results from a parametric analysis which obviously does not allow a rigorous optimal efficiency of the solar cell. Indeed, parametric optimization can optimize only one parameter at a time, it thus does not allow to take into account the dependencies of optimized parameters.

We thus propose here a new approach which consists in performing a global simulation of the complete InGaN-based double-junction solar cell, including the tunnel junction, using rigorous physical models and empirical material parameters, and including structural defects and polarization effects.

Indeed InGaN is usually grown on lattice-mismatched Sapphire (Al₂O₃) or GaN/Sapphire substrates by MOVPE [10]. Thus, during the growth, a relaxation occurs beyond a certain thickness called "critical thickness"[11–13], resulting in a high density of dislocations and other structural defects in the InGaN layers, which impacts the solar cell performances. In addition to the defects, InGaN presents a spontaneous polarization due to the intrinsic asymmetry of the bonds in the wurtzite crystal structure, as well as a piezoelectric polarization mainly due to the lattice mismatch between the layers [14, 15].

The device is simulated in the framework of a drift-diffusion model using the Atlas device simulation framework from the Silvaco company. It is the most commonly used model to accurately describe the fundamental physics of semiconductor-based devices. For the carrier mobility, we use the classic model of Caughey-Thomas in which the carrier mobility is a function of doping and temperature [16]. In order to model the losses in the structure, we take into account the Shockley-Read-

Hall (SRH) [17] and direct Auger recombination models using Fermi statistics. An empirical model of the absorption of light using and interpolating experimental data was used [18] and combined to a precise model for the refractive index initially developed for GaN [19]. When the direct experimental data are not available, the InGaN parameters values are determined using Vegard's law linear interpolation between InN and GaN binaries values, except for the bandgap (E_g) and the electronic affinity (χ), where we used a quadratic modified Vegard's law with a bowing parameter of $b = 1.43 \text{ eV}$ for the bandgap and $b = 0.8 \text{ eV}$ for the electronic affinity[18]. All these physical models and dependent material parameters are detailed in our previous work [6, 7] and summarized in table 1.

The optimization of the device is done with rigorous mathematical methods. The objective is to find the optimal efficiency of the solar cell by simultaneously optimizing a given set of physical and geometrical parameters of the cell. We have, for this solar cell, optimized simultaneously eleven parameters. With that many parameters to optimize, a careful fine tuning of the previously used optimization algorithms was necessary in order to ensure their convergence while keeping a reasonable computing time, which can be as high as 30 hours per simulation.

This optimization work was realized with a Python package that we developed using the SciPy optimizers[20]. This packaged source code will be released to the photovoltaic community soon.

The paper is organized as follows. In the next section 2, we describe the solar cell structure as well as the InGaN structural properties that can limit the performance of the solar cell, including structural defects and spontaneous and piezoelectric polarizations. In section 3, we present the optimal results obtained and discuss the effects of the polarization and structural defects on the performance of the solar cell, before concluding.

Parameters [18]	InN	GaN
Bandgap E_g (eV)	0.7	3.42
Electron affinity χ (eV)	5.6	4.1
Hole and electron lifetime (ns)	1	1
Dielectric constant	15.3	8.9
N_c (cm^{-3})	9.1×10^{17}	2.3×10^{18}
N_v (cm^{-3})	5.3×10^{19}	4.6×10^{17}
Physical model		
Charge carrier mobility	Caughey-Thomas Model [21]	
Absorption Coefficient	Phenomenological Model [18]	
Refraction Index	Adachi Model [19]	

Table 1: Parameters of InN an GaN and significant physical model used in the simulations of the solar cell structure

2. Solar cell, polarization and defect modeling

2.1. Solar cell structure

The modeled solar cell is based on an InGaN-made double junction structure shown on Figure 1. This structure is composed of two vertically stacked $p-n$ junctions with a top wide bandgap cell, named *subcell-1* in the following, and a bottom narrower bandgap cell, named *subcell-2*, both being connected by a tunnel junction. The difference in the bandgaps is obtained by changing the Indium composition x of InGaN. The structure is studied under an AM1.5 illumination (about $100mW/cm^2$). The high-energy photons are absorbed by the top subcell (*subcell-1*) while those of lower energies are transmitted through the junction to be absorbed by the bottom subcell (*subcell-2*). The solar cell structure therefore absorbs a larger solar spectrum and can thus produce more energy than a single junction cell, provided the current can be matched between the two subcells.

Indeed, the two subcells being connected in series, the subcell that produces the smallest current limits the maximum current that can flow through the solar cell. Therefore, the optimum performance of the cell strongly depends on the subcells electrical interconnection and their respective optimal currents.

To achieve that fine tuning, the two subcells are connected by a tunnel junction, consisting of yet another $p-n$ junction built in the opposite direction to that of the two subcells which are designed for photocurrent generation. We designed here a $p-n$ junction whose Indium composition can be different from that

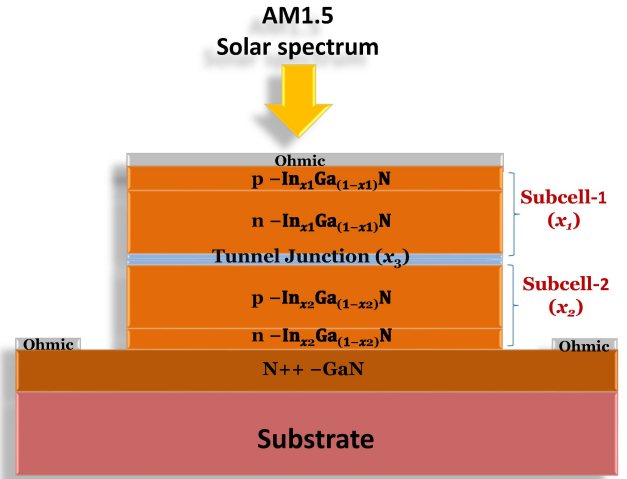


Figure 1: Solar cell structure with two $p-n$ junctions (*subcell-1* and *subcell-2*) connected by a tunnel junction. The cell is realized on a substrate via a heavily doped GaN buffer layer. x_1 represents the composition of the *subcell-1*, x_2 the Indium composition of *subcell-2* and x_3 the indium composition of the tunnel junction. The structure is under AM1.5 illumination.

of the subcells and that includes highly doped $p++$ and $n++$ layers. Therefore, This junction space-charge region is spatially narrow and its Fermi level is located both in the conduction and valence band: the bottom of the conduction band is on the same level as the top of the valence band. This way, the charge carriers can cross the potential barrier by tunneling.

Finally, the device proposed here can be realized in practice using the conventional InGaN growth on insulating substrates such as Sapphire (Al_2O_3) or SiC via a GaN or similar heavily doped buffer layer.

2.2. Spontaneous and piezoelectric polarizations

The InGaN crystal structure belongs to the wurtzite group [22]. It is a non-centrosymmetric structure having two types of polarities according to the arrangement of atoms of Gallium and Nitrogen during the epitaxial growth. They are called Ga-face polarity if the last epitaxially grown atomic layer is composed of Gallium atoms and N-face polarity if the last epitaxially grown atomic layer is composed of Nitrogen atoms [23]. This polarity will generate a spontaneous polarization (PSP) in the structure, spatially separating electrons and holes and thus creating an accumulation of charges at the InGaN surface.

	<i>InN</i>	<i>GaN</i>
$e_{31}(C/m^2)$	-0.41	-0.34
$e_{33}(C/m^2)$	0.81	0.67
$c_{13}(GPa)$	94	100
$c_{33}(GPa)$	200	392
$P_{sp}(C/m^2)$	-0.042	-0.034

Table 2: Experimental data of piezoelectric parameters and spontaneous polarization values for InN and GaN [26].

In addition to the spontaneous polarization, the lattice mismatched between different composition layers induces strain and thus piezoelectric polarization. This polarization also results in a displacement of electrons and holes at the surface of the layers and therefore may either reduce or enhance the total polarization [23]. Thus, without external polarization, the total polarization of the structure is equal to the sum of the spontaneous and piezoelectric polarizations.

Charge accumulation at the extremities of the structure can then change the electrical properties of the solar cell and the resulting photovoltaic performances. Indeed, if the electric field associated with the total polarization is in the same direction as the field generated in the space-charge region of the PN junction, the polarization could improve the performance of the cell, or reduce it otherwise.

The spontaneous polarization (P_{sp}) is quadratically interpolated from the modified Vegard's law with a bowing factor of $b = 0,038C/m^2$ [24, 25]:

$$P_{sp}^{InGaN} = xP_{sp}^{InN} + (1-x)P_{sp}^{GaN} - bx(1-x), \quad (1)$$

where x is the Indium composition of $In_xGa_{1-x}N$, P_{sp}^{InGaN} is the spontaneous polarization of $In_xGa_{1-x}N$ while P_{sp}^{InN} and P_{sp}^{GaN} are the spontaneous polarizations of InN and GaN binaries. Their values are given in Table 2.

The piezoelectric polarization (P_{pi}) is linearly interpolated by Vegard's law [24]:

$$P_{pi}^{InGaN} = xP_{pi}^{InN} + (1-x)P_{pi}^{GaN}, \quad (2)$$

where P_{pi}^{InGaN} is the piezoelectric polarization of $In_xGa_{1-x}N$,

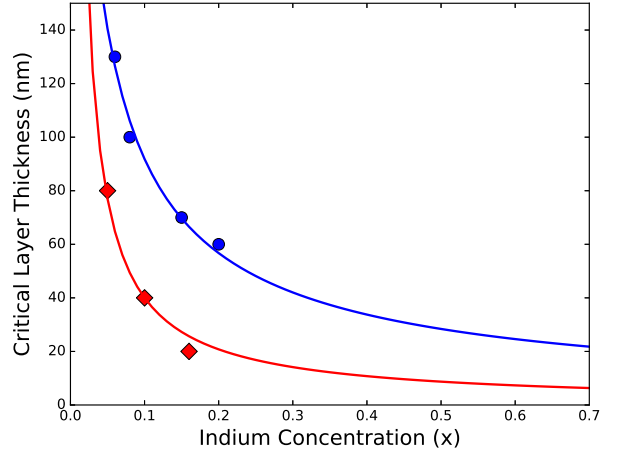


Figure 2: The critical layer thickness for $In_xGa_{(1-x)}N/GaN$ layer as a function of the Indium composition using a Fischer model [27]. Blue circles represent experimental measurements from [13] and the red diamonds the experimental measurements from [12].

P_{pi}^{InN} and P_{pi}^{GaN} are the piezoelectric polarizations of InN and GaN binaries. The latter are evaluated from piezoelectric constants and lattice parameters [22]:

$$P_{pi} = 2 \times \frac{a_s - a_0}{a_0} \times (e_{31} - \frac{c_{13}}{c_{33}} \times e_{33}), \quad (3)$$

where e_{31} and e_{33} are piezoelectric constants; c_{13} and c_{33} are elastic constants; a_s and a_0 are the lattice parameters of the substrate and of the material layer respectively. Their values are summarized in Table 2.

2.3. Critical thickness

In our modeling conditions, the spontaneous polarization is applied throughout the structure, while the piezoelectric polarization is applied only to the strained layers. In order to determine if a layer is strained or relaxed, we refer to the experimental measurements of critical thicknesses as a function of Indium composition. Since no measurements has been done for high Indium content layers, we have fitted the experimental data reported in [12] and [13] by a Fischer model [27] to extrapolate the critical thicknesses for high Indium content layers (Figure 2). For the extrapolation based on the measurements reported in [13], the critical thickness corresponding to an Indium composition of $x = 0.6$ would be about 30nm, while it would be closer

to 10nm for that based on [12]. Therefore, for our simulations, we considered that all layers with a thickness less than 10nm are strained and thus that layers having a thickness greater than 10nm are relaxed. This layer relaxation induces a high density of dislocations and other structural defects in the epitaxial layer.

2.4. Structural defects

Among the defects in III-Nitride alloys, and particularly evidenced in GaN, one can cite point antisites and interstitials defects, threading dislocations [28], stacking faults [29], stacking mismatch boundaries [30], grain boundaries [31], nanopipes and inversion domains[32]. All these defects cause perturbations in the periodicity of the crystal and affect the electronic properties of the devices. The defect whose associated energy lies in the bandgap, exchange with the conduction and valence bands through the emission and capture of electrons or/and holes. Garni *et al.* [33], showed that threading dislocations in GaN act as nonradioactive centers and diffusion centers in electron transport[34]. Moreover, Garni *et al.* also showed that these dislocations defects cause nonradioactive recombination which causes crystal heating up and thus significantly influences the performances of the device.

According to calculations by Neugebauer and Van de Walle [35], the formation energy of GaN point defects depends on the position of the Fermi level. The type of the defects usually observed in GaN thus depends on the doping type of the layer. This study shows that, for p-type doping, the Nitrogen vacancy have the lowest energy and are therefore most likely to be present in this layer while for n-type doping, Gallium vacancy are dominating. The Nitrogen vacancies being n-type and the Gallium vacancies being p-type, one can conclude that on p-type layers, the donor-type defects will be dominant, while, in n-type layers, the acceptor-type defects will.

Stampf *et al.* [36] also showed that, for InN, the Nitrogen vacancies have a lower energy than other defects in the p-type layers and thus are more likely to be present. It is then assumed that InGaN has the same types of defects as GaN and InN. We therefore consider in our simulations the donor-type defects in

p-type layers and acceptor-type defects in n-type layers.

Furthermore and to take into account the disorder due to dislocations, the defect density of states is defined as a continuous function of the energy. It is a combination of exponentially decaying band tail states and Gaussian distributions states within the bandgap. Band tail states represent the shallow states in the bandgap and Gaussian distributions represent the deep ones. They both impact the effective charge mobility through the capture and recombination of opposite charges. The total density of states is thus composed of four bands: two band tails and two bands of deep levels.

3. Influence of defect and polarization on efficiency

3.1. Solar cell optimization

The simulated and optimized cell structure is shown schematically in Figure 1. It includes three different Indium compositions: *subcell-1* has an x_1 Indium composition and *subcell-2* has an x_2 one while the junction tunnel can exhibit a yet different x_3 Indium composition.

The 11 parameters for which we sought the optimal values are the the 2 layer thicknesses and dopings for each PN junction along with their compositions. These 10 parameters are included globally in the optimization process along with the tunnel junction Indium composition. The thicknesses and dopings of the tunnel junction have not been included in the optimization process because our simulations have shown that they need to be as highly doped and thin as the technology allows it.

All the 11 optimization parameters have been constrained within a variation domain, which is the interval in which they are meaningful in physical and technological terms. The Indium compositions are allowed within the full range: $x_i \in [0, 1], \forall i \in \{1, 2, 3\}$. The InGaN layer n-type and p-type dopings (N_d and N_a) minimum allowed values are chosen to be the minimal residual doping of $10^{16}cm^{-3}$, as usually measured in this material. Their maximum values are, on their side, set to $10^{19}cm^{-3}$, which is the highest doping value used. To put it shortly: $N_d \in [10^{16}, 10^{19}]cm^{-3}$ and $N_a \in [10^{16}, 10^{19}]cm^{-3}$. For

Subcell 1	* $x_1 = 0.40$	* $L_p(\mu\text{m})$	0.01
		* $L_n(\mu\text{m})$	0.58
		* $N_a(\text{cm}^{-3})$	8.04×10^{18}
		* $N_d(\text{cm}^{-3})$	2.38×10^{16}
Tunnel cell	* $x_3 = 0.5$	$L_n(\mu\text{m})$	0.005
		$L_p(\mu\text{m})$	0.005
		$N_d(\text{cm}^{-3})$	5.00×10^{19}
		$N_a(\text{cm}^{-3})$	5.00×10^{19}
Subcell 2	* $x_2 = 0.59$	* $L_p(\mu\text{m})$	0.57
		* $L_n(\mu\text{m})$	0.01
		* $N_a(\text{cm}^{-3})$	2.51×10^{18}
		* $N_d(\text{cm}^{-3})$	1.00×10^{17}
GaN Layer		$L_n(\mu\text{m})$	0.16
		$N_d(\text{cm}^{-3})$	1.00×10^{19}

Table 3: Optimized parameters of the heterojunction InGaN solar cell. The parameters marked with an asterisk are the eleven simultaneously optimized parameters, the other parameters being fixed. x_1, x_2 and x_3 are the Indium compositions for *subcell-1*, *subcell-2* and tunnel junction respectively (see figure 1). L_p and L_n are the InGaN layer thicknesses of the P and N layers respectively. N_a and N_p are the dopings of the P and N layers respectively. The optimum efficiency found is 24.4%, associated to a $J_{SC} = 12.92\text{mA}/\text{cm}^2$ short circuit current, a $V_{OC} = 2.287\text{V}$ open circuit voltage and a $FF = 82.55\%$ fill factor.

the InGaN layer thicknesses (L_p and L_n), we chose to set the minimum values to 10nm and the maximum values to $1\mu\text{m}$. Therefore: $L_p \in [0.01, 1.00]\mu\text{m}$ and $L_n \in [0.01, 1.00]\mu\text{m}$.

Finally, for the optimizing algorithm to work as expected, the sought parameters should have similar variation ranges. Therefore, the Indium composition and the layer thicknesses are normalized to 1, their maximum value. For doping, the normalization is logarithmic (\log_{10}) and normalized to 19, its maximum value. The mathematical optimization methods used for optimization are detailed in our recent published work [6, 7]. Optimization is run several times with several starting points randomly chosen to ensure that the optimum obtained is absolute.

Using this procedure, a global maximum cell efficiency of 24.4% is found. Associated to this maximum efficiency, the dual junction solar cell main electrical characteristics at the optimum are the short circuit current $J_{SC} = 12.92\text{mA}/\text{cm}^2$, the open circuit voltage $V_{OC} = 2.287\text{V}$ and the fill factor $FF = 82.55\%$. The optimal values obtained for the free parameters are given in Table 3. Figure 3 shows the current-voltage and power-voltage characteristics of the this optimal solar cell.

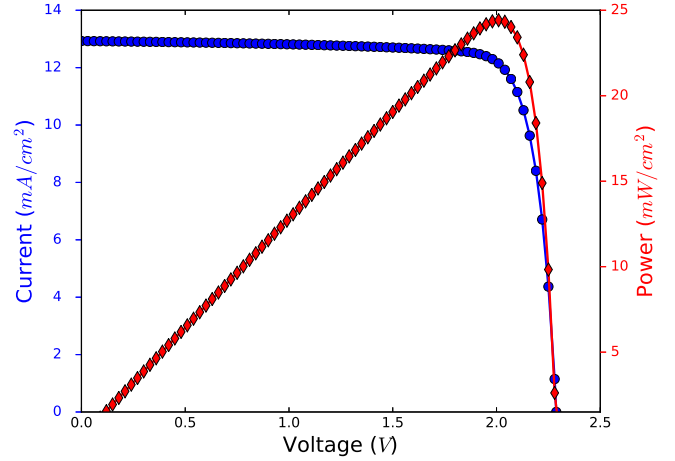


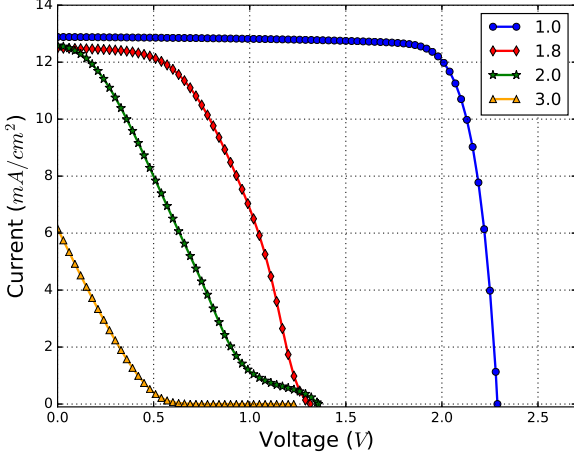
Figure 3: Current-voltage and power-voltage characteristics of the solar cell at the optimum.

By comparison, previous simulation works for InGaN-based double-junction solar cells obtained a higher efficiency [8, 9] of 35.1%. This is due to the fact that this work did not take into account the tunnel junction as itself. The two subcells were simulated separately and the effects of the series resistance and tunnel junction ignored.

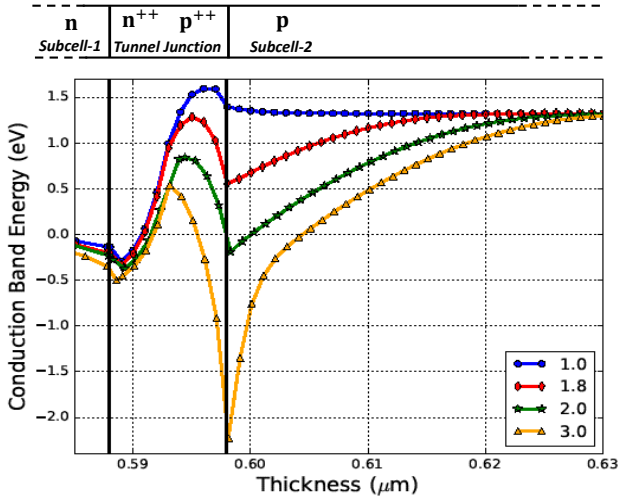
In addition, our work is based on an empirical conduction and absorption model directly deduced from experimental data, associated to a rigorous optimization procedure. Therefore, we expect the obtained optimal efficiency to be on one side actually optimal and, on the other side, not overestimated.

3.2. Spontaneous and Piezoelectric polarization

Once the optimal solar cell found, the effects of the polarization charges on its performances have been analyzed. As pointed out above, the spontaneous polarization is applied throughout the structure while the piezoelectric polarization is applied only to the constrained layers. *subcell-2* has an optimum thickness of $0.56\mu\text{m}$, and is thus relaxed, as deduced from figure 2 and previously discussed in section 2.3. The tunnel junction, at the bottom, has a thickness of 10nm, so it is a constrained layer. The *subcell-1* has a lattice parameter very close to that of the tunnel junction, it is also constrained. In short, the piezoelectric polarization is considered for the tunnel junction and *subcell-1* only.



(a) Current-voltage characteristics with different polarization charges.



(b) Conduction band diagrams in the tunnel junction with different polarization charges.

Figure 4: Impact of total polarization on the performances of the solar cell. The normalized value of 1.0 corresponds to the value in C/m^2 , obtained by the sum of the spontaneous and piezoelectric polarization given respectively by equations (1) and (2). To evidence the impact of a strong polarization, it is gradually increased up to 3 times the total polarization. *subcell-1* is on the left of the figure, while *subcell-2* is on its right, the tunnel junction being sandwiched between the two.

Figure 4(a) shows the I-V curves for solar cells with different polarization charge. The normalized value of 1.0 corresponds to the value in C/m^2 , obtained by the sum of the spontaneous and piezoelectric polarization given respectively by equations (1) and (2). For a small polarization, a normalized value of 1.0, the solar cell performances barely change, the optimal efficiency dropping from 24.4% to 23.3%. However, should the polarization charges be higher, a sharp decrease in the open circuit voltage V_{OC} would be seen while the short circuit current J_{SC} would remains constant. In addition, for high polarizations, a change in the shape of the curve would be observed, considerably degrading the fill factor.

To explain these behaviors, we have analyzed the conduction band in the tunnel junction with various polarization charges, as is shown in figure 4(b). The polarization charges bend the conduction band at the interface between the tunnel junction and *subcell-2*; the bend curvature increasing with the polarization. This causes charges to accumulate at this interface, which in turn create an electric field. This field, induced by the polarization, is in the direction opposite to that of the PN junction internal field. It thus tends to lower the total integrated electric field in the cell, which explains the decrease in the open circuit voltage V_{OC} .

Moreover, the presence of this induced field reveals a reverse biased diode, which opposes the efficient collection of the carriers. A threshold effect appears: when the induced field is less than the electric field of the cell, a deformation of the I-V curve is observed which severely degrades the fill factor and thus decreases the solar cell efficiency though no change on the short circuit current is J_{SC} is observed; on the contrary, if the induced field is largely greater than that of the cell, the deformation of the I-V curve also results a decrease of the J_{SC} . Indeed, the induced field opposes efficient collection of the carriers. This leads to a decrease of the current in the *subcell-2*. This latter then imposes its low current to the whole structure.

3.3. Deep and band tail defects

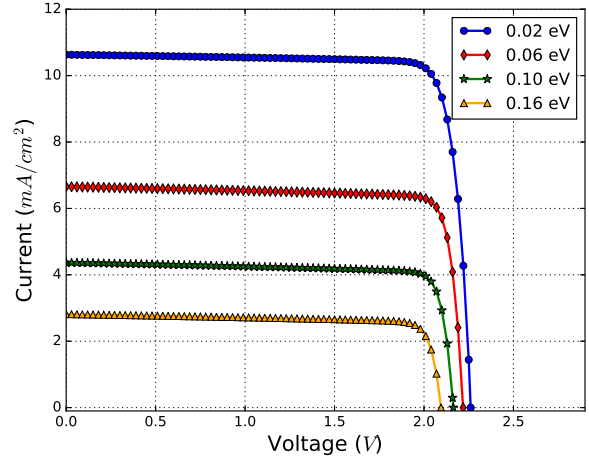
As for polarization, structural defects play a central role in an InGaN solar cell operation. The InGaN layer composing *subcell-2* being relaxed, the dislocations generated are similar to line charges and are an important source of defects, considerably reducing the mobility of the charge carriers and the efficiency of the solar cell.

In figure 5, the impact of the Urbach band tail energy on the cell performances was studied. In figure 6, it is the impact of the deep defects Gaussian distribution that was studied. Figures 5(a) and 6(a) show the different current-voltage characteristics of the solar cell, while figures 5(b) and 6(b) show the variations of J_{SC} and V_{OC} .

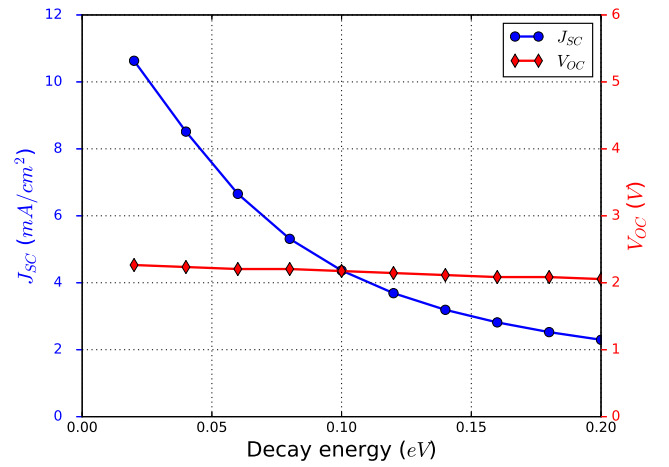
Band tails reflect the state of disorder of the material. The characteristic energies of the band tails and deep defects usually increase with the Indium composition. This is the reason why their influence on the solar cell performances was studied.

The increase of the Urbach energy as well as that of the decay energy (Figures 5(a) and 6(a)) leads to a decrease of the J_{SC} while the V_{OC} remains almost constant (Figures 5(b) and 6(b)). Indeed, an increase of the Urbach energy implies a reduction in the optical gap owing to the presence of states in the InGaN bandgap, leading to a higher probability of carrier recombination. In the same way, for the deep states, the increase of the decay energy also leads to an increase of the state density in the bandgap and therefore to a higher probability of recombination of the carriers. Both these effects thus result in the decrease of the short circuit current J_{SC} .

However, it can be noted the an increase in the disorder – in the Urbach energy– degrades the solar cell performances much more than a deep states decay energy increase. Indeed, for an Urbach energy of 60meV, the short circuit current J_{SC} drops from more than 10 to $6.65\text{mA}/\text{cm}^2$, while the variation of the decay energy, between 20meV and 160meV only induces a slight decrease in the same short circuit current J_{SC} from 12.8 down to $10.5\text{mA}/\text{cm}^2$. In InGaN solar cells, material disorder is therefore more detrimental to the performances than deep defects states.

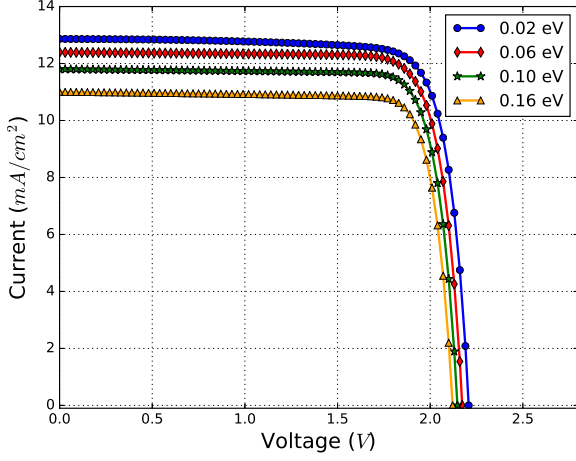


(a) Current-voltage characteristics with various Urbach energy of the band tail.

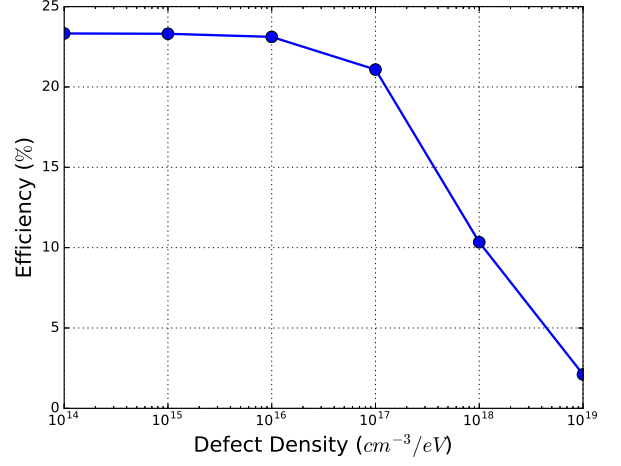


(b) Variation of J_{SC} and V_{OC} as a function of the Urbach energy of the band tail.

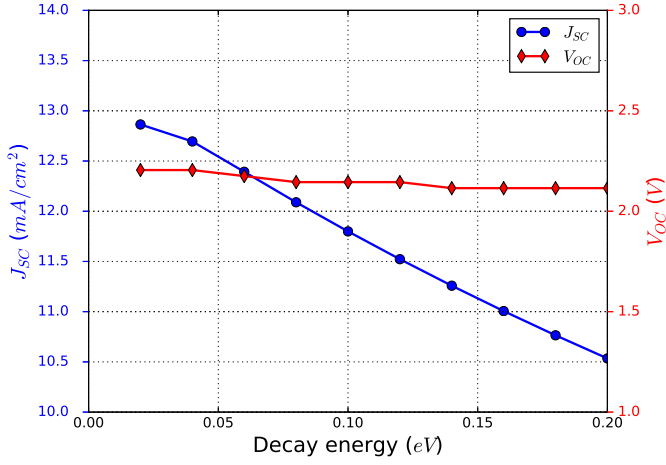
Figure 5: Impact of the band tail Urbach energy on the performances of the solar cell. The efficiency drops from 20.54% down to 4.89% for an Urbach energy varying from 0.02eV up to 0.16eV. The fill factor remains unchanged around 82%. The capture cross-sections are taken equal to $1 \times 10^{-14}\text{cm}^2$, the band tail state densities are taken equal to $2.3 \times 10^{18}\text{cm}^{-2}$ for the valence band and $4.6 \times 10^{19}\text{cm}^{-2}$ for the conduction band.



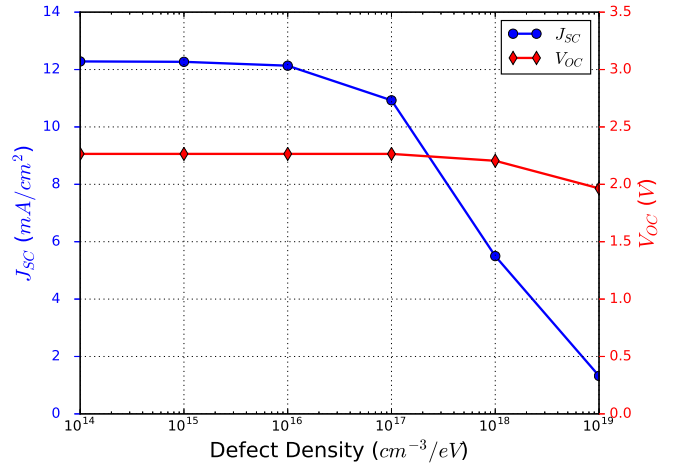
(a) Current-voltage characteristics with various decay energy of the Gaussian distribution.



(a) Solar cell efficiency as a function of the defect density of donor-type states in a Gaussian distribution.



(b) Variation of J_{SC} and V_{OC} as a function of the decay energy of the Gaussian distribution.



(b) Variation of J_{SC} and V_{OC} as a function of the defect density of donor-type states in a Gaussian distribution.

Figure 6: Impact of the Gaussian distribution decay energy on the performances of the solar cell. The efficiency is slightly degraded from 22.90% down to 19.45% for a defect decay energy varying from 0.02eV up to 0.20eV, the fill factor remaining unchanged around 81%. The capture cross-sections are taken equal to $1 \times 10^{-14} \text{cm}^2$. The energies that correspond to the Gaussian distribution peaks are fixed at 0.6eV, the Urbach energy of the band tail is set at 0.05eV, the defect density of donor-type is equal to $1 \times 10^{18} \text{cm}^{-3}$. The defect density of acceptor-type is equal to $1 \times 10^{17} \text{cm}^{-3}$.

Figure 7: Impact of the defect density of donor-type states in a Gaussian distribution on the performance of the solar cell. The defect density of acceptor-type states is constant and equal to 10^{17}cm^{-3} . The capture cross-sections are taken equal to 10^{-14}cm^2 . The energies that correspond to the Gaussian distribution peaks are set to 0.6eV.

Finally, in Figure 7, the effect of the density of donor-type defects was studied. Figure 7(a) shows its impact on the efficiency whereas figure 7(b) shows its impact on the short circuit current J_{SC} and the open circuit voltage V_{OC} . For donor-like defects, that are mostly found in the p-type layer, it can be observed that, when the defects density is less than the doping density in the layer, the efficiency remains close to its optimal value. On the contrary, as soon as the defects density exceeds the doping density, the efficiency of the cell dramatically decreases as shown in Figure 7(a). Indeed, by their ability to emit or capture free carriers, defects can compensate for intentionally introduced donor or acceptor levels. This reduces the mobility of the photogenerated carriers, thus reducing the short circuit current J_{SC} (Figure 7(b)) and the solar cell overall efficiency.

4. Conclusion

In this work, a detailed investigation on the performance of an InGaN-based double-junction solar cell was carried out. We have globally simulated the solar cell using empirical InGaN material parameters, to avoid any overestimation in the solar cell performances. In order to take into account the interdependence of the solar cell physical and geometrical parameters, ensuring the absoluteness of the optimized photovoltaic performances, the cell was optimized using a multivariate optimization algorithm that simultaneously optimizes eleven physical and geometrical parameters. We obtained an optimal efficiency of 24.4%, with a short circuit current $J_{SC} = 12.92\text{mA}/\text{cm}^2$, an open circuit voltage $V_{OC} = 2.287\text{V}$ and a fill factor $FF = 82.55\%$.

We quantitatively investigated the impact on the solar cell performances of the internal polarization and structural defects in InGaN. We have shown that the internal polarization reduces the performance of the cell by inducing an electric field which does not favor an efficient collection of photo-generated carriers; showing the importance of controlling this polarization when designing the lattice-mismatched solar cell heterostructure. We have also investigated the impact of structural defects

in InGaN, including disorder and deep defects, and correlated their effect to the InGaN doping concentration.

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