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Nonmonotonic aftereffect measurements in perpendicular synthetic ferrimagnetsT. Fache,¹ H. S. Tarazona,² J. Liu,³ G. L'vova,⁴ M. J. Applegate,³ J. C. Rojas-Sanchez,¹ S. Petit-Watlot,¹ C. V. Landauro,^{2,5} J. Quispe-Marcatoma,^{2,5} R. Morgunov,^{4,6} C. H. W. Barnes,³ and S. Mangin¹¹*Institut Jean Lamour, (UMR-CNRS 7198), Université de Lorraine, Nancy, France*²*Facultad de Ciencias Físicas, Universidad Nacional Mayor de San Marcos, P.O.-Box 14-0149, Lima 14, Peru*³*Cavendish Laboratory, University of Cambridge, J. J. Thomson Avenue, Cambridge CB3 0HE, United Kingdom*⁴*Institute of Problems of Chemical Physics, 142432, Chernogolovka, Russia*⁵*Centro de Investigaciones Tecnológicas, Biomédicas y Medioambientales, Calle José Santos Chocano 199, Bellavista, Callao-Peru*⁶*Tambov State Technical University, 392000, Tambov, Russia*

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Aftereffect measurements have been performed on synthetic ferrimagnets showing strong perpendicular magnetic anisotropy, namely (Co/Pt)/Ir/(Co/Pt), by measuring the magnetization of the sample as a function of time for several different applied magnetic fields. Unexpected magnetic relaxation has been observed. Indeed, for some particular applied magnetic fields, the magnetization as a function of time first plummets and then increases. This nonmonotonic aftereffect can be understood by considering the possible magnetic states and the transitions between those states. This understanding has been confirmed and detailed using magneto-optical Kerr effect (MOKE) microscopy. Indeed, the measurements show nucleation and propagation of a transient metastable magnetic configuration. Furthermore, we were able to obtain a good qualitative agreement between a simple one-dimensional model and the experimental observation. We could strongly support the hypothesis that this peculiar nonmonotonic behavior could be a very general feature that should be observed in any antiferromagnetically coupled system (synthetic ferrimagnets, synthetic antiferromagnets) with perpendicular magnetic anisotropy.

DOI: [10.1103/PhysRevB.98.064410](https://doi.org/10.1103/PhysRevB.98.064410)**I. INTRODUCTION**

Perpendicularly magnetized synthetic ferrimagnets made of two antiferromagnetically exchange-coupled magnetic layers are very attractive materials for spintronic devices, and have gathered a great deal of attention in the past decades. Numerous materials, such as [Pt/Co]Ru[Co/Pt] [1], CoFeB/(Ta or MgO)/CoFeB [2–4], [Co/Pd]Ru[Co/Pd] [4], or even Heusler alloys [5] or graphene [6] are suitable for the growth of these multilayers that consist of two ferromagnetic layers separated by a nonmagnetic spacer providing an antiferromagnetic coupling between them. Since there are plenty of useful applications to these magnetic structures, various studies explore in details the properties of these materials that are nowadays routinely used in magnetic memory devices, mainly as pinned layers in perpendicular magnetic tunnel junctions [7–9]. Another main application field of synthetic ferrimagnets is the concept of racetrack memories [10] developed by IBM in the past decades, using the magnetic structure to increase the domain wall speeds [11]. Other applications of synthetic ferrimagnets include also superconducting devices [12], sensors [13–17], and even cells for neuromorphic computing. The dynamic properties of the magnetization configurations in these systems can thus be of great interest in order to consider the stability of such layers when subjected to a magnetic field and/or temperature.

The evolution of the magnetization of thin films as a function of time is a very challenging issue that cannot be tackled by finite element algorithms due to the sizes and timescales involved in such phenomena. Neither is it possible

to address that question using exact theoretical methods, since the analytical models cannot take into account all the interactions of such large samples, and all the defects of atomic and magnetic structures which can hardly be avoided. Several approaches and models provide clues in order to tackle this issue considering the type of materials involved [18–25]. These methods and macroscopic descriptions predict a continuous switching of the magnetization from one state to another, leading to monotonic evolution of the magnetization as a function of time. Since aftereffect measurements probe the evolution of the magnetic system under thermal activation, the system will reach a lower energy state during the measurement. We then expect the magnetic system to evolve monotonically toward the most stable state. One would then predict a usual exponential decay corresponding to the crossing of a single energy barrier, as it has already been studied in Ref. [26].

In this paper, we display magnetic relaxation curves that show a nonmonotonic behavior when relaxing from an antiparallel state to the other one. Furthermore, since we attribute that unusual behavior to the energetic landscape that we obtain in our case, we can also link the existence of such a transition to the temperature, which plays a well-known role on the magnetic configuration of the sample [2].

II. METHODS AND SAMPLES

The samples involved in this study are Co/Pt bilayers separated by an Ir spacer. The exact stacks are silicon/Ta(3 nm)/Pt(3.2 nm)/Co(1 nm)/Ir(1.4 nm)/Co(9 nm)/

Pt(3.2 nm) (which will be referred as Pt/Co/Ir/Co/Pt sample 1), silicon/Ta(3 nm)/Pt(3 nm)/Co(0.7 nm)/Ir(1.4 nm)/Co (1 nm)/Pt(3 nm) (sample 2), and silicon/Ta(3nm)/Pt(3 nm)/Co(0.95 nm)/Ir(1.4 nm)/Co(1 nm)/Pt(3 nm) (sample 3). We deposited all the layers using a dc magnetron sputtering with elemental sources at room temperature. The thicknesses of the cobalt layers as well as the platinum layers that share an interface with the cobalt are chosen so as to obtain perpendicular magnetic anisotropy (PMA). The iridium spacer thickness is the most sensitive parameter of this sample: it determines the coupling between the ferromagnetic layers, by the RKKY interaction. We have measured that by changing the iridium thickness in our structure; we can obtain a local maximum antiferromagnetic coupling for a thickness of 14 Å. This coupling has been widely investigated, and the thicknesses that we found experimentally that maximize the antiferromagnetic coupling are very similar to what can be found in the literature [27]. We can point out that the global maximum of antiferromagnetic coupling is reached around 6 Å. Nonetheless, we chose to grow a 14-Å-thick iridium spacer so as to avoid the contribution of the roughness of our sample. The size of the samples went from $8 \times 10 \text{ mm}^2$ for the MOKE and MOKE microscopy measurements, whereas we cut these samples in $4 \times 4 \text{ mm}^2$ squares in order to perform the SQUID measurements.

Since we have two antiferromagnetically coupled cobalt layers, there are four accessible magnetic states given the combinations of orientations of both of the magnetic layers. In this paper, we will call P^+ (respectively, P^-) the parallel configuration where both magnetizations are oriented up (respectively, down), and AP^+ (respectively, AP^-) the antiparallel configuration of larger (respectively, smaller) net magnetization.

The relaxation measurements that we performed were carried out in the following conditions. First, the sample temperature is stabilized at a temperature under $H = +1000 \text{ Oe}$ in order to saturate the sample in the P^+ state. Then, we apply a field H^* close to the switching field deduced from the magnetic hysteresis obtained from SQUID measurements. It takes 60 s to stabilize the field, and then we could acquire the data corresponding to the net magnetization of the whole sample using a SQUID measurement.

The MOKE microscopy was performed at room temperature using a MOKE microscope with tuneable region of interest [28]. We usually choose the maximum image size in order to catch the magnetic signal of a surface of $100 \times 100 \mu\text{m}^2$. The low-temperature MOKE measurements were carried out in a handmade setup at the Cavendish Laboratory, where the size of the probed area is $3 \times 3 \mu\text{m}^2$. The temperature was set precisely with helium cooling, making it possible to achieve temperatures in the range of 1.6 K.

III. ROOM-TEMPERATURE MAGNETIC RELAXATION

The first set of data that we have recorded corresponds to sample 1. We performed all the measurements concerning that sample at room temperature. We display in Fig. 1 the hysteresis loop corresponding to this sample. We can observe that the four magnetic states are accessible by sweeping the field. However, a closer look shows us that it is difficult, at this sweeping rate,

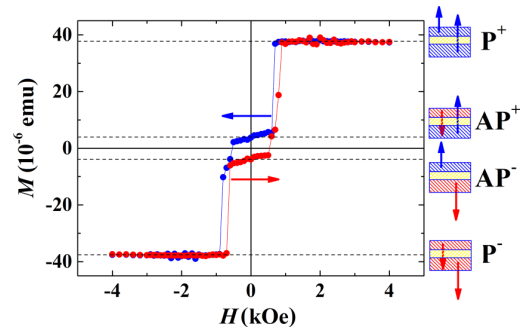


FIG. 1. Magnetic characteristics of Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.9 nm)/Pt: Magnetization M as a function of the applied magnetic field H at room temperature.

to access those states in a single sweeping from negative to positive fields, or vice versa: for an increasing (respectively, decreasing) field, the magnetization jumps directly from the AP^- (respectively, AP^+) configuration to the P^+ (respectively, P^-) one.

Our following work stems from that special feature of the hysteresis loop. We set the field around this AP^+ to P^- transition in order to see if we could observe a relaxation between the AP^+ to AP^- configurations. We disclose the relaxation curves obtained in Fig. 2. We can distinguish two types of relaxation. For H^* applied magnetic fields between -500 and -550 Oe , we can observe a relaxation between states AP^+ and AP^- , with a monotonic decrease of magnetization as a function of time. For lower fields, between -600 and -650 Oe , we can see a very different behavior that displays an increase towards the AP^- state. This feature is counterintuitive given the fact that the magnetization is initially set to the AP^+ state. This is a sign that in less than 60 s, which is the time for the field to stabilize, the magnetization decreases from AP^+ to a level that is lower than AP^- .

In order to have an idea of the phenomena involved in the unusual relaxation curves that we are able to obtain, we used MOKE microscopy to image the magnetic configurations. We could sweep the applied magnetic field until observing nucleation and propagation of magnetic domains. One can understand from the steepness of the transition $AP^+ \rightarrow AP^- \rightarrow P^-$

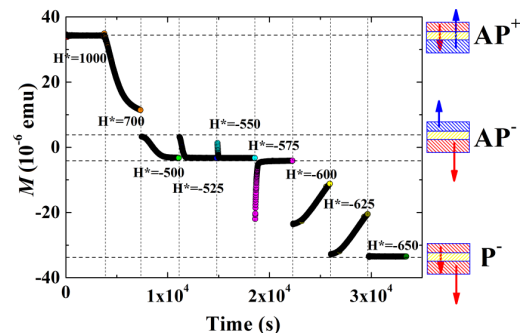


FIG. 2. Relaxation curves at room temperature under several magnetic fields applied on sample 1, Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.9 nm)/Pt: magnetization as a function of time during 3600 seconds for fields ranging from 1000 to -700 Oe

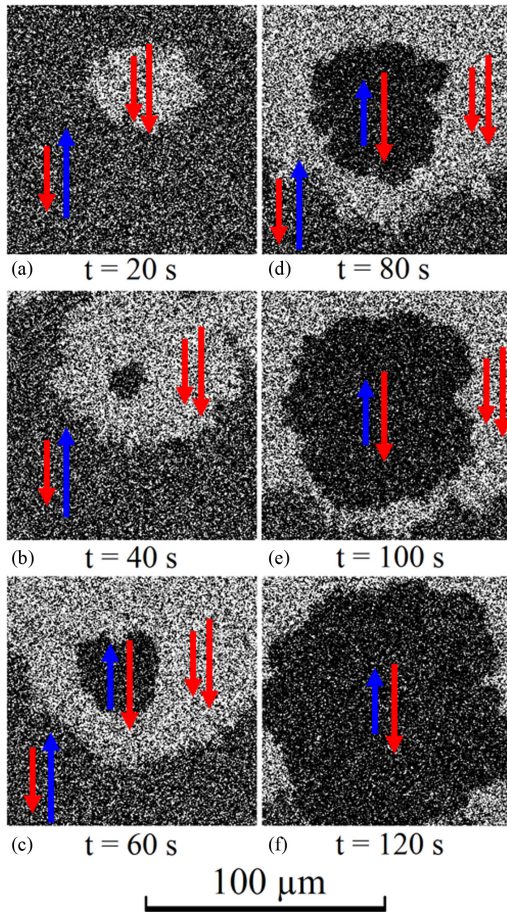


FIG. 3. MOKE imaging of the $AP^+ \rightarrow P^- \rightarrow AP^-$ transition at successive times at a constant applied magnetic field $H^* = -60$ mT after 20 (a), 40 (b), 60 (c), 80 (d), 100 (e), and 120 s (f). The shade of grey on the sides of images (a)–(c) represents the AP^+ magnetic configuration (initial state); the light shade of grey appearing in the middle of images (a)–(c) and on the sides of (d)–(e) represents the P^- magnetic configuration (transient state); the darker shade of grey growing in the middle of images (b)–(f) represents the AP^- magnetic configuration (final state).

in Fig. 2 that the field-induced magnetization switching is very fast. We could take a movie of this reversal on a very small timescale and extract several images showing the relaxation process.

The images that we get show clearly a two-step process: for a constant value of applied magnetic field (-60 mT), starting from AP^+ configuration (middle grey $\downarrow\uparrow$ domain on the sides), the nucleation and slow spread of P^- domains (lightest shade of grey $\downarrow\downarrow$) occur [Figs. 3(a)–3(c)]. Then an AP^- domain (darkest shade of grey $\uparrow\downarrow$) can nucleate from within the P^- domain and grow slowly [Figs. 3(b)–3(f)]. As time goes by, the initial AP^+ domains disappear quite fast, leading to a large proportion of P^- domains, and afterwards the P^- domains are slowly replaced by the final AP^- configuration.

In the MOKE microscopy that we disclose in Fig. 3, we can see what happens around one nucleation point. It turns out that there are very few nuclei like this one in our samples.

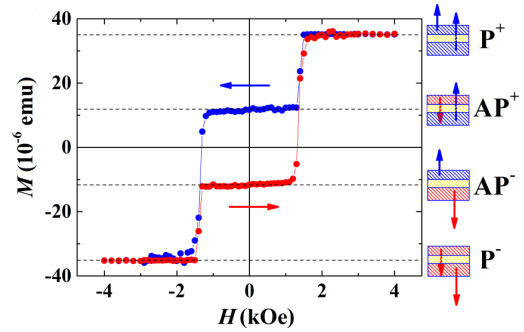


FIG. 4. Magnetic characteristics of Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.7 nm)/Pt: magnetization M as a function of the applied magnetic field H at a temperature of 100 K.

We can only find a few of them (less than 5) in our sample that is 10×8 mm² large. Consequently, the phenomenon that we are studying is starting with these nucleations, which are very sensitive to the initial conditions of the experiment, and is evolving according to the domain wall propagation throughout the sample.

The propagation of the domain walls that we are observing actually only consists in a domain wall in one cobalt layer: $AP^+ \rightarrow P^-$ corresponds to the switching of the lower cobalt layer magnetization, and $P^- \rightarrow AP^-$ corresponds to the switching of the upper cobalt layer magnetization. These domain wall motions are easily stopped by pinning points and spread slowly, compared to the $AP^+ \leftrightarrow AP^-$ transitions that can be observed for other temperatures and fields. In the latter case, both magnetizations are switching at the same time and pinning points must be at the same location for both cobalt layers to stop the motion, which is very unlikely. However, we can see that the shapes of the domains that are shown at different times in Fig. 3 are mostly round and smooth, indicating that the film seems homogenous with very few defects. We can consider that from these similar behaviors of propagations for transitions $AP^+ \rightarrow P^-$ and $P^- \rightarrow AP^-$ stem the slow nonmonotonic relaxation that we report in this article.

The dynamics of the phenomenon that we report here is very sensitive to the applied magnetic field. The most important feature that must be pointed out is that the phenomenon is qualitatively the same for different values of the field, changing only the timescale of the transition.

IV. LOW-TEMPERATURE MAGNETIC RELAXATION

In order to have a better understanding of the phenomenon that is occurring in these experiments, we decided to grow the same kind of samples with different cobalt thicknesses. This would result in a change in the hysteresis loop concerning the magnetization of the different states, and the values of the magnetic fields needed to observe the transition between given states. We could then use the knowledge of the magnetic behavior of these synthetic ferrimagnets with respect to the temperature [2] in order to obtain a hysteresis cycle that would compare well with the one of sample 1 at room temperature (Fig. 1). Thus we grew sample 2 that displays the magnetic behavior depicted in Fig. 4.

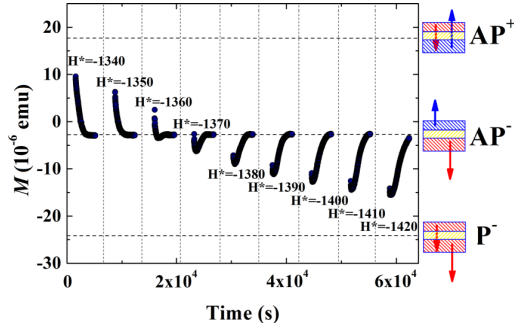


FIG. 5. Relaxation curves at a temperature of 100 K under several magnetic fields applied on sample 2, Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.7 nm)/Pt: magnetization as a function of time during 3600 seconds for fields ranging from -1340 to -1420 Oe. All field values of H^* written in the figure are in oersted.

In that case, again, the transition between states AP^+ and AP^- while decreasing the field does not seem to enable a stable AP^- configuration and the magnetization jumps nearly directly to the P^- state. The corresponding relaxation curves at low temperature (150 K) presented in Fig. 5 show the same kind of nonmonotonic behavior.

We can even see more precisely the magnetization minimum reached a few minutes after the beginning of each measurement. We provide a zoom of the most significant curves in Fig. 6 for two different values of the field.

It is important to notice that we could not get the very first values of the magnetization because we need to wait for the magnetic field to be stable in order to get a precise measurement of the magnetization of the sample. Thus the first point of the curves should correspond to a magnetization of AP^+ . We can especially notice on the relaxation at $H^* = -1350$ Oe that the magnetization decreases from a value that is higher than the level of the AP^- state.

For this low-temperature observation, we also wanted to be certain of the local phenomenon occurring. We performed local MOKE probing of the magnetization on sample 3, showing the same type of hysteresis loop with what we present above at a temperature of 232 K, and we could observe the evolution of

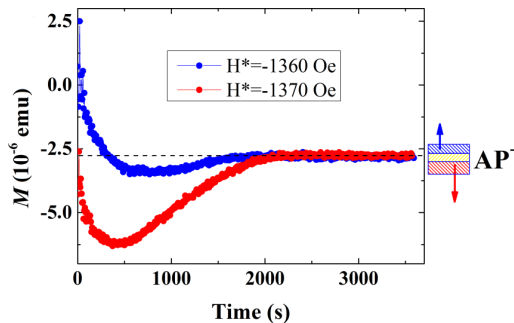


FIG. 6. Relaxation curves at a temperature of 100 K under several magnetic fields applied on sample 2 Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.7 nm)/Pt: magnetization as a function of time during 3600 seconds for $H^* = -1360$ Oe (blue) and $H^* = -1370$ Oe (red).

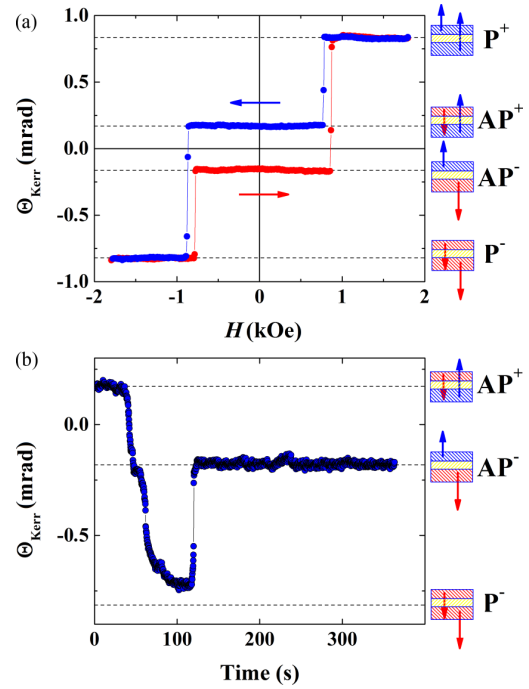


FIG. 7. (a) Magnetic characteristics of Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.95 nm)/Pt: Kerr signal (Θ_{Kerr}) as a function of the applied magnetic field H at a temperature of 232 K. (b) Local relaxation curve at a temperature of 232 K under an applied magnetic field $H^* = 880$ Oe. MOKE signal (Θ_{Kerr}) as a function of time during 350 s corresponding to a $30 \mu\text{m}^2$ part of sample 3 Pt/Co(1 nm)/Ir(1.4 nm)/Co(0.95 nm)/Pt.

the magnetic configuration with respect to time that we disclose in Fig. 7.

This local measurement of the magnetic configuration gives us the proof that the magnetization switches completely from the AP^+ state to the P^- state in a first stage, and afterwards another transition between the P^- state and the AP^- one occurs with a delay that depends on the position. We can especially pay attention to the fact that the sooner the first transition the shorter the delay between the $AP^+ \rightarrow P^-$ and the $P^- \rightarrow AP^-$ transitions. We can link that observation to Fig. 3 where we can clearly see that close to a nucleation center, both transitions happen in a small spatial range, corresponding to a short delay [Fig. 3(a)], whereas far from nucleation centers, the first transition $AP^+ \rightarrow AP^-$ happens long before the $P^- \rightarrow AP^-$ transition [Figs. 3(b)–3(d)].

V. TWO-SPIN MODEL

The transitions that are observed thanks to MOKE microscopy (Fig. 3) show that the magnetic reversal occurs from very few nucleation points, and propagates through slow domain wall motion. The fact that, for samples of around 80 nm^2 , only one or two nucleation points can be found leads us to think that this phenomenon is dominated by propagation. Furthermore, the nearly perfect symmetry of the bubbles that are observed during the relaxation support the idea that the magnetic properties of the film are very homogeneous. Thus

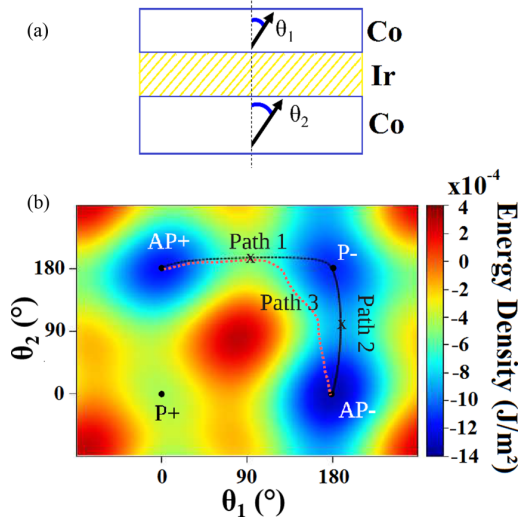


FIG. 8. (a) Definition of the angles θ_1 and θ_2 between the magnetization of top and bottom cobalt layers, and the perpendicular-to-the-plane direction. (b) Energy landscape of the two-macrospin model of the SAF bilayer for an applied magnetic field of 1500 Oe. Paths 1 and 2 show the two-step transition $AP^+ \rightarrow P^- \rightarrow AP^-$, whereas path 3 represents the direct $AP^+ \rightarrow AP^-$ transition.

we chose to use a two-spin model for this material, with homogeneous properties.

In order to understand the physical origin of this two-step relaxation, we established an energetic model based on a macrospin approximation for each magnetic layer. Even though this approach does not provide an exact picture of the relaxation experiment, it can still give some insight on the main differences between the direct $AP^+ \rightarrow AP^-$ transition and the consecutive $AP^+ \rightarrow P^- \rightarrow AP^-$ reversals.

Within the macrospin approach, the magnetization of the system can be entirely described by four independent angles ($\theta_1, \phi_1, \theta_2, \phi_2$), which define the magnetization directions in polar coordinates. However, due to symmetry reasons, only cases where the magnetizations lie in the same vertical plane describe equilibrium states. Thus it is sufficient to consider the polar angles θ_1 and θ_2 , which describe the orientation of the magnetization of the top and bottom layers with respect to the perpendicular to the film plane axis, respectively [see Fig. 8(a)]. Then, the energy of the system for any configuration of these angles can be written as the sum of the three contributions (anisotropy, Zeeman, and interlayer exchange coupling) detailed as follows:

$$\begin{aligned} E_A &= -K_1^{\text{eff}} \cos^2 \theta_1 - K_2^{\text{eff}} \cos^2 \theta_2, \\ E_Z &= -\mu_0 H (t_1 M_1 \cos \theta_1 + t_2 M_2 \cos \theta_2), \\ E_{EX} &= -J_{\text{ex}} \cos(\theta_1 - \theta_2), \end{aligned} \quad (1)$$

where t_1 and t_2 are the thickness of the top and bottom layer and S becomes the normal surface area of each magnetic layer. Here, $K_i^{\text{eff}} = \mu_0 t_i M_i^2 / 2 - K_{si}$ ($i = 1, 2$).

In order to take into account, the perpendicular magnetic anisotropy, we can use the fact that $K_i^{\text{eff}} < 0$. In addition, considering that the external field H is applied perpendicular to the film plane, it is straightforward to demonstrate that the

TABLE I. Values of the magnetic parameters chosen for the calculations.

	t (nm)	M (kA m ⁻¹)	Ks (mJ m ⁻²)
Co ₁ (top)	0.7	1300	1.1
Co ₂ (bottom)	1.0	1400	1.9

energy given by Eq. (1) presents four minima corresponding to the four accessible states of the magnetization shown above in the hysteresis loops. In order to obtain a picture of the energy landscape corresponding to the hysteresis loop of sample 2 in Fig. 4, we used the magnetic parameters listed in Table I. The interlayer exchange coupling constant was obtained from the switching condition, $|J_{\text{ex}}| = 2t_2 M_2 B_C$, with $B_C = 1200$ Oe and thus evaluated as $J_{\text{ex}} = -0.2352$ mJ m⁻². The magnetocrystalline surface anisotropy constants were estimated from Koplak's calculations for hysteresis loops of type II [2], which is given by

$$K_2^{\text{eff}} \frac{t_1 M_1}{t_2 M_2} - 2J_{\text{ex}} \frac{t_1 M_1 - t_2 M_2}{t_2 M_2} < K_1^{\text{eff}} < 2J_{\text{ex}} \frac{t_1 M_1 - t_2 M_2}{t_2 M_2}.$$

Using these parameters, we can plot the energy surface for any value of the applied field H , and thus analyze the thermally induced transitions between any two states. We studied in detail the case of the $AP^+ \rightarrow AP^-$ transition. We plotted the energy landscape for $H = -1500$ Oe in Fig. 8(b). In this case, the global minimum of the energy corresponds to the AP^- configuration, but there is no path going from the AP^+ to the AP^- without any large energy barrier. We consider that the system will evolve to the stable state AP^- through the minimum energy path (MEP). Although there is an infinite number of paths between these two states we define the minimum energy path as the path of highest statistical weight for transitions. Since the MEPs are very important in our theoretical analysis, we have implemented a procedure to calculate it numerically following the nudged elastic band method (NEB) [29].

Two possibilities are relevant for this transition. On the one hand, the combination of paths 1 and 2 (i.e., the $AP^+ \rightarrow P^-$ and $P^- \rightarrow AP^-$ successive transitions) gives the lowest energy transition. The combination of these paths becomes a two-step process. On the other hand, the transition following path 3, which is going directly from AP^+ to AP^- state, is worth being considered. This path has only one energy barrier given by the saddle point on path 1. Consequently, it is only qualitative, since the only parameter that matters is the height of the energy barrier of the transition, and the fact that the variation of the energy as a function of curvilinear coordinates is monotonic between the initial state and the maximum of energy of this path. We represent the energy as a function of the curvilinear coordinates for these paths in Fig. 9.

Now we assume that the thermal fluctuations may cause spontaneous jumps of the system from one metastable state to another to undergo a type of Brownian rotation and consequent relaxation. In order to calculate the dynamics of the transition occurring in our case, we propose a model that takes into account the probability of switching as a function of the transition paths drawn in Fig. 9.

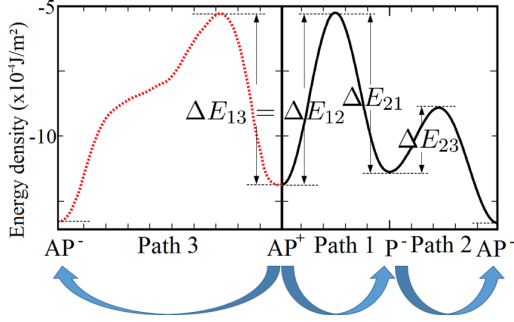


FIG. 9. Energetic paths corresponding to the two considered transitions between AP^+ and AP^- states: energy of the system as a function of curvilinear coordinates of the transition paths drawn in Fig. 8(b).

We considered the trilayer Co/Ir/Co systems as an assembly of small Co/Ir/Co systems characterized by surface area S whose energy is given by Eq. (1). Each system can be found in one of the configurations amongst AP^+ , AP^- , and P^- with probabilities P_{AP^+} , P_{AP^-} , and P_{P^-} , respectively. They must satisfy the equation

$$P_{AP^+} + P_{AP^-} + P_{P^-} = 1.$$

We then assume that the energy barriers that are given by the paths of Fig. 9 define the probability of transition between the different states. We can write down a differential system that defines the probabilities P_{AP^+} , P_{AP^-} , and P_{P^-} . Because the jumps from one to other state are Markov processes we can use the master equation, which describes the temporal evolution of the probabilities, calling $\mathbf{X} = \begin{pmatrix} P_{AP^+} \\ P_{AP^-} \\ P_{P^-} \end{pmatrix}$, and ω_{ij} the transition terms from state i towards j :

$$\frac{\partial \mathbf{X}}{\partial t} = \begin{pmatrix} -(\omega_{12} + \omega_{13}) & \omega_{21} & \omega_{31} \\ \omega_{12} & -(\omega_{21} + \omega_{23}) & \omega_{32} \\ \omega_{13} & \omega_{23} & -(\omega_{31} + \omega_{32}) \end{pmatrix} \cdot \mathbf{X}.$$

In a purpose of clarity, we assimilate in this notation 1, 2, and 3 to AP^+ , P^- , and AP^- , respectively. In that master equation, we introduce transition terms from state i to j of the form ω_{ij} . We link this parameter to the energy barriers that we found thanks to the energy landscape, Fig. 8(b): $\omega_{ij} = \tau_{i \rightarrow j} \exp\left(\frac{-\Delta E_{i \rightarrow j}}{k_B T}\right)$. We include here another parameter $\tau_{i \rightarrow j}$, which is the frequency attempt of switching from state i to j .

It is worth noting that the system starts from the AP^+ state, whereas in the statistical equilibrium we expect the system to be in the AP^- state. Therefore, making use of the principle of detailed balance, we obtain $\omega_{31} = \omega_{32} = 0$ (irreversible transition from to AP^- state). We can eventually solve the master equation (see Appendix for details), and access the statistical average value of the projection of the magnetization on the out of plane axis of the sample given by the magnetic moment of each state multiplied by the probability of getting each state:

$$\langle M_Z \rangle = P_{AP^+} \cdot M_{AP^+} + P_{AP^-} \cdot M_{AP^-} + P_{P^-} \cdot M_{P^-},$$

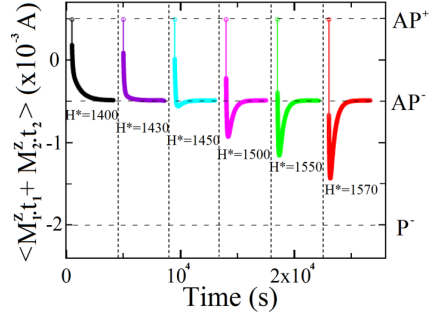


FIG. 10. Evolution of the average magnetization according to the macrospin model for various values of applied fields, with an initial configuration AP^+ .

where

$$M_{AP^+} = t_2 M_2 - t_1 M_1,$$

$$M_{AP^-} = t_1 M_1 - t_2 M_2,$$

$$M_{P^-} = -t_1 M_1 - t_2 M_2.$$

By using the material parameters shown in Table I, we obtain the magnetic relaxation curves displayed in Fig. 10.

In this calculation, we can clearly see a qualitative agreement with the experimental results shown previously [Figs. 2, 5, 6, and 7(b)]. A slight change in the field (~ 10 Oe) produces a significant change in the relaxation curve just as in the experiments. We have chosen the adequate parameters (attempt frequency τ_{ij} and surface area S) such that the simulation time matches the measurement time window (~ 120 s). Thus the attempt frequencies were $\tau_{AP^+ \rightarrow P^-} = 1.0 \times 10^{15}$, $\tau_{P^- \rightarrow AP^+} = 2.0 \times 10^{13}$, $\tau_{P^- \rightarrow AP^-} = 1.0 \times 10^4$, $\tau_{AP^+ \rightarrow AP^-} = 1.0 \times 10^{10}$, and $S \approx 20 \times 20 \text{ nm}^2$.

It is worth noting that we were forced to use a very low attempt frequency for the transition rate between state P^- and AP^- compared to the values usually estimated ($10^{10} - 10^{11}$ Hz) [30]. This could reveal that the real energy profile has a lower energy barrier than the energy barrier obtained from the macrospin model (see energy barrier on path 2 Fig. 9), which is possibly due to the domain formations at the Co layers.

We can especially notice that from a certain value of the field, the steepness of the initial magnetization decreases, stemming from a very fast $AP^+ \rightarrow P^-$ transition, and a slower $P^- \rightarrow AP^-$ relaxation. However, we can see that both transitions are much faster in the simulation than in the experiment, and we can attribute this to three factors: low attempt frequency, small surface area assumed in the model, and overestimated energy barriers. Despite the fact that we used a small surface area ($\sim 20 \times 20 \text{ nm}^2$) in the simulations, we relate this to thermally induced local switching where nucleation processes take place. Thus, if the exchange length becomes higher than 20 nm, the coherent rotation would be plausible.

The plot represented in Fig. 11 shows these probabilities calculated for the particular case of $H = -1500$ Oe. We can see that at $t = 0$ the system starts in the AP^+ state with probability 1 (solid line with circles). As the time goes by, the probability to stay in AP^+ state decreases smoothly, whereas the probability of state P^- increases faster than the state AP^- .

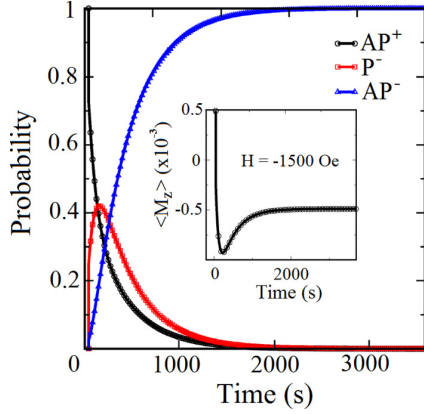


FIG. 11. Probability for the simulated elements to be in the AP^+ , P^- , or AP^- states depending on time. This simulation corresponds to an applied field of -1500 Oe. The inset represents the corresponding evolution of the projection of the magnetization on the out of plane axis with respect to time.

One of the main points made in this article is that with the assumption of a two-spin model, and neglecting the dipolar interaction, we give an explanation to the unusual nonmonotonic magnetic aftereffect that we are reporting. By doing so, we show that this phenomenon is not specific to the very stack that was studied: this property should be present for any synthetic antiferro- and ferrimagnets that display a large perpendicular magnetic anisotropy.

VI. CONCLUSION

In this paper, we report nonmonotonic magnetic relaxation curves that stem from specific energetic configurations, leading to the nucleation and growth of metastable magnetic domains. We base our observations on SQUID techniques for magnetization relaxation of the whole samples, and we perform MOKE microscopy at room temperature and MOKE measurement at low temperature in order to get accurate local information on the magnetic configuration with respect to time.

We also propose a simple theoretical model that gives an understanding of the phenomena that are observed, and corresponding numerical simulations giving a good agreement with the experimental results. We support this model with experimental evidence showing that this phenomenon can

occur provided that the temperature and field conditions can allow such energetic configurations that can be linked to the values of magnetic transitions read on the hysteresis loops. Consequently, the feature that we report here is very likely to be a common property of synthetic ferrimagnets and synthetic antiferromagnets showing perpendicular magnetic anisotropy.

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APPENDIX

From the detailed balance, we got $\omega_{31} = \omega_{32} = 0$ and hence Eq. (3) can be reduced to a simple form:

$$\frac{\partial \mathbf{X}}{\partial t} = \begin{pmatrix} -(\omega_{12} + \omega_{13}) & \omega_{21} & 0 \\ \omega_{12} & -(\omega_{21} + \omega_{23}) & 0 \\ \omega_{13} & \omega_{23} & 0 \end{pmatrix} \cdot \mathbf{X}. \quad (\text{A1})$$

Because in our model we are considering constant rates ω_{ij} , it is possible after some algebra using Laplace transform to find an explicit solution of Eq. (A1), which can be written as

$$P_{AP^+} = \frac{\omega_{21} + \omega_{23} + r_1}{r_1 - r_2} e^{r_1 t} - \frac{\omega_{21} + \omega_{23} + r_2}{r_1 - r_2} e^{r_2 t},$$

$$P_{P^-} = \frac{\omega_{12}}{r_1 - r_2} e^{r_1 t} - \frac{\omega_{12}}{r_1 - r_2} e^{r_2 t},$$

$$P_{AP^-} = 1 + \frac{\omega_{13} + r_2}{r_1 - r_2} e^{r_1 t} - \frac{r_1 + \omega_{13}}{r_1 - r_2} e^{r_2 t},$$

where

$$r_1 = -\frac{\omega_{12} + \omega_{13} + \omega_{21} + \omega_{23}}{2} - \frac{\sqrt{(\omega_{12} + \omega_{13} + \omega_{21} + \omega_{23})^2 - 4(\omega_{12}\omega_{23} + \omega_{13}\omega_{21} + \omega_{13}\omega_{23})}}{2},$$

$$r_2 = -\frac{\omega_{12} + \omega_{13} + \omega_{21} + \omega_{23}}{2} + \frac{\sqrt{(\omega_{12} + \omega_{13} + \omega_{21} + \omega_{23})^2 - 4(\omega_{12}\omega_{23} + \omega_{13}\omega_{21} + \omega_{13}\omega_{23})}}{2}.$$

We can verify that these probabilities satisfy the condition of Eq. (2). It is worth noting that in the limit $t \rightarrow \infty$ the stationary distributions are

$$P_{AP^+} = 0,$$

$$P_{P^-} = 0,$$

$$P_{AP^-} = 1.$$

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