



HAL
open science

Controlling All-Optical Helicity-Dependent Switching in Engineered Rare-Earth Free Synthetic Ferrimagnets

Jung-Wei Liao, Pierre Vallobra, Liam O'Brien, Unai Atxitia, Victor Raposo, Dorothée Petit, Tarun Vemulkar, Grégory Malinowski, Michel Hehn, Eduardo Martínez, et al.

► **To cite this version:**

Jung-Wei Liao, Pierre Vallobra, Liam O'Brien, Unai Atxitia, Victor Raposo, et al.. Controlling All-Optical Helicity-Dependent Switching in Engineered Rare-Earth Free Synthetic Ferrimagnets. *Advanced Science*, 2019, pp.1901876. 10.1002/advs.201901876 . hal-02337686

HAL Id: hal-02337686

<https://hal.univ-lorraine.fr/hal-02337686v1>

Submitted on 29 Oct 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Controlling All-Optical Helicity-Dependent Switching in Engineered Rare-Earth Free Synthetic Ferrimagnets

Jung-Wei Liao, Pierre Vallobra, Liam O'Brien, Unai Atxitia, Victor Raposo, Dorothée Petit, Tarun Vemulkar, Gregory Malinowski, Michel Hehn, Eduardo Martínez, Stéphane Mangin,* and Russell P. Cowburn

All-optical helicity-dependent switching in ferromagnetic layers has revealed an unprecedented route to manipulate magnetic configurations by circularly polarized femtosecond laser pulses. In this work, rare-earth free synthetic ferrimagnetic heterostructures made from two antiferromagnetically exchange coupled ferromagnetic layers are studied. Experimental results, supported by numerical simulations, show that the designed structures enable all-optical switching which is controlled, not only by light helicity, but also by the relative Curie temperature of each ferromagnetic layer. Indeed, through the antiferromagnetic exchange coupling, the layer with the larger Curie temperature determines the final orientation of the other layer and so the synthetic ferrimagnet. For similar Curie temperatures, helicity-independent back switching is observed and the final magnetic configuration is solely determined by the initial magnetic state. This demonstration of electrically-detected, optical control of engineered rare-earth free heterostructures opens a novel route toward practical opto-spintronics.

1. Introduction

Controlling magnetic state without the application of a magnetic field draws substantial research interest, particularly as a means to avoid localized magnetic field generation within magnetic logic and data-storage devices. One method for magnetic reversal that holds great promise is the use of circularly polarized optical laser pulses:^[1] Spurred by the first observation of all-optical switching (AOS) in rare earth (RE)-transition metal (TM) alloys, such as GdFeCo,^[2] all-optical helicity-dependent switching (AO-HDS) represents a fast and facile method to manipulate magnetization. AO-HDS has been observed in a high number of materials, and the term AOS now encompasses several possible micro-

scopic mechanisms for controlling magnetization with polarized light pulses. These range from the helicity-independent switching (AO-HIS) seen in several Gd-based samples, for example, GdCo alloys—arising from picosecond-duration angular momentum transfer between the two antiferromagnetically (AF) oriented sublattices in RE-TM alloys^[3]—through to AO-HDS in RE-free ferromagnets (FMs),^[4] such as Pt/Co/Pt multilayers,^[5] where switching occurs even in the absence of AF coupling. In this case, both laser-induced magnetic domain nucleation and domain-wall (DW) motion^[4–7] are shown to play an important role.

Despite the numerous systems where AO-HDS has been observed, in all cases any helicity dependence is intrinsic to the material composition under investigation and cannot be tuned, limiting both our ability to understand the mechanisms at play during reversal, and their ultimate use in technological applications. Synthetic ferrimagnets (SF_i), where a composite ferrimagnet is fabricated from two (or more) antiferromagnetically (AF) exchange coupled thin film FM layers, are a potential solution to this issue.^[8,9] In SF_i's, the interfacial AF coupling and the thermal response of the magnetization (i.e., nucleation of domains) within each FM layer can be tailored independently,^[10] and so, in principle, the AO-HDS mechanism. Within the wider class of composite SF_i materials, RE-free TM/FM multilayers^[8] are particularly appealing due to the low cost and earth abundance of their constituent materials. The role of magnetic anisotropy^[11] and a magnetic compensation temperature^[8] are closely scrutinized to understand AO-HDS in


Dr. J.-W. Liao, Dr. D. Petit, Dr. T. Vemulkar, Prof. R. P. Cowburn
 Cavendish Laboratory
 University of Cambridge
 JJ Thomson Avenue, Cambridge CB3 0HE, UK

Dr. P. Vallobra, Dr. G. Malinowski, Dr. M. Hehn, Prof. S. Mangin
 Institute Jean Lamour
 UMR CNRS 7198
 Université de Lorraine
 2 allée André Guinier-BP 50840, 54011 Nancy, France
 E-mail: stephane.mangin@univ-lorraine.fr

Dr. L. O'Brien
 Department of Physics
 University of Liverpool
 Liverpool L69 7ZE, UK

Dr. U. Atxitia
 Department of Physics
 Freie Universität Berlin
 Arnimalle 14, 14195 Berlin, Germany

Prof. V. Raposo, Prof. E. Martínez
 Department of Applied Physics of the Faculty of Science
 University of Salamanca
 37008 Salamanca, Spain

 The ORCID identification number(s) for the author(s) of this article can be found under <https://doi.org/10.1002/adv.201901876>.

© 2019 The Authors. Published by WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim. This is an open access article under the terms of the Creative Commons Attribution License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

DOI: 10.1002/adv.201901876

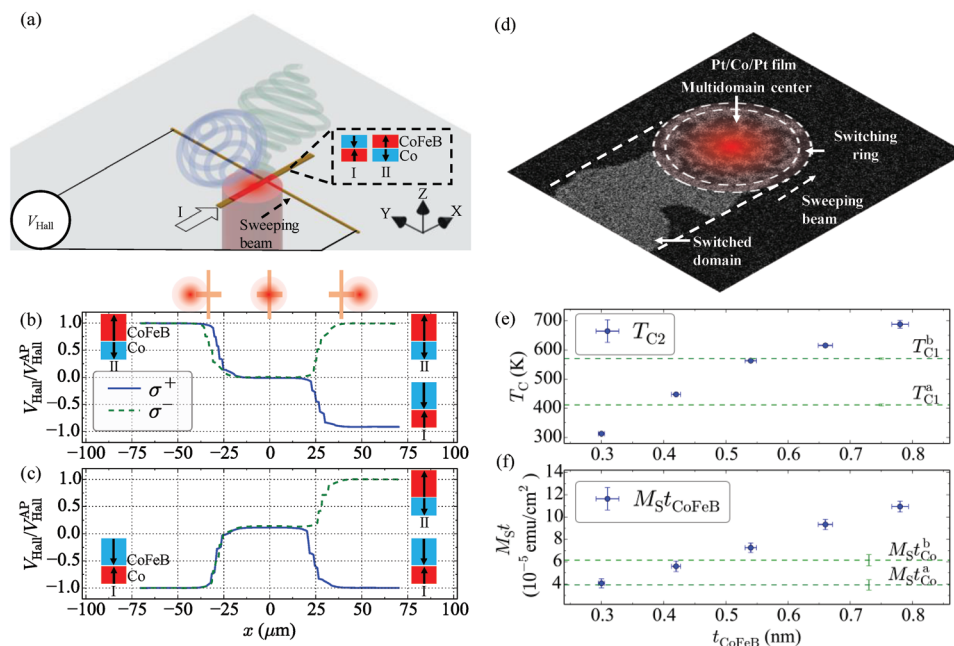


Figure 1. a) Schematics of the laser sweeping measurement, with initial State I and II shown. b,c) Measured normalized Hall voltage, $V_{\text{Hall}}/V_{\text{Hall}}^{\text{AP}}$, versus beam position, x (relative to cross center) for an SFi Hall bar with $t_{\text{Co}} = 0.73$ nm and $t_{\text{CoFeB}} = 0.54$ nm. Initial magnetic state is (b) State II and (c) State I. The measurement is repeated with two different circular polarizations, σ^+ (solid line) and σ^- (dashed line). d) Magneto-optical Kerr image of a single Pt/Co/Pt film during the laser beam sweep. e,f) Variation of Curie temperature, T_C , and the magnetization per unit area, $M_S t$, as a function of the FM2 layer thickness. The horizontal lines indicate the properties of the Co layer. T_{Cl}^a corresponds to $t_{\text{Co}}^a = 0.49$ nm, and T_{Cl}^b to $t_{\text{Co}}^b = 0.78$ nm.

this material class, while the underlying mechanism driving AO-HDS in RE-free SFi's is still debated. To this end the full control of a given composite material, for example, modulating the HD, is yet to be achieved.

In this paper, we demonstrate direct control of AO-HDS in RE-free SFi's through manipulation of the Curie temperature, T_C . AO-HDS due to irradiation by circularly polarized femto-second laser pulses is investigated using anomalous Hall effect (AHE) measurements of SFi microbars (Figure 1a).^[12] Through experiment and temperature-dependent micromagnetic simulation we determine that the reversal of the SFi is set by the distinct thermal response of the two coupled magnetic layers. A model prevails where the laser fluence heats the FM layer with the lower T_C above its T_C , while the helicity sets the magnetization direction of the other FM layer, through a mechanism such as the inverse Faraday effect (IFE) acting as a magneto-optical field on the FM layer,^[13] breaking the symmetry and so causing AO-HDS. Ultimately, by altering the thermal response of each layer, through modification of the individual layer T_C , we demonstrate a simple means to achieve all-optical control of magnetic heterostructures composed of commonly used materials, which may provide a promising route to practical opto-spintronics.^[14]

2. Results and Discussion

We perform measurements on micrometer-scale SFi Hall bars consisting of two perpendicularly magnetized, AF-coupled FM layers, termed FM1 (Co/Pt) and FM2 (CoFeB/Pt/CoFeB) and with variable thickness t_{Co} and t_{CoFeB} . Varying thickness allows us to control the T_C of each layer^[15] and therefore thermal response. Vibrating sample magnetometry (VSM) confirms two possible

AF remnant states (see Figure 1a; State I (II), FM1 points in $+z$ ($-z$)), with magnetization, M , of each FM uniformly aligned out-of-plane and antiparallel to one another (Section S1, Supporting Information). In the SFi Hall bars, AHE measurements, V_{Hall} (Section S1, Supporting Information) provide an electrical measurement sensitive to M_z of each layer, and confirm remnant states I and II are maintained in the patterned films.

Figure 1a illustrates the sweeping laser beam method used to investigate the response of the SFi to fs light pulses (discussed in the Experimental Section). The SFi is initialized in a known state (I or II) before pulsed, circularly polarized laser light (either σ^+ and σ^-), focused to an ≈ 46 μm spot, is swept along the x -axis arm. V_{Hall} is monitored as the laser passes over the Hall cross. Figure 1b (initial state I) and (c) (initial state II) illustrate the resulting V_{Hall} versus relative beam position, x , for each incident polarization, and normalized to the remnant value, $V_{\text{Hall}}^{\text{AP}}$, for $t_{\text{Co}} = 0.73$ nm and $t_{\text{CoFeB}} = 0.54$ nm. Here, one can clearly see that V_{Hall} decreases toward zero as the laser beam becomes centered on the cross. This is attributed to an initial helicity-independent demagnetization, as observed in some SFi's and single FM layers,^[12] resulting from the central portion of the laser heating FM1 and FM2 above T_C , causing loss of long-range order. With no locally favored state (I or II), upon cooling the irradiated film reorders multidomain. This demagnetized state is stable, persisting after removal of the laser beam, and so is presumed to occur only when $T > T_C$ for both individual layers, T_{Cl} and T_{C2} . As the laser spot passes the cross a helicity-dependent remagnetization occurs, that is, net magnetization is restored, with the σ^+ beam switching to final State I, while σ^- polarized beam results in State II. It is apparent from this behavior that AO-HDS occurs in these SFi's not within the beam center (which remains demagnetized), but at the trailing

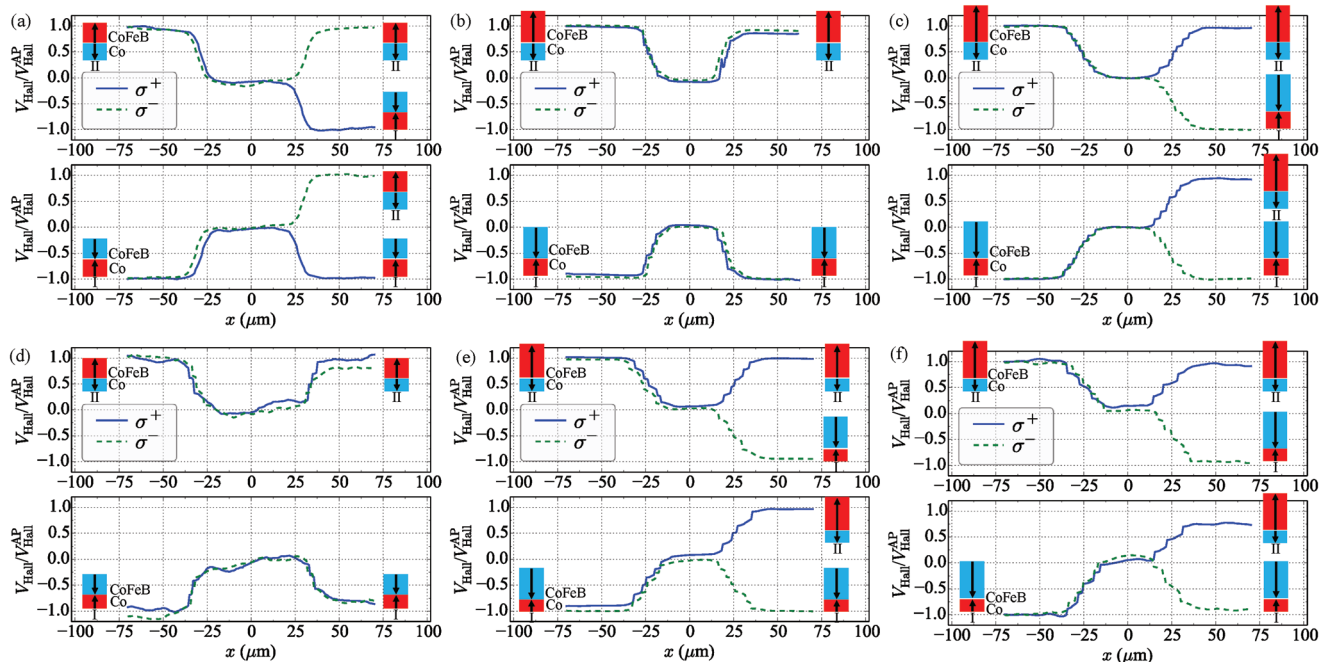


Figure 2. a–c) Variation of the normalized Hall voltage, $V_{\text{Hall}}/V_{\text{Hall}}^{\text{AP}}(x)$, for SFi with $t_{\text{Co}} = 0.73$ nm and varying t_{CoFeB} . (a) $t_{\text{CoFeB}} = 0.42$ nm, (b) 0.66 nm, and (c) 0.78 nm. Initial State I (bottom panel) or State II (top). The measurements are repeated using σ^+ (solid line) and σ^- (dashed) polarization. d–f) $V_{\text{Hall}}/V_{\text{Hall}}^{\text{AP}}(x)$, for SFi with $t_{\text{Co}} = 0.49$ nm and varying t_{CoFeB} . (d) $t_{\text{CoFeB}} = 0.42$ nm, (e) 0.66 nm, and (f) 0.78 nm.

edge of the beam profile. Changing the initial state from State II to State I, in Figure 1c we observe the same dependence of the final state on laser helicity, demonstrating both AO-HDS in RE-free SFi's and an independence of the mechanism on initial state.

On first inspection, the composite SFi appears to behave as a single FM layer (such as Pt/Co/Pt):^[12] Provided the beam fluence is sufficiently large, $T > T_C$ and thermal demagnetization occurs, with multiple domains nucleated. A ring around the central region exists where $T < T_C$ (so the film remains magnetized), but exceeds the threshold for AO-HDS.^[7] On sweeping the beam, the helicity favors one magnetization direction over the other, remagnetizing the SFi (Figure 1d). While this argument qualitatively explains data presented thus far, and follows the prevailing wisdom for single FM reversal, the situation must be more complex for SFi's where $T_{C1} \neq T_{C2}$, as is the case here. For these materials, while the central beam may still demagnetize the bar, a region exists where, for example, $T_{C1} < T < T_{C2}$, that is, FM2 remains magnetized while FM1 has lost long-range order. Should this also be sufficiently high that AO-HDS can occur, within this ring the laser fluence and so any HD effects (such as magnetic circular dichroism (MCD), or the inverse Faraday effect (IFE)) now act not on the composite SFi as a whole, but on an individual FM layer (in this case FM2). Consequently, AO-HDS is intimately linked with the precise ordering of the SFi on cooling, and tuning each T_C offers the opportunity to tailor any helicity dependence.

To test this hypothesis, we independently vary T_C in each layer. To minimize the effect of changing other parameters, we maintain the same materials for FM1 and FM2 and instead tune T_C by varying FM thickness. Figure 1e,f shows the resulting variation in (e) T_{C2} as a function of t_{CoFeB} , and (f) $M_S t_{\text{CoFeB}}$, the layer moment per unit area. Reducing t_{CoFeB} reduces T_{C2} from

bulk, due to a transition toward 2D magnetic behavior,^[15] while also linearly decreasing $M_S t_{\text{CoFeB}}$ (i.e., with minimal changes to M_S). For comparison, in Figure 1e,f T_{C1} and $M_S t_{\text{Co}}$ are shown for two different tested thicknesses of FM1, $t_{\text{Co}} = 0.49$ nm and $t_{\text{Co}} = 0.73$ nm. Examining the two plots, one can easily see for a given t_{Co} the critical FM2 thicknesses where T_{C2} exceeds T_{C1} , which in the case of $t_{\text{Co}} = 0.73$ nm occurs above $t_{\text{CoFeB}} \approx 0.6$ nm.

Investigating AO-HDS in this range, shown in Figure 2a–c for $t_{\text{Co}} = 0.73$ nm, a pattern emerges: At low t_{CoFeB} (e.g., $t_{\text{CoFeB}} < 0.66$ nm, Figure 2a and Section S2, Supporting Information), that is, while $T_{C1} > T_{C2}$, we observe a similar AO-HDS dependence as Figure 1b,c, with σ^+ light remagnetizing the SFi toward State I. Moving to the point where $T_{C1} \approx T_{C2}$ (Figure 2b, $t_{\text{CoFeB}} = 0.66$ nm) a complete loss of AO-HDS is surprisingly found, with helicity-independent reversion to the initialized state. Finally, further increasing t_{CoFeB} until $T_{C1} < T_{C2}$, for example, for $t_{\text{CoFeB}} = 0.77$ nm (Figure 2c) restores AO-HDS, however, now the helicity dependence is inverted and the σ^+ beam leads to final State II. Through tuning the T_C within each layer, we demonstrate direct control of AO-HDS.

We may compare these results against samples with reduced $t_{\text{Co}} = 0.49$ nm. Here, $T_{C1} \approx T_{C2}$ occurs at lower FM2 thickness ($t_{\text{CoFeB}} \approx 0.4$ nm) which, in theory, should be reflected in any crossover of HD. Figure 2d–f shows the resulting $V_{\text{Hall}}/V_{\text{Hall}}^{\text{AP}}(x)$ for $t_{\text{CoFeB}} = 0.42, 0.66,$ and 0.78 nm, respectively. Again, when $T_{C1} = T_{C2}$ (Figure 2d, $t_{\text{CoFeB}} \approx 0.42$ nm) helicity-independent reversion occurs, while for $T_{C2} > T_{C1}$ identical AO-HDS is found for all samples, with σ^+ leading to final State II (Figure 2e,f and Section S2, Supporting Information).

These results cannot be understood by current models of AO-HDS in single FMs, where switching is intrinsically set by the light helicity and is T_C independent.^[5] Here, we propose

a scenario to account for the observed T_C dependence: When irradiating an SFi, for example, with $T_{C1} > T_{C2}$, due to the continuous T distribution across the beam, a narrow ring exists where $T_{C1} > T > T_{C2}$ and FM1 remains ordered, but demagnetized, while FM2 loses any long-range order. AO-HDS occurs in FM1 throughout the ring, without the presence of any inter-layer exchange coupling (since $M_2 = 0$), and the layer begins to remagnetize. As the laser pulse is removed and the sample cools, the reordering of FM2 and restoration of AF coupling forces FM2 antiparallel to FM1: The SFi ordering is dictated by FM1. Should $T_{C2} > T_{C1}$, the situation becomes reversed and the helicity dependence of FM2 dictates ordering of the SFi. When $T_{C1} \approx T_{C2}$ the ring width collapses, no domain orientation direction favored (the helicity favors FM1 and FM2 moments to be parallel, while the AF coupling favors those moments to be antiparallel), and no, or at least weak, AO-HDS occurs. Simple checks of the reversal symmetry of the SFi stack with $T_{C1} > T_{C2}$, as compared with single Co FM layers ($t_{Co} = 0.73$ nm, see Section S3, Supporting Information), confirm a consistent helicity dependence between the two. Furthermore, measurements replacing the circular polarization with linearly polarized light and an externally applied field ($H_Z \approx 3$ Oe) shows an equivalent symmetry to AO-HDS (see Section S4, Supporting Information). We interpret the interaction of the circularly polarized light with M to be equivalent to an effective field in this case, potentially as a result of the IFE.^[13]

We performed micromagnetic simulations to study the magnetization dynamics induced by the femtosecond laser pulses. Here, as opposed to conventional zero-temperature simulations, a framework based on the Landau–Lifshitz–Bloch equation is used which is valid for highly nonequilibrium micromagnetism (see Section S5, Supporting Information, for details).^[13] Importantly, this model accounts for all relevant parameters of the experiment, including geometry, magnetic parameters, and thermal and magneto-optical effects due to the IFE. From simulations we extract values for $V_{Hall}/V_{Hall}^{AP}(x)$, which are in good qualitative agreement with the experimental results (see Section S6, Supporting Information). In **Figure 3a**, we show simulated snapshots of the transient magnetization in the Hall cross, with $T_{C1} > T_{C2}$. The trailing edge of the σ^- beam restores the initial magnetic state and results in no-switching (top panel of **Figure 3a**). For σ^+ , switching occurs (bottom panel of **Figure 3a**) creating a reversed domain in the illuminated area. The results further validate the proposed switching behavior explanation: The FM layer with the highest T_C controls the helicity dependence of the SFi ordering.

An interesting property of a SFi is that by appropriate engineering of layer structure one can control the so-called magnetization compensation point, that is, create a synthetic AF, where $M_1 = M_2$. This point plays an important role for single-shot switching in GdFeCo^[3] and could play a similar role in AO-HDS. To consider this, we first examine **Figure 1f**, where the critical t_{CoFeB} for SFi moment compensation is readily seen. Throughout our measurements the compensation point is systematically at higher t_{CoFeB} than the HD inversion thickness, suggesting a limited role of magnetization compensation in this system. Further numerical simulations for parameters at and around the compensation point (varying both T_C and $M_S t_{CoFeB}$ for the two FMs) also show a dominant role of T_C in determining AO-HDS in these SFi's, over any compensation point (see Section S7, Supporting Information).

We further studied the response of the SFi to light pulses with different laser fluence using the sweeping beam method (see Section S8, Supporting Information). In this case, $t_{Co} = 0.73$ nm and $t_{CoFeB} = 0.54$ nm (the switching behavior with a laser fluence of 3.9 mJ cm⁻² is shown in **Figure 1b,c**). The results reveal that the demagnetization area (where $V_{Hall} \approx 0$) reduces as the beam fluence is decreased, effectively indicating that the switching is focused to the center of the beam where the laser provides sufficient energy to create a demagnetized state. With the beam fluence reduced to 1.5 mJ cm⁻², a minimum in the demagnetization area and subsequent helicity-dependent switching is observed. Using this laser fluence, the response of the SFi to light pulses using a fixed beam method is also investigated (see Section S9, Supporting Information). The results show that as the laser pulse number increases up to 1×10^3 pulses, V_{Hall} decreases toward zero. Partial helicity-dependent remagnetization occurs when the pulse number exceeds 5×10^3 but, in general, a sweeping beam is required to ensure complete HD-AOS.

As a final point, we address the origin of the helicity-independent back switching that occurs when $T_{C1} \approx T_{C2}$. When $T_{C1} = T_{C2}$, the AO-HDS ring width collapses, resulting in a similar thermal response for each layer. Strong AF coupling remains and consequently no net expansion of one domain structure over the other should take place. Still, a heat gradient is present along the cross arms, due to the localized laser pulse, which is predicted to drive domain walls toward the hotter region (even in ideal antiferromagnets).^[16] The sweeping of the laser spot should therefore cause domain expansion, as DWs propagate along this gradient. In this case, any regions that have not otherwise been demagnetized, for example, the vertical arms of the Hall cross, are expanded, restoring the initial magnetization of the structure. In **Figure 3b**, we show the calculated transient magnetic configurations in the Hall cross, for the case of $T_{C1} = T_{C2}$. While in the horizontal wire we obtained a multidomain state, the vertical arm breaks this randomizing effect and restores the initial magnetic state of the cross position. Additional simulations without the vertical arm (see Section S10, Supporting information) show multiple domain states as the laser beam is swept over the wire, confirming that the cross geometry is essential for the backswitching effect.

In addition to explaining AO-HDS effects in SFi's, the presented results provide valuable information on the AO-HDS mechanism, in general. Indeed, a picture prevails where the laser pulse must provide sufficient heat for the FM to reach a temperature approaching T_C , at which point light helicity can induce a symmetry-breaking reversal mechanism, such as the IFE, and lead to switching.

3. Conclusion

To summarize, we studied all-optical helicity-dependent switching in planar Hall crosses made from rare-earth free synthetic ferrimagnetic heterostructures (two FM layers, AF exchange coupled). We demonstrate that the AOS not only depends on the light helicity but also on the relative Curie temperature of each ferromagnetic layer. Indeed, by varying the thickness of the individual FM layer constituting the SFi, the helicity dependence of magnetic switching can be modulated. This behavior can be explained by the interaction between the

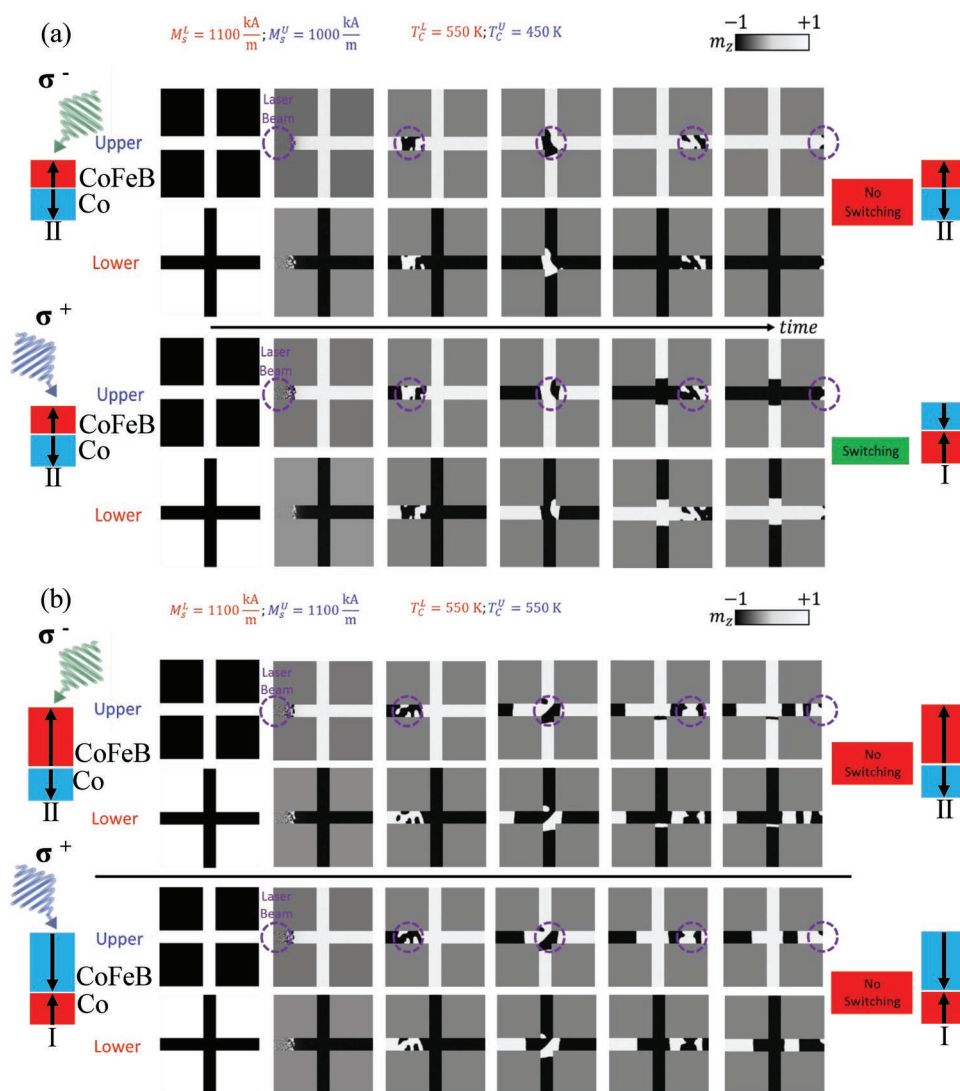


Figure 3. Modeled transient snapshots of magnetic configurations in the Hall crosses. The crosses are composed of two ferromagnetic layers with a) $M_S^{\text{Co}} < M_S^{\text{CoFeB}}$ and $T_C^{\text{Co}} < T_C^{\text{CoFeB}}$ and b) $M_S^{\text{Co}} = M_S^{\text{CoFeB}}$ and $T_C^{\text{Co}} = T_C^{\text{CoFeB}}$. Here, laser beam (position indicated by the purple circle) sweeps from left to right, with the SFi initialized in State II. The top panel corresponds to σ^- laser helicity, bottom to σ^+ . The final magnetic state (shown on the right) is the state at the cross center.

laser polarization and the net moment at switching temperatures, which is governed by the relative T_C of each layer. When the Curie temperatures of each layer are matched, we observe a helicity-independent backswitching, with the final configuration determined only by the initial state. In most materials to-date helicity dependence is intrinsic, with backswitching demonstrations limited to select rare-earth alloys,^[17] or numerical calculations.^[9] Here, we successfully demonstrate each of these switching behaviors in Hall devices, made from the same constituent materials and commonly used ferromagnetic materials. Furthermore, Co/Pt and CoFeB/Pt are magnetically high anisotropy materials commonly used in spintronic devices and magnetic recording media. As the current developed synthetic ferrimagnets are composed of these materials, we believe that the developed heterostructures with distinct all-optical switching behaviors have opened a pathway toward practical opto-spintronics.

4. Experimental Section

Sample Fabrication: Samples were grown on glass substrates at room temperature by DC magnetron sputtering, base pressure 6×10^{-8} mbar. Two sets of samples were prepared: (a) single FM layers, Ta (3 nm)/Pt (4 nm)/FM₁₍₂₎/Pt capping (4 nm), where FM₁ = Co ($t_{\text{Co}} = 0.49$ or 0.73 nm), and FM₂ = CoFeB (t_{CoFeB})/Pt (0.4 nm)/CoFeB (t_{CoFeB}) ($0.30 \leq t_{\text{CoFeB}} \leq 0.78$ nm). In FM2, both CoFeB layers are FM-coupled and act as a single FM; (b) SFi samples with structure: Ta (3 nm)/Pt (4 nm)/FM₁/Pt (0.4 nm)/Ru (0.9 nm)/Pt (0.4 nm)/FM₂/Pt capping (4 nm). The Pt/Ru/Pt stack provides the antiferromagnetic Ruderman–Kittel–Kasuya–Yosida (RKKY) interlayer exchange coupling between the adjacent FMs.^[10] All layers were grown at 8×10^{-3} mbar Ar partial pressure.

Measurements: Magnetic properties of the films were studied by vibrating sample magnetometry. The films were patterned into electrical Hall devices by optical lithography and ion milling. The Hall voltage was measured along the y-direction, while a DC current was injected along the x-direction. The width of the current carrying wire is $\approx 5 \mu\text{m}$, giving a DC current density of $\approx 6 \times 10^9 \text{ A m}^{-2}$. A Ti:sapphire fs-laser was used with wavelength 800 nm, pulse duration

43 fs, and laser repetition rate of 5 kHz. The Gaussian laser beam spot was focused with a full-width at half-maximum, FWHM $\approx 46 \mu\text{m}$. A quarter-wave plate was used to create circularly [right- (σ^+) and left-handed (σ^-)] or linearly polarized (π) beam. The laser fluence is 3.9 mJ cm^{-2} , unless otherwise stated, with a half-wave plate used to adjust beam power. For all experiments, the initial magnetic configuration (state I or II) was set by applying an external saturating magnetic field ($\pm z$ -direction), before returning to remanence. Hall voltage was constantly monitored (0.2 s per point), and unless stated otherwise, no magnetic field was applied during the measurement. For each initial state, the measurement was repeated using each circular polarization of the incoming light, σ^+ and σ^- . Switching was investigated via two methods: (1) Sweeping beam method—the laser beam was swept in discrete steps across the Hall device in the x -direction, at a step rate of $1.4 \mu\text{m s}^{-1}$; (2) Fixed beam method—beam position is fixed at the center of the cross, with the number of incident pulses controlled by a pulse picker.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

The authors acknowledge Spencer Brennan, Juan-Carlos Rojas-Sanchez, and Mohammed Salah El Hadri for assisting device fabrication and measurements. The authors thank Ulrich Nowak, Joerg Wunderlich, Amalio Fernández-Pacheco, Alex Welbourne, and Dedalo Sanz-Hernandez for useful discussions. J.-W.L. acknowledges the support from the Newton International Fellowship scheme NF150217. U.A. acknowledges financial support by the Deutsche Forschungsgemeinschaft through SFB/TRR 227 “Ultrafast Spin Dynamics,” Project A08. L.O.B. acknowledges support from the UK EPSRC (Grant No. EP/P005713/1). P.V., M.H., G.M., and S.M. were supported by the Institut Carnot ICEEL for the project “Optic-switch” and Matelas and by the French PIA project “Lorraine Université d’Excellence,” (Ref. No. ANR-151DEX-04-LUE). Experiments were performed using equipment funded by FEDER (EU), ANR, Région Grand Est, and Métropole Grand Nancy. The work by V.R. and E.M. was supported by Project No. MAT2017-87072-C4-1-P from the (Ministerio de Economía y Competitividad) Spanish government and Project No. SA299P18 from Consejería de Educación Junta de Castilla y León.

Conflict of Interest

The authors declare no conflict of interest.

Keywords

all-optical switching, ferromagnets, synthetic ferrimagnets

Received: July 21, 2019

Revised: September 25, 2019

Published online:

- [1] a) A. Kirilyuk, A. V. Kimel, T. Rasing, *Rev. Mod. Phys.* **2010**, *82*, 2731 and the references therein; b) G. P. Zhang, T. Latta, Z. Babyak, Y. H. Bai, T. F. George, *Mod. Phys. Lett. B* **2016**, *30*, 16300052.
- [2] C. D. Stanciu, F. Hansteen, A. V. Kimel, A. Kirilyuk, A. Tsukamoto, A. Itoh, T. Rasing, *Phys. Rev. Lett.* **2007**, *99*, 047601.
- [3] a) T. Ostler, J. Barker, R. Evans, R. Chantrell, U. Atxitia, O. Chubykalo-Fesenko, S. El Moussaoui, L. Le Guyader, E. Mengotti, L. Heyderman, F. Nolting, A. Tsukamoto, A. Itoh, D. Afanasiev, B. Ivanov, A. Kalashnikova, K. Vahaplar, J. Mentink, A. Kirilyuk, T. Rasing, A. Kimel, *Nat. Commun.* **2012**, *3*, 666 and the references therein; b) J. H. Mentink, J. Hellsvik, D. V. Afanasiev, B. A. Ivanov, A. Kirilyuk, A. V. Kimel, O. Eriksson, M. I. Katsnelson, T. Rasing, *Phys. Rev. Lett.* **2012**, *108*, 057202.
- [4] a) M. O. A. Ellis, E. E. Fullerton, R. W. Chantrell, *Sci. Rep.* **2016**, *6*, 30522; b) R. John, M. Berritta, D. Hinzke, C. Müller, T. Santos, H. Ulrichs, P. Nieves, J. Walowski, R. Mondal, O. Chubykalo-Fesenko, J. McCord, P. M. Oppeneer, U. Nowak, M. Münzenberg, *Sci. Rep.* **2017**, *7*, 4114.
- [5] a) C.-H. Lambert, S. Mangin, B. S. D. C. S. Varaprasad, Y. K. Takahashi, M. Hehn, M. Cinchetti, G. Malinowski, K. Hono, Y. Fainman, M. Aeschlimann, E. E. Fullerton, *Science* **2014**, *345*, 1337; b) K. C. Kuiper, T. Roth, A. J. Schellekens, O. Schmitt, B. Koopmans, M. Cinchetti, M. Aeschlimann, *Appl. Phys. Lett.* **2014**, *105*, 202402; c) R. Medapalli, D. Afanasiev, D. K. Kim, Y. Quessab, S. Manna, S. A. Montoya, A. Kirilyuk, T. Rasing, A. V. Kimel, E. E. Fullerton, *Phys. Rev. B* **2017**, *96*, 22442.
- [6] T. Janda, P. E. Roy, R. M. Otxoa, Z. Šobáň, A. Ramsay, A. C. Irvine, F. Trojaneck, M. Surýnek, R. P. Campion, B. L. Gallagher, P. Němec, T. Jungwirth, J. Wunderlich, *Nat. Commun.* **2017**, *8*, 15226.
- [7] a) G.-M. Choi, A. Schleife, D. G. Cahill, *Nat. Commun.* **2017**, *8*, 15085; b) Y. Quessab, R. Medapalli, M. S. E. Hadri, M. Hehn, G. Malinowski, E. E. Fullerton, S. Mangin, *Phys. Rev. B* **2018**, *97*, 054419.
- [8] S. Mangin, M. Gottwald, C.-H. Lambert, D. Steil, V. Uhlř, L. Pang, M. Hehn, S. Alebrand, M. Cinchetti, G. Malinowski, Y. Fainman, M. Aeschlimann, E. E. Fullerton, *Nat. Mater.* **2014**, *13*, 286.
- [9] a) R. F. L. Evans, T. A. Ostler, R. W. Chantrell, I. Radu, T. Rasing, *Appl. Phys. Lett.* **2014**, *104*, 082410; b) S. Gerlach, L. Oroszlany, D. Hinzke, S. Sievering, S. Wienholdt, L. Szunyogh, U. Nowak, *Phys. Rev. B* **2017**, *95*, 224435.
- [10] R. Lavrijsen, A. Fernández-Pacheco, D. Petit, R. Mansell, J. H. Lee, R. P. Cowburn, *Appl. Phys. Lett.* **2012**, *100*, 052411.
- [11] G. P. Zhang, Y. H. Bai, T. F. George, *J. Phys.: Condens. Matter* **2017**, *29*, 425801.
- [12] M. S. El Hadri, P. Pirro, C.-H. Lambert, S. Petit-Watelot, Y. Quessab, M. Hehn, F. Montaigne, G. Malinowski, S. Mangin, *Phys. Rev. B* **2016**, *94*, 064412.
- [13] U. Atxitia, D. Hinzke, U. Nowak, *J. Phys. D: Appl. Phys.* **2016**, *50*, 033003.
- [14] P. Němec, M. Fiebig, T. Kampfrath, A. V. Kimel, *Nat. Phys.* **2018**, *14*, 229.
- [15] C. M. Schneider, P. Bressler, P. Schuster, J. Kirschner, J. J. de Miguel, R. Miranda, *Phys. Rev. Lett.* **1990**, *64*, 1059.
- [16] S. Selzer, U. Atxitia, U. Ritzmann, D. Hinzke, U. Nowak, *Phys. Rev. Lett.* **2016**, *117*, 107201.
- [17] A. R. Khorsand, M. Savoini, A. Kirilyuk, A. V. Kimel, A. Tsukamoto, A. Itoh, T. Rasing, *Phys. Rev. Lett.* **2013**, *110*, 107205.