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CRITICAL CRITERIA OF FUJITA TYPE FOR A SYSTEM OF INHOMOGENEOUS WAVE INEQUALITIES IN EXTERIOR DOMAINS

MOHAMED JLELI, BESSEM SAMET, AND DONG YE

ABSTRACT. We consider blow-up results for a system of inhomogeneous wave inequalities in exterior domains. We will handle three type boundary conditions: Dirichlet type, Neumann type and mixed boundary conditions. We use a unified approach to show the optimal criteria of Fujita type for each case. Our study yields naturally optimal nonexistence results for the corresponding stationary wave system and equation. We provide many new results and close some open questions.

1. INTRODUCTION

This paper is concerned with the study of existence and nonexistence of global weak solutions to the system of wave inequalities

$$(1.1) \quad \square u \geq |x|^a |v|^p, \quad \square v \geq |x|^b |u|^q \quad \text{in } (0, \infty) \times \Omega^c.$$

Here $\square := \partial_{tt} - \Delta$ is the wave operator, Ω^c denotes the complement of Ω , with Ω a bounded smooth open set in \mathbb{R}^N containing the origin and $N \geq 2$. Let $p, q > 1$ and $a, b \geq -2$.

We will study (1.1) under three types of boundary conditions: the Dirichlet type condition:

$$(1.2) \quad (u(t, x), v(t, x)) \succeq (f(x), g(x)), \quad \text{on } (0, \infty) \times \partial\Omega;$$

the Neumann type condition:

$$(1.3) \quad \left(\frac{\partial u}{\partial \nu}(t, x), \frac{\partial v}{\partial \nu}(t, x) \right) \succeq (f(x), g(x)), \quad \text{on } (0, \infty) \times \partial\Omega;$$

and the mixed boundary condition:

$$(1.4) \quad \left(u(t, x), \frac{\partial v}{\partial \nu}(t, x) \right) \succeq (f(x), g(x)), \quad \text{on } (0, \infty) \times \partial\Omega,$$

where $f, g \in L^1(\partial\Omega, \mathbb{R}_+)$ are two fixed functions and ν is the outward unit normal vector on $\partial\Omega$, relative to Ω^c . By the notation \succeq , we mean the partial order on \mathbb{R}^2 , that is

$$(y_1, y_2) \succeq (z_1, z_2) \iff y_i \geq z_i, \quad i = 1, 2.$$

We write $y \succ z$, for $y, z \in \mathbb{R}^2$ if $y \succeq z$ and $y \neq z$.

The large-time behavior of solutions to the wave equation

$$(1.5) \quad \square u = |u|^p \quad \text{in } [0, \infty) \times \mathbb{R}^N$$

has been studied extensively since four decades. Inspired by the seminal work of John [7] in \mathbb{R}^3 , Strauss conjectured in [15] that for each $N \geq 2$, there exists a critical exponent $p_c(N)$ of

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Fujita type for the global existence question to (1.5) with compactly supported data, and it should be the positive root of the polynomial

$$(1.6) \quad (N-1)p^2 - (N+1)p - 2 = 0.$$

This conjecture is finally showed to be true for all dimensions $N \geq 2$ after twenty-five years of efforts, see for instance [7, 4, 3, 13, 12, 5, 17, 19] and the references therein. More precisely, let $N \geq 2$ and

$$p_c(N) = \frac{N+1 + \sqrt{N^2 + 10N - 7}}{2(N-1)},$$

then

- for any $(u, \partial_t u)|_{t=0}$ compactly supported with positive average, the solution to (1.5) blows up in a finite time if $1 < p \leq p_c(N)$;
- if $p > p_c(N)$, there are compactly supported initial conditions $(u, \partial_t u)|_{t=0} \succ (0, 0)$ such that the solution to (1.5) exists globally in time.

The wave inequality in the whole space was firstly studied by Kato [8]:

$$(1.7) \quad \square u \geq |u|^p \quad \text{in } [0, \infty) \times \mathbb{R}^N.$$

He found another critical exponent $\tilde{p}_c(N) = \frac{N+1}{N-1}$. Pohozaev & Veron [11] generalized Kato's work and pointed out the sharpness of \tilde{p}_c for (1.7). More precisely, they proved that,

- for any $N \geq 2$ and $1 < p \leq \tilde{p}_c(N)$, there is no global weak solution to (1.7), if

$$(1.8) \quad \int_{\mathbb{R}^N} \partial_t u(0, x) dx > 0;$$

- inversely, if $p > \tilde{p}_c(N)$, there are positive global solutions satisfying (1.7) and (1.8).

A natural question is to understand the wave equation or inequality on other unbounded domains of \mathbb{R}^N . The study of blow-up for wave equation on exterior domains was initiated by Zhang in [18]. Among many other things, he considered the inhomogeneous equation

$$(1.9) \quad \square u = |x|^\alpha |u|^p \quad \text{in } (0, \infty) \times \Omega^c,$$

where $N \geq 3$, $\alpha > -2$ and $\Omega \subset \mathbb{R}^N$ is a smooth bounded set. Under the Neumann boundary condition $\frac{\partial u}{\partial \nu} = f \geq 0$ on $(0, \infty) \times \partial\Omega$, Zhang showed that the critical exponent becomes now $\frac{N+\alpha}{N-2}$.

- when $1 < p < \frac{N+\alpha}{N-2}$, (1.9) has no global solution if $f \not\equiv 0$;
- when $p > \frac{N+\alpha}{N-2}$, problem (1.9) has global solutions for some $f > 0$.

However, the Dirichlet boundary condition case was left open, see Remark 1.5 of [18]. Recently the special case with $\alpha = 0$ and $\Omega = B_r$ was studied in [6]. Here and after, B_r denotes the ball centered at 0 with radius $r > 0$. Our study for (1.1) will yield an optimal answer for (1.9) under the Dirichlet boundary condition, see Corollary 1.9 below.

Here we are interested to understand the blow-up of solutions to (1.1) under various boundary conditions (1.2), (1.3) and (1.4). We will determine the critical criteria of Fujita type for (p, q) in each case, without any assumption on the initial data. As far as we know, we are not aware of such results concerning system of wave equations or inequalities. The study for (1.1) yields natural consequences for the corresponding stationary system, which seem also to be new for the Neumann type condition and the mixed boundary condition, see Corollary

1.8 below. We are confident that our ideas can be adapted for other situations, as damped wave operators, parabolic operators or higher order operators.

Before stating our results, let us mention in which sense the solutions are considered. Denote

$$Q = (0, \infty) \times \Omega^c \quad \text{and} \quad \Gamma = (0, \infty) \times \partial\Omega.$$

We introduce the test function space

$$\mathcal{D} = \left\{ \varphi \in C_{\text{cpt}}^2(Q, \mathbb{R}_+) : \varphi|_{\Gamma} = 0, \frac{\partial\varphi}{\partial\nu}|_{\Gamma} \leq 0 \right\}.$$

Here, $C_{\text{cpt}}^2(Q, \mathbb{R}_+)$ means the space of nonnegative C^2 functions compactly supported in Q . Notice that Ω^c is closed and $\Gamma \subset Q$.

Definition 1.1. A pair $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is a global weak solution to (1.1)-(1.2), if for any $\varphi \in \mathcal{D}$,

$$(1.10) \quad \int_Q |x|^a |v|^p \varphi dxdt - \int_{\Gamma} \frac{\partial\varphi}{\partial\nu} f d\sigma dt \leq \int_Q u \square \varphi dxdt$$

and

$$(1.11) \quad \int_Q |x|^b |u|^q \varphi dxdt - \int_{\Gamma} \frac{\partial\varphi}{\partial\nu} g d\sigma dt \leq \int_Q v \square \varphi dxdt.$$

For Neumann boundary problem, we consider the test function space

$$\mathcal{N} = \left\{ \varphi \in C_{\text{cpt}}^2(Q, \mathbb{R}_+) : \frac{\partial\varphi}{\partial\nu}|_{\Gamma} = 0 \right\}.$$

Definition 1.2. A pair $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is called a global weak solution to (1.1)-(1.3), if for any $\psi \in \mathcal{N}$,

$$(1.12) \quad \int_Q |x|^a |v|^p \psi dxdt + \int_{\Gamma} \psi f d\sigma dt \leq \int_Q u \square \psi dxdt$$

and

$$(1.13) \quad \int_Q |x|^b |u|^q \psi dxdt + \int_{\Gamma} \psi g d\sigma dt \leq \int_Q v \square \psi dxdt.$$

For the mixed boundary problem, the natural test function space is then $\mathcal{D} \times \mathcal{N}$.

Definition 1.3. A pair $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is a global weak solution to (1.1)-(1.4), if for any $(\varphi, \psi) \in \mathcal{D} \times \mathcal{N}$, there holds (1.10) and (1.13).

Define

$$I_f = \int_{\partial\Omega} f d\sigma, \quad \text{for any } f \in L^1(\partial\Omega).$$

Let sgn denote the standard sign function over \mathbb{R} . Our main result is the following.

Theorem 1.4. Assume that $(a, b) \succ (-2, -2)$, $f, g \in L^1(\partial\Omega)$, $(I_f, I_g) \succ (0, 0)$ and $p, q > 1$. Let either $N = 2$; or $N \geq 3$ and

$$(1.14) \quad \max \left\{ \text{sgn}(I_f) \times \frac{2p(q+1) + pb + a}{pq - 1}, \text{sgn}(I_g) \times \frac{2q(p+1) + qa + b}{pq - 1} \right\} > N.$$

Then

- (i) there exists no global weak solution to (1.1)–(1.2) if $f, g \geq 0$;
- (ii) there exists no global weak solution to (1.1)–(1.3);
- (iii) there exists no global weak solution to (1.1)–(1.4) if $p > 2$ and $f \geq 0$.

Furthermore, if $\Omega = B_r$, the sign condition for f, g can be erased in (i) and (iii).

Remark 1.5. The condition (1.14) is equivalent to

$$I_f > 0 \text{ and } \delta > N - 2; \quad \text{or} \quad I_g > 0 \text{ and } \gamma > N - 2,$$

where

$$(1.15) \quad \delta = \frac{a + 2 + p(b + 2)}{pq - 1}, \quad \gamma = \frac{b + 2 + q(a + 2)}{pq - 1}.$$

Therefore, (1.14) always holds true when $N = 2$, $p, q > 1$ and $(a, b) \succ (-2, -2)$.

In fact, the constants δ, γ come from the scaling transform of the stationary problem

$$(1.16) \quad -\Delta u = |x|^a v^p, \quad -\Delta v = |x|^b u^q.$$

Let (u, v) be a solution to the system (1.16), then for any $\lambda > 0$, $u_\lambda(x) = \lambda^\delta u(\lambda x)$, $v_\lambda(x) = \lambda^\gamma v(\lambda x)$ satisfy still (1.16).

Remark 1.6. Assume that $N \geq 3$, $p, q > 0$, $pq > 1$, and

$$(1.17) \quad 0 < \min(\delta, \gamma) \leq \max(\delta, \gamma) < N - 2.$$

Let $(u_*, v_*)(x) = (A_u |x|^{-\delta}, A_v |x|^{-\gamma})$ with $A_u, A_v > 0$ given by

$$A_u^{pq-1} = \delta(N - 2 - \delta)[\gamma(N - 2 - \gamma)]^p, \quad A_v^{pq-1} = \gamma(N - 2 - \gamma)[\delta(N - 2 - \delta)]^q.$$

We can check that (u_*, v_*) is a positive solution to (1.16) in $\mathbb{R}^N \setminus \{0\}$. If Ω is star-shaped with respect to the origin, there holds $\frac{\partial u_*}{\partial \nu}, \frac{\partial v_*}{\partial \nu} \geq 0$ on $\partial\Omega$ with respect to Ω^c . So (u_*, v_*) is a stationary solution to (1.1) and satisfies all the boundary conditions (1.2), (1.3) and (1.4) for suitable $f, g \geq 0$. This means that the condition (1.14) is optimal for the nonexistence of global solution to the wave system (1.1).

Remark 1.7. Assume that $B_{r_1} \subset \Omega$ with $r_1 > 0$. Let $a, b \leq 0$, $p, q > 0$, $pq > 1$. Similarly as above, there are suitable $A_1, A_2 > 0$ such that

$$u(t, x) = A_1(t + 1)^{-\frac{2(p+1)}{pq-1}}, \quad v(t, x) = A_2(t + 1)^{-\frac{2(q+1)}{pq-1}}$$

satisfy $\square u = r_1^a v^p$ and $\square v = r_1^b u^q$ in $\mathbb{R}_+ \times \mathbb{R}^N$. Therefore, (u, v) resolves (1.1) and satisfies all the boundary conditions (1.2), (1.3) and (1.4) with $f = g = 0$. This means the necessity of the assumption $(I_f, I_g) \succ (0, 0)$ in Theorem 1.4 when $a, b \leq 0$.

Clearly, Theorems 1.4 yields nonexistence results for the corresponding stationary problem

$$(1.18) \quad -\Delta u \geq |x|^a |v|^p, \quad -\Delta v \geq |x|^b |u|^q \quad \text{in } \Omega^c.$$

Corollary 1.8. Let $N \geq 2$, $f, g \in L^1(\partial\Omega)$ and $(a, b) \succ (-2, -2)$. Assume that $(I_f, I_g) \succ (0, 0)$ and $p, q > 1$ satisfy (1.14). Then (1.18) has no weak solution if one of the following conditions holds true:

- (i) $f, g \in L^1(\partial\Omega, \mathbb{R}_+)$, $(u, v) \succeq (f, g)$ on $\partial\Omega$;
- (ii) $(\frac{\partial u}{\partial \nu}, \frac{\partial v}{\partial \nu}) \succeq (f, g)$ on $\partial\Omega$;
- (iii) $f \in L^1(\partial\Omega, \mathbb{R}_+)$, $p > 2$ and $(u, \frac{\partial v}{\partial \nu}) \succeq (f, g)$ on $\partial\Omega$.

We rekind Corollary 1.3 in [16] for the Dirichlet boundary condition case, where $a, b > -2$ was assumed. It seems to be the first time that such nonexistence results are showed for (1.18) under the Neumann type condition or the mixed boundary condition. Similarly, the sign condition for f, g can be erased if $\Omega = B_r$.

Theorems 1.4 yields also new result for the following wave inequality in exterior domain

$$(1.19) \quad \square u \geq |x|^a |u|^p \text{ in } (0, \infty) \times \Omega^c, \quad u(t, x) \geq f(x) \text{ on } (0, \infty) \times \partial\Omega,$$

and answers an open question proposed in Remark 1.5 of [18].

Corollary 1.9. *Let $a > -2$, $f \in L^1(\partial\Omega, \mathbb{R}_+)$ and $N \geq 3$. If*

$$(1.20) \quad I_f > 0 \quad \text{and} \quad 1 < p < \frac{N+a}{N-2},$$

there is no global weak solution in $L^p_{loc}(Q)$ to (1.19). In other words, $p^ = \frac{N+a}{N-2}$ is the Fujita critical exponent for (1.19) if $N \geq 3$, $a > -2$.*

Indeed, when $b = a$, $q = p > 1$

$$N < \frac{2p(q+1) + pb + a}{pq-1} = \frac{2p+a}{p-1} \iff p < \frac{N+a}{N-2}.$$

Taking $(v, b, q, g) = (u, a, p, f)$ in (1.1)–(1.2), we deduce the above nonexistence result from part (i) of Theorem 1.4. Again the condition $f \geq 0$ is not necessary if $\Omega = B_r$. On the other hand, (1.19) admits positive solution for $a > -2$, $p > \frac{N+a}{N-2}$, $N \geq 3$ and $f > 0$ with $\|f\|_\infty$ is sufficiently small (see [18, Proposition 6.1]).

Remark 1.10. *Similarly, for the exterior Neumann inequality*

$$\square u \geq |x|^a |u|^p \text{ in } (0, \infty) \times \Omega^c, \quad \frac{\partial u}{\partial \nu}(t, x) \geq f(x) \text{ on } (0, \infty) \times \partial\Omega,$$

we rekind the critical exponent $p^ = \frac{N+a}{N-2}$ as indicated by [18, Theorem 1.4].*

Let us say some words for our approach which is based on suitable test functions and integral estimates. At first glance it looks like the method in [18, 16] or similar works for the blow-up study in exterior domains, however some key choices are completely different.

- In most previous works, we use cut-off functions with fixed scaling for the time variable t , we obtain then integral estimates on cylinder type domain $Q_D = \Sigma_t \times \Sigma_x$ where $|\Sigma_x| \sim R^N$ and Σ_t is of length CR or CR^2 . Here we consider a large scale for t by choosing $|\Sigma_t| \sim R^\theta$ with θ large enough.
- In [18, 16], they often use test functions with support away from the boundary $\partial\Omega$, hence it's more difficult to observe the effect of the Dirichlet boundary condition. In this work, we make use of harmonic function on Ω^c with zero boundary condition, which permits to cut off only at infinity.

These ideas make our method more transparent, for example we avoid the iterative step used in [18, 16].

The paper is organized as follows. In section 2, we establish some preliminary estimates that will be used in the proof of our main results. In Section 3, we prove Theorem 1.4 in two dimensional case. The proof of Theorem 1.4 for $N \geq 3$ is given in Section 4. Finally, some open questions are raised in Section 5.

The symbols C or C_i denote always generic positive constants, which are independent of the scaling parameter T and the solutions u, v . Their values could be changed from one line to another. We will write $B := B_1$ for the unit ball, and we will use the notation $h \sim k$ for two positive functions or quantities, which satisfy $C_1 h \leq k \leq C_2 h$.

2. PRELIMINARY ESTIMATES

Let $N \geq 2$. We introduce the following harmonic function in Ω^c :

$$-\Delta H_\Omega = 0 \text{ in } \Omega^c, \quad H_\Omega = 0 \text{ on } \partial\Omega;$$

and

$$\lim_{|x| \rightarrow \infty} \frac{H_\Omega(x)}{\ln|x|} = 1 \text{ if } N = 2; \quad \lim_{|x| \rightarrow \infty} H_\Omega(x) = 1 \text{ if } N \geq 3.$$

Clearly H_Ω is uniquely determined and $H_\Omega > 0$ in $\bar{\Omega}^c$.

We need also two cut-off functions. Let $\xi \in C^\infty(\mathbb{R}^N)$ satisfies

$$0 \leq \xi \leq 1; \quad \xi \equiv 1 \text{ in } B; \quad \xi(x) \equiv 0 \text{ if } |x| \geq 2.$$

Fix also $\vartheta \in C^\infty(\mathbb{R})$ such that

$$\vartheta \geq 0, \quad \vartheta \not\equiv 0, \quad \text{supp}(\vartheta) \subset (0, 1).$$

For $0 < T < \infty$, let

$$\Xi_T(x) = H_\Omega(x) \xi\left(\frac{x}{T}\right)^k \text{ in } \Omega^c$$

and

$$\vartheta_T(t) = \vartheta\left(\frac{t}{T^\theta}\right)^k \text{ in } (0, \infty).$$

Here, $k \geq 2$ and $\theta > 0$ are constants to be chosen later.

Consider

$$D_T(t, x) = \vartheta_T(t) \Xi_T(x), \quad (t, x) \in (0, \infty) \times \Omega^c$$

and

$$N_T(x) = \vartheta_T(t) \xi\left(\frac{x}{T}\right)^k, \quad (t, x) \in (0, \infty) \times \Omega^c.$$

Obviously, for any $T > \text{dist}(0, \partial\Omega)$ and $\theta > 0$,

$$(D_T, N_T) \in \mathcal{D} \times \mathcal{N}.$$

Denote $H := H_B$, i.e.

$$H(x) = \begin{cases} \ln|x| & \text{if } N = 2, \\ 1 - |x|^{2-N} & \text{if } N \geq 3. \end{cases}$$

In the following, we will give some integral estimates for D_T and N_T . Our approach uses only the asymptotic behavior of H_Ω and its derivatives at infinity. For simplicity, we will detail our proof only for the unit open ball B . The readers can be convinced easily that the same ideas work well for general smooth open sets Ω . More precisely, as $B_{r_2} \subset \Omega \subset B_{r_1}$ with $r_1 > r_2 > 0$, thanks to the maximum principle, we have $H_{B_{r_1}} \leq H_\Omega \leq H_{B_{r_2}}$ in $B_{r_1}^c$. The standard elliptic theory yields that $|\nabla^k H_\Omega|(x) \sim |\nabla^k H|(x)$ as $|x| \rightarrow \infty$, for all $k \geq 0$.

The following estimates follow from standard calculations.

Lemma 2.1. *Let $N = 2$, $\alpha \in \mathbb{R}$ and $\beta > -1$. There holds, as $T \rightarrow +\infty$,*

$$\int_{1 < |x| < T} |x|^\alpha (\ln |x|)^\beta dx \sim \begin{cases} 1 & \text{if } \alpha < -2, \\ (\ln T)^{\beta+1} & \text{if } \alpha = -2, \\ T^{\alpha+2} (\ln T)^\beta & \text{if } \alpha > -2; \end{cases}$$

and for any $\alpha, \beta \in \mathbb{R}$, we have

$$\int_{T < |x| < 2T} |x|^\alpha (\ln |x|)^\beta dx \sim T^{\alpha+2} (\ln T)^\beta, \quad \text{as } T \rightarrow +\infty.$$

Lemma 2.2. *Let $N \geq 3$, $\alpha \in \mathbb{R}$ and $\beta > -1$. There holds, as $T \rightarrow +\infty$.*

$$\int_{1 < |x| < T} |x|^\alpha (1 - |x|^{2-N})^\beta dx \sim \begin{cases} 1 & \text{if } \alpha < -N; \\ \ln T & \text{if } \alpha = -N; \\ T^{\alpha+N} & \text{if } \alpha > -N; \end{cases}$$

and for any $\alpha, \beta \in \mathbb{R}$, we have

$$\int_{T < |x| < 2T} |x|^\alpha (1 - |x|^{2-N})^\beta dx \sim T^{\alpha+N}, \quad \text{as } T \rightarrow +\infty.$$

2.1. Estimates involving D_T . By the definitions of Ξ_T and ϑ_T , there holds

$$\Delta \Xi_T(x) = \frac{2k}{T} \xi \left(\frac{x}{T} \right)^{k-1} \nabla H_\Omega(x) \nabla \xi \left(\frac{x}{T} \right) + \frac{H_\Omega(x)}{T^2} [k(k-1) \xi^{k-2} |\nabla \xi|^2 + k \xi^{k-1} \Delta \xi] \left(\frac{x}{T} \right)$$

and

$$\vartheta_T''(t) = \frac{1}{T^{2\theta}} \vartheta \left(\frac{t}{T^\theta} \right)^{k-2} [k(k-1) \vartheta'^2 + k \vartheta \vartheta''] \left(\frac{t}{T^\theta} \right).$$

We deduce then

Lemma 2.3.

$$|\Delta \Xi_T(x)| \leq C_k \left(\frac{H_\Omega(x)}{T^2} + \frac{|x|^{1-N}}{T} \right) \xi \left(\frac{x}{T} \right)^{k-2} \chi_{\{T < |x| < 2T\}} \quad \text{in } \Omega^c.$$

and

$$|\vartheta_T''(t)| \leq \frac{C_k}{T^{2\theta}} \vartheta \left(\frac{t}{T^\theta} \right)^{k-2} \chi_{\{0 < t < T^\theta\}} \quad \text{in } (0, \infty).$$

As the harmonic function H_Ω has very different behaviors in dimension two comparing to higher dimensions, we separate the study in two cases: $N = 2$ and $N \geq 3$.

Lemma 2.4. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N = 2$. We have, as $T \rightarrow +\infty$,*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\partial_{tt} D_T|^{\frac{m}{m-1}} dx dt = \begin{cases} O \left(T^{2 - \frac{\tau + (m+1)\theta}{m-1}} \ln T \right) & \text{if } \tau < 2(m-1), \\ O \left(T^{-\frac{(m+1)\theta}{m-1}} (\ln T)^2 \right) & \text{if } \tau = 2(m-1), \\ O \left(T^{-\frac{(m+1)\theta}{m-1}} \right) & \text{if } \tau > 2(m-1). \end{cases}$$

Proof. Without loss of generality, let $\Omega = B$. By the definition of D_T and Lemma 2.3, we get

$$\begin{aligned} \int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\partial_{tt} D_T|^{\frac{m}{m-1}} dx dt &= \int_0^\infty \vartheta_T(t)^{\frac{-1}{m-1}} |\vartheta_T''(t)|^{\frac{m}{m-1}} dt \times \int_{B^c} |x|^{\frac{-\tau}{m-1}} \Xi_T(x) dx \\ &\leq CT^{\frac{-2\theta m}{m-1}} \int_0^{T^\theta} \vartheta^{k - \frac{2m}{m-1}} \left(\frac{t}{T^\theta} \right) dt \times \int_{1 < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x) dx \\ &\leq CT^{-\frac{(m+1)\theta}{m-1}} \int_{1 < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x) dx. \end{aligned}$$

Using Lemma 2.1 with $\alpha = \frac{-\tau}{m-1}$ and $\beta = 1$, we obtain the claimed estimate. \square

Similarly, we deduce from Lemma 2.2 that

Lemma 2.5. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N \geq 3$. There holds, as $T \rightarrow +\infty$,*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\partial_{tt} D_T|^{\frac{m}{m-1}} dx dt = \begin{cases} O\left(T^{N - \frac{\tau + (m+1)\theta}{m-1}}\right) & \text{if } \tau < N(m-1), \\ O\left(T^{-\frac{(m+1)\theta}{m-1}} \ln T\right) & \text{if } \tau = N(m-1), \\ O\left(T^{-\frac{(m+1)\theta}{m-1}}\right) & \text{if } \tau > N(m-1). \end{cases}$$

Furthermore, there holds

Lemma 2.6. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N = 2$. Then*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta D_T|^{\frac{m}{m-1}} dx dt = O\left(T^{\theta - \frac{\tau+2}{m-1}} \ln T\right), \quad \text{as } T \rightarrow +\infty.$$

Proof. Consider still $\Omega = B$. By the definition of D_T ,

$$\begin{aligned} &\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta D_T|^{\frac{m}{m-1}} dx dt \\ (2.1) \quad &= \int_0^\infty \vartheta_T(t) dt \times \int_{B^c} |x|^{\frac{-\tau}{m-1}} H(x)^{\frac{-1}{m-1}} \xi\left(\frac{x}{T}\right)^{\frac{-k}{m-1}} |\Delta \Xi_T(x)|^{\frac{m}{m-1}} dx \\ &= CT^\theta \int_{B^c} |x|^{\frac{-\tau}{m-1}} H(x)^{\frac{-1}{m-1}} \xi\left(\frac{x}{T}\right)^{\frac{-k}{m-1}} |\Delta \Xi_T(x)|^{\frac{m}{m-1}} dx. \end{aligned}$$

Applying Lemma 2.3, there holds, for any $|x| > 1$,

$$\begin{aligned} &H(x)^{\frac{-1}{m-1}} \xi\left(\frac{x}{T}\right)^{\frac{-1}{m-1}} |\Delta \Xi_T(x)|^{\frac{m}{m-1}} \\ (2.2) \quad &\leq CH(x)^{\frac{-1}{m-1}} \xi^{k - \frac{2m}{m-1}} \left(\frac{x}{T}\right) \left(\frac{H(x)}{T^2} + \frac{|x|^{1-N}}{T}\right)^{\frac{m}{m-1}} \chi_{\{T < |x| < 2T\}} \\ &\leq C \left[T^{-\frac{2m}{m-1}} H(x) + T^{-\frac{m}{m-1}} H(x)^{\frac{-1}{m-1}} |x|^{\frac{(1-N)m}{m-1}} \right] \chi_{\{T < |x| < 2T\}}. \end{aligned}$$

Combining (2.1)–(2.2), we obtain

$$\begin{aligned} & \int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta D_T|^{\frac{m}{m-1}} dx dt \\ & \leq CT^{\theta - \frac{2m}{m-1}} \int_{T < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x) dx + CT^{\theta - \frac{m}{m-1}} \int_{T < |x| < 2T} H(x)^{\frac{-1}{m-1}} |x|^{\frac{-\tau + (1-N)m}{m-1}} dx \\ & = O\left(T^{\theta - \frac{\tau+2}{m-1}} \ln T\right), \end{aligned}$$

as T goes to $+\infty$. The last line is given by $N = 2$ and Lemma 2.1. \square

Very similarly, using the expression of H and Lemma 2.2, we have

Lemma 2.7. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N \geq 3$, then*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta D_T|^{\frac{m}{m-1}} dx dt = O\left(T^{N-2+\theta - \frac{\tau+2}{m-1}}\right), \quad \text{as } T \rightarrow +\infty.$$

2.2. Estimates involving N_T .

Lemma 2.8. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N \geq 2$. There holds, as $T \rightarrow +\infty$,*

$$\int_Q |x|^{\frac{-\tau}{m-1}} N_T^{\frac{-1}{m-1}} |\partial_{tt} N_T|^{\frac{m}{m-1}} dx dt = \begin{cases} O\left(T^{-\frac{(m+1)\theta}{m-1}} \ln T\right) & \text{if } \tau \geq N(m-1), \\ O\left(T^{N - \frac{\tau + (m+1)\theta}{m-1}}\right) & \text{if } \tau < N(m-1). \end{cases}$$

Proof. Consider $\Omega = B$. By the definition of N_T and Lemma 2.3, we get

$$\begin{aligned} \int_Q |x|^{\frac{-\tau}{m-1}} N_T^{\frac{-1}{m-1}} |\partial_{tt} N_T|^{\frac{m}{m-1}} dx dt &= \int_0^\infty \vartheta_T(t)^{\frac{-1}{m-1}} |\vartheta_T''(t)|^{\frac{m}{m-1}} dt \times \int_{B^c} |x|^{\frac{-\tau}{m-1}} \xi^k \left(\frac{x}{T}\right) dx \\ &\leq CT^{-\frac{(m+1)\theta}{m-1}} \int_{1 < |x| < 2T} |x|^{\frac{-\tau}{m-1}} dx, \end{aligned}$$

The desired estimate follows directly from Lemmas 2.1–2.2 with $\alpha = \frac{-\tau}{m-1}$ and $\beta = 0$. \square

Lemma 2.9. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N \geq 2$. Then*

$$\int_Q |x|^{\frac{-\tau}{m-1}} N_T^{\frac{-1}{m-1}} |\Delta N_T|^{\frac{m}{m-1}} dx dt = O\left(T^{N-2+\theta - \frac{\tau+2}{m-1}}\right), \quad \text{as } T \rightarrow +\infty.$$

Proof. As in Lemma 2.3, there holds

$$(2.3) \quad \left| \Delta \left[\xi^k \left(\frac{x}{T}\right) \right] \right| \leq C_k T^{-2} \xi \left(\frac{x}{T}\right)^{k-2} \chi_{\{T < |x| < 2T\}}.$$

We can claim the mentioned estimate similarly as for Lemmas 2.6. \square

2.3. Estimates involving D_T and N_T . The following are some estimates necessary to handle the mixed boundary problem (1.1)–(1.4).

Lemma 2.10. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 2$, $k > \frac{2m}{m-1}$ and $N \geq 2$. There holds, as $T \rightarrow +\infty$,*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\partial_{tt} N_T|^{\frac{m}{m-1}} dx dt = \begin{cases} O\left(T^{-\frac{(m+1)\theta}{m-1}} \ln T\right) & \text{if } \tau \geq N(m-1), \\ O\left(T^{N - \frac{\tau + (m+1)\theta}{m-1}}\right) & \text{if } \tau < N(m-1). \end{cases}$$

Proof. Without loss of generality, consider $\Omega = B$. By the definitions of D_T and N_T , thanks to Lemma 2.3, we get

$$\begin{aligned} & \int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\partial_{tt} N_T|^{\frac{m}{m-1}} dx dt \\ &= \int_0^\infty \vartheta_T(t)^{\frac{-1}{m-1}} |\vartheta_T''(t)|^{\frac{m}{m-1}} dt \int_{B^c} |x|^{\frac{-\tau}{m-1}} H(x)^{-\frac{1}{m-1}} \xi^k \left(\frac{x}{T} \right) dx \\ &\leq CT^{-\frac{(m+1)\theta}{m-1}} \int_{1 < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x)^{-\frac{1}{m-1}} dx. \end{aligned}$$

Applying Lemmas 2.1–2.2 with $\alpha = \frac{-\tau}{m-1}$ and $\beta = \frac{-1}{m-1} \in (-1, 0)$ (here $m > 2$ was used), we obtain the desired estimate. \square

Similarly, we have

Lemma 2.11. *Let $\tau \in \mathbb{R}$, $\theta > 0$, $m > 1$, $k > \frac{2m}{m-1}$ and $N \geq 2$. Then*

$$\int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta N_T|^{\frac{m}{m-1}} dx dt = O\left(T^{N-2+\theta-\frac{\tau+2}{m-1}}\right), \quad \text{as } T \rightarrow +\infty.$$

Proof. Using (2.3), there holds, for large T ,

$$\begin{aligned} & \int_Q |x|^{\frac{-\tau}{m-1}} D_T^{\frac{-1}{m-1}} |\Delta N_T|^{\frac{m}{m-1}} dx dt \\ &= \int_0^\infty \vartheta_T(t) dt \times \int_{\Omega^c} |x|^{\frac{-\tau}{m-1}} H(x)^{\frac{-1}{m-1}} \xi \left(\frac{x}{T} \right)^{\frac{-k}{m-1}} \left| \Delta \left[\xi^k \left(\frac{x}{T} \right) \right] \right|^{\frac{m}{m-1}} dx \\ &\leq CT^{\theta-\frac{2m}{m-1}} \int_{T < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x)^{\frac{-1}{m-1}} \xi \left(\frac{x}{T} \right)^{k-\frac{2m}{m-1}} dx \\ &\leq CT^{\theta-\frac{2m}{m-1}} \int_{T < |x| < 2T} |x|^{\frac{-\tau}{m-1}} H(x)^{\frac{-1}{m-1}} dx \\ &\leq CT^{\theta-\frac{2m}{m-1}} \int_{T < |x| < 2T} |x|^{\frac{-\tau}{m-1}} dx \\ &= CT^{N-2+\theta-\frac{\tau+2}{m-1}}. \end{aligned}$$

So we are done. \square

3. TWO DIMENSIONAL SITUATION

In this section, we prove successively the parts (i), (ii) and (iii) of Theorem 1.4 for $N = 2$. We will detail the proof for (i). The proofs for parts (ii) and (iii) are similar, so we proceed more quickly. Let $p, q > 1$ and fix

$$(3.1) \quad k > \max \left\{ \frac{2p}{p-1}, \frac{2q}{q-1} \right\}.$$

As mentioned above, we consider only $\Omega = B$, and we explain in Remarks 3.1–3.2 how the same ideas work for general case.

3.1. Proof of part (i). We argue by contradiction by assuming that the pair $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is a global weak solution to (1.1)–(1.2). For $T > 1$ and $\theta > 0$, taking $\varphi = D_T$ in (1.10), then

$$\int_Q |x|^a |v|^p D_T dxdt - \int_{\Gamma} \frac{\partial D_T}{\partial \nu} f d\sigma dt \leq \int_Q |u| |\square D_T| dxdt.$$

Moreover, as $\partial_\nu H$ is constant on ∂B ,

$$(3.2) \quad - \int_{\Gamma} \frac{\partial D_T}{\partial \nu} f d\sigma dt = C \int_0^\infty \vartheta \left(\frac{s}{T^\theta} \right)^k ds \times \int_{\partial B} f(x) d\sigma = C I_f T^\theta,$$

where C is a constant depending only on H and ϑ . This yields

$$(3.3) \quad \int_Q |x|^a |v|^p D_T dxdt + I_f T^\theta \leq C \int_Q |u| |\square D_T| dxdt.$$

Similarly, taking $\varphi = D_T$ in (1.11), we get

$$(3.4) \quad \int_Q |x|^b |u|^q D_T dxdt + I_g T^\theta \leq C \int_Q |v| |\square D_T| dxdt.$$

By Hölder's inequality, there holds

$$(3.5) \quad \int_Q |u| |\partial_{tt} D_T| dxdt \leq \left(\int_Q |x|^b |u|^q D_T dxdt \right)^{\frac{1}{q}} \left(\int_Q |x|^{\frac{-b}{q-1}} D_T^{\frac{-1}{q-1}} |\partial_{tt} D_T|^{\frac{q}{q-1}} dxdt \right)^{\frac{q-1}{q}}.$$

Using Lemma 2.4 with $\tau = b$ and $m = q$, we obtain

$$(3.6) \quad \int_Q |x|^{\frac{-b}{q-1}} D_T^{\frac{-1}{q-1}} |\partial_{tt} D_T|^{\frac{q}{q-1}} dxdt = \begin{cases} O \left(T^{2 - \frac{b+(q+1)\theta}{q-1}} \ln T \right) & \text{if } b < 2(q-1), \\ O \left(T^{-\frac{(q+1)\theta}{q-1}} (\ln T)^2 \right) & \text{if } b = 2(q-1), \\ O \left(T^{-\frac{(q+1)\theta}{q-1}} \right) & \text{if } b > 2(q-1), \end{cases}$$

as $T \rightarrow +\infty$.

On the other hand,

$$(3.7) \quad \int_Q |u| |\Delta D_T| dxdt \leq \left(\int_Q |x|^b |u|^q D_T dxdt \right)^{\frac{1}{q}} \left(\int_Q |x|^{\frac{-b}{q-1}} D_T^{\frac{-1}{q-1}} |\Delta D_T|^{\frac{q}{q-1}} dxdt \right)^{\frac{q-1}{q}}.$$

Applying Lemma 2.6 with $\tau = b$ and $m = q$, we have

$$(3.8) \quad \left(\int_Q |x|^{\frac{-b}{q-1}} D_T^{\frac{-1}{q-1}} |\Delta D_T|^{\frac{q}{q-1}} dxdt \right)^{\frac{q-1}{q}} = O \left(T^{\frac{\theta(q-1)-b-2}{q}} (\ln T)^{\frac{q-1}{q}} \right), \quad \text{as } T \rightarrow +\infty.$$

Combining (3.3) with (3.5)–(3.8), for T large enough, there holds

$$(3.9) \quad J_T(a, v) + I_f T^\theta \leq C [J_T(b, u)]^{\frac{1}{q}} \alpha(T),$$

where

$$J_T(a, v) = \int_Q |x|^a |v|^p D_T dxdt, \quad J_T(b, u) = \int_Q |x|^b |u|^q D_T dxdt$$

and

$$(3.10) \quad \alpha(T) = T^{\frac{\theta(q-1)-b-2}{q}} (\ln T)^{\frac{q-1}{q}} + \begin{cases} T^{-\frac{(q+1)\theta}{q}} (\ln T)^{\frac{2q-2}{q}} & \text{if } b \geq 2(q-1), \\ T^{\frac{2(q-1)-b-(q+1)\theta}{q}} (\ln T)^{\frac{q-1}{q}} & \text{if } b < 2(q-1). \end{cases}$$

Exchanging now the roles of u and v , using (3.4), we have also

$$(3.11) \quad J_T(b, u) + I_g T^\theta \leq C[J_T(a, v)]^{\frac{1}{p}} \beta(T),$$

where

$$(3.12) \quad \beta(T) = T^{\frac{\theta(p-1)-a-2}{p}} (\ln T)^{\frac{p-1}{p}} + \begin{cases} T^{\frac{-(p+1)\theta}{p}} (\ln T)^{\frac{2p-2}{p}} & \text{if } a \geq 2(p-1), \\ T^{\frac{2(p-1)-a-(p+1)\theta}{p}} (\ln T)^{\frac{p-1}{p}} & \text{if } a < 2(p-1). \end{cases}$$

Without loss of generality, we assume $I_f > 0$, as $(I_f, I_g) \succ (0, 0)$. Combining (3.9) and (3.11), there holds, for large T ,

$$J_T(a, v) + T^\theta \leq C J_T(a, v)^{\frac{1}{pq}} \beta(T)^{\frac{1}{q}} \alpha(T).$$

Using Young's inequality, we get

$$(3.13) \quad T^{-\theta} \alpha(T)^{\frac{pq}{pq-1}} \beta(T)^{\frac{p}{pq-1}} \geq C > 0, \quad \text{for large } T.$$

However, we claim that with large $\theta > 0$,

$$(3.14) \quad \lim_{T \rightarrow +\infty} T^{-\theta} \alpha(T)^{\frac{pq}{pq-1}} \beta(T)^{\frac{p}{pq-1}} = 0.$$

By (3.10) and (3.12), for $\theta > 0$ large enough, there hold

$$(3.15) \quad \alpha(T) \sim T^{\frac{\theta(q-1)-b-2}{q}} (\ln T)^{\frac{q-1}{q}}, \quad \beta(T) \sim T^{\frac{\theta(p-1)-a-2}{p}} (\ln T)^{\frac{p-1}{p}}, \quad \text{as } T \rightarrow +\infty.$$

Therefore

$$(3.16) \quad T^{-\theta} \alpha(T)^{\frac{pq}{pq-1}} \beta(T)^{\frac{p}{pq-1}} \sim T^{-\frac{(b+2)p+(a+2)}{pq-1}} \ln T, \quad \text{as } T \rightarrow +\infty,$$

hence (3.14) holds true (with large but fixed θ) since $(a, b) \succ (-2, -2)$. Obviously, (3.14) is not compatible with (3.13), which means that no global weak solution exists. This proves part (i) of Theorem 1.4 for $N = 2$.

Remark 3.1. For general smooth open sets Ω , we have no longer $\partial_\nu H_\Omega \equiv \text{constant}$ on $\partial\Omega$, hence we have no longer the equality (3.2) for all $f \in L^1(\partial\Omega)$. However, by Hopf's Lemma, $\partial_\nu H_\Omega \leq -C_\Omega < 0$ on $\partial\Omega$. If now $f \geq 0$ and $T > \text{dist}(0, \partial\Omega)$, there holds

$$-\int_\Gamma \frac{\partial D_T}{\partial \nu} f d\sigma dt \geq C_\Omega \int_0^\infty \vartheta \left(\frac{s}{T^\theta} \right)^k ds \times \int_{\partial\Omega} f(x) d\sigma \geq C I_f T^\theta.$$

where C depends only on Ω and ϑ . It's easy to see that all the arguments are still valid for $f, g \geq 0$.

3.2. Proof of part (ii). Assume that $(u, v) \in L^q_{loc}(Q) \times L^p_{loc}(Q)$ is a global weak solution to (1.1)–(1.3). Let

$$K_T(a, v) = \int_Q |x|^a |v|^p N_T dx dt, \quad K_T(b, u) = \int_Q |x|^b |u|^q N_T dx dt.$$

By Hölder's inequality,

$$(3.17) \quad \int_Q |u| |\square N_T| dx dt \leq C K_T(b, u)^{\frac{1}{q}} \left(\int_Q |x|^{\frac{-b}{q-1}} N_T^{\frac{-1}{q-1}} |\square N_T|^{\frac{q}{q-1}} dx dt \right)^{\frac{q-1}{q}}.$$

Applying Lemmas 2.8–2.9 with $\tau = b$, $m = q$ and $N = 2$, remarking that the involved estimates are exactly of the same order or better than those in Lemmas 2.4 and 2.6, we deduce that for T large,

$$(3.18) \quad \left(\int_Q |x|^{\frac{-b}{q-1}} N_T^{\frac{-1}{q-1}} |\square N_T|^{\frac{q}{q-1}} dx dt \right)^{\frac{q-1}{q}} \leq C\alpha(T)$$

where $\alpha(T)$ is given by (3.10). Similarly, there holds, for T large,

$$(3.19) \quad \int_Q |v| |\square N_T| dx dt \leq CK_T(a, v)^{\frac{1}{p}} \left(\int_Q |x|^{\frac{-a}{p-1}} N_T^{\frac{-1}{p-1}} |\square N_T|^{\frac{p}{p-1}} dx dt \right)^{\frac{p-1}{p}} \\ \leq CK_T(a, v)^{\frac{1}{p}} \beta(T),$$

where $\beta(T)$ is given by (3.12). Moreover, by the definition of N_T , for T large,

$$(3.20) \quad \int_{\Gamma} f N_T d\sigma dt = \int_{\Gamma} f \vartheta_T(t) d\sigma dt = CI_f T^\theta, \quad \int_{\Gamma} g N_T d\sigma dt = CI_g T^\theta.$$

Take $\psi = N_T$ in (1.12)–(1.13), combining with (3.17)–(3.19), we get

$$(3.21) \quad K_T(a, v) + I_f T^\theta \leq CK_T(b, u)^{\frac{1}{q}} \alpha(T), \quad K_T(b, u) + I_g T^\theta \leq CK_T(a, v)^{\frac{1}{p}} \beta(T).$$

Remark that (3.21) is just (3.9) and (3.11), if we replace K_T by J_T . Assuming without loss of generality $I_f > 0$, repeating the previous arguments for part (i), (3.13) still holds true. However, we can always choose $\theta > 0$ large to get (3.14), which makes (3.13) impossible. We reach a contradiction.

Remark 3.2. To get (3.20), we used only $\xi\left(\frac{x}{T}\right) \equiv 1$ on $\partial\Omega$ when T is large, which is true for general bounded open sets Ω .

3.3. Proof of part (iii). We use again the method of contradiction. Assume that $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is a global weak solution to (1.1)–(1.4), with now $p > 2$. We take (D_T, N_T) as a couple of test functions, and use the same notations J_T , K_T , $\alpha(T)$ and $\beta(T)$ as before.

Inserting $\varphi = D_T$ in (1.12), we obtain, for T large,

$$(3.22) \quad J_T(a, v) + I_f T^\theta \leq C[J_T(b, u)]^{\frac{1}{q}} \alpha(T) \leq C[K_T(b, u) \ln T]^{\frac{1}{q}} \alpha(T).$$

The key point here is to estimate $\|v \square N_T\|_{L^1(Q)}$ using $J_T(a, v)$. By Hölder's inequality,

$$(3.23) \quad \int_Q |v| |\square N_T| dx dt \leq J_T(a, v)^{\frac{1}{p}} \left(\int_Q |x|^{\frac{-a}{p-1}} D_T^{\frac{-1}{p-1}} |\square N_T|^{\frac{p}{p-1}} dx dt \right)^{\frac{p-1}{p}} \\ \leq CJ_T(a, v)^{\frac{1}{p}} \beta(T).$$

The last inequality follows from Lemmas 2.10–2.11 with $\tau = a$, $m = p$ and $N = 2$. Moreover, let $\psi = N_T$ in (1.13), using (3.23), there holds

$$(3.24) \quad K_T(b, u) + I_g T^\theta \leq C \int_Q |v| |\square N_T| dx dt \leq CJ_T(a, v)^{\frac{1}{p}} \beta(T).$$

- Assume first $I_f > 0$, combining (3.22) and (3.24), we deduce that

$$J_T(a, v) + T^\theta \leq CJ_T(a, v)^{\frac{1}{pq}} \alpha(T) [\beta(T) \ln T]^{\frac{1}{q}}.$$

Applying Young's inequality, there holds

$$T^{-\theta} \alpha(T)^{\frac{pq}{pq-1}} [\beta(T) \ln T]^{\frac{p}{pq-1}} \geq C > 0, \quad \text{for large } T.$$

However, fix $\theta > 0$ large, we have still (3.16), which is impossible seeing the above estimate.

- Assume now $I_g > 0$. Always using (3.22) and (3.24), there holds

$$K_T(b, u) + T^\theta \leq CK_T(b, u)^{\frac{1}{pq}} \beta(T) [\alpha(T) (\ln T)^{\frac{1}{q}}]^{\frac{1}{p}},$$

hence

$$T^{-\theta} [\alpha(T) \ln T]^{\frac{q}{pq-1}} \beta(T)^{\frac{pq}{pq-1}} \geq C > 0, \quad \text{for large } T.$$

Moreover, fixing a large θ such that (3.15) is valid, we get, as $T \rightarrow +\infty$,

$$T^{-\theta} [\alpha(T) \ln T]^{\frac{q}{pq-1}} \beta(T)^{\frac{pq}{pq-1}} \sim T^{-\frac{(b+2)+(\alpha+2)q}{pq-1}} (\ln T)^{1+\frac{q}{pq-1}}.$$

This contradicts the previous inequality.

To conclude, if $(I_f, I_g) \succ (0, 0)$ and $(a, b) \succ (-2, -2)$, there exists always a contradiction if a global weak solution exists. The proof of part (iii) is completed for $N = 2$. \square

4. PROOF OF THEOREM 1.4 FOR $N \geq 3$

Let $N \geq 3$, $p, q > 1$ and k satisfy (3.1). As above, we can consider just $\Omega = B$. The proof is very similar to the case $N = 2$.

4.1. Proof of parts (i)–(ii). Without restriction of the generality, suppose $I_f > 0$ and

$$(4.1) \quad \delta + 2 = \frac{2p(q+1) + pb + a}{pq-1} > N,$$

where δ is defined by (1.15). Assume that $(u, v) \in L_{loc}^q(Q) \times L_{loc}^p(Q)$ is a global weak solution to (1.1)–(1.2). Proceeding as above, by Lemmas 2.5 and 2.7 with $\tau = b$ and $m = q$, Hölder and Young's inequalities, we obtain again (3.13) with now

$$(4.2) \quad \alpha(T) = T^{\frac{(N-2+\theta)(q-1)-b-2}{q}} + \begin{cases} T^{\frac{-(q+1)\theta}{q}} (\ln T)^{\frac{q-1}{q}} & \text{if } b \geq N(q-1), \\ T^{\frac{N(q-1)-b-(q+1)\theta}{q}} & \text{if } b < N(q-1) \end{cases}$$

and

$$(4.3) \quad \beta(T) = T^{\frac{(N-2+\theta)(p-1)-a-2}{p}} + \begin{cases} T^{\frac{-(p+1)\theta}{p}} (\ln T)^{\frac{p-1}{p}} & \text{if } a \geq N(p-1), \\ T^{\frac{N(p-1)-a-(p+1)\theta}{p}} & \text{if } a < N(p-1). \end{cases}$$

Taking θ large enough, when $T \rightarrow +\infty$, there holds

$$(4.4) \quad \alpha(T) \sim T^{\frac{(N-2+\theta)(q-1)-b-2}{q}} \quad \text{and} \quad \beta(T) \sim T^{\frac{(N-2+\theta)(p-1)-a-2}{p}}.$$

Hence

$$T^{-\theta} \alpha(T)^{\frac{pq}{pq-1}} \beta(T)^{\frac{p}{pq-1}} \sim T^{N-2-\frac{(b+2)p+(\alpha+2)}{pq-1}} = T^{N-2-\delta}.$$

Thanks to (4.1), (3.14) follows by choosing a large θ .

The contradiction between (3.13) and (3.14) means that no global weak solution exists for (1.1)–(1.2). The nonexistence result for (1.1)–(1.3) can be derived by similar arguments, so we omit the proof.

4.2. Proof of part (iii). Let $p > 2$ and suppose that $(u, v) \in L^q_{loc}(Q) \times L^p_{loc}(Q)$ is a global weak solution to (1.1)–(1.4). For $T > 1$, using $\varphi = D_T$ in (1.12), we can claim that

$$(4.5) \quad J_T(a, v) + I_f T^\theta \leq C[J_T(b, u)]^{\frac{1}{q}} \alpha(T) \leq C[K_T(b, u)]^{\frac{1}{q}} \alpha(T).$$

Here we used $H(x) \leq 1$ as $N \geq 3$.

Proceeding as in the proof of part (iii) for $N = 2$, taking $\psi = N_T$ in (1.13), we get (3.24). Here $\alpha(T)$ and $\beta(T)$ are given by (4.2) and (4.3). Assume first $I_f > 0$ and (4.1) holds. Using (4.5) and (3.24), we have still (3.13), but also the claim (3.14) for θ large enough, which is impossible.

Assume now $I_g > 0$ and

$$\gamma + 2 = \frac{2q(p+1) + qa + b}{pq - 1} > N,$$

with γ given by (1.15). Combining (4.5) and (3.24), there holds

$$K_T(b, u) + T^\theta \leq CK_T(b, u)^{\frac{1}{pq}} \beta(T) \alpha(T)^{\frac{1}{p}},$$

hence

$$T^{-\theta} \alpha(T)^{\frac{q}{pq-1}} \beta(T)^{\frac{pq}{pq-1}} \geq C > 0, \quad \text{as } T \rightarrow \infty.$$

We can conclude if

$$(4.6) \quad \lim_{T \rightarrow +\infty} T^{-\theta} \alpha(T)^{\frac{q}{pq-1}} \beta(T)^{\frac{pq}{pq-1}} = 0.$$

Taking θ large enough, by (4.4), there holds $T^{-\theta} \alpha(T)^{\frac{q}{pq-1}} \beta(T)^{\frac{pq}{pq-1}} \sim T^{N-2-\gamma}$ for T large, so (4.6) holds true and the proof of part (iii) is completed. \square

5. FURTHER REMARKS

It's worthy to mention that the system of wave equations in the whole space, i.e.

$$\square u = |v|^p, \quad \square v = |u|^q \quad \text{in } (0, \infty) \times \mathbb{R}^N, \quad p, q > 1, \quad N \geq 2$$

has been extensively studied since the seminal work [1]. It is showed that for compactly supported initial data with positive averages for $\partial_t u(0, x)$, $\partial_t v(0, x)$, there exists a critical curve for the global existence, which is

$$\max \left\{ \frac{p+2+q^{-1}}{pq-1}, \frac{q+2+p^{-1}}{pq-1} \right\} = \frac{N-1}{2}.$$

The corresponding system of inequalities was studied in [11], where Theorem 6 (see also Application 2) proves the nonexistence of nontrivial global solution if

$$1 < p, q < \frac{N+1}{N-1}, \quad \int_{\mathbb{R}^N} \partial_t u(0, x) dx \geq 0 \quad \int_{\mathbb{R}^N} \partial_t v(0, x) dx \geq 0.$$

We can see that the critical criteria in the above cases are quite different for our situation. This phenomenon is similar to comparing Strauss's critical exponent $p_c(N)$ for (1.5), Kato's exponent $\tilde{p}_c(N)$ for (1.7) and Zhang's exponent p^* for (1.9). In other words, the blow-up for inequalities on exterior domains is of very different nature comparing to the whole space situation.

The critical case $N \geq 3$,

$$\max \left\{ \operatorname{sgn}(I_f) \times \frac{2p(q+1) + pb + a}{pq-1}, \operatorname{sgn}(I_g) \times \frac{2q(p+1) + qa + b}{pq-1} \right\} = N$$

for the system (1.1) is not investigated here. It should be interesting to decide whether this critical curve in (p, q) -plan belongs to the blow-up situation.

For the mixed boundary condition case (1.4), we supposed that $p > 2$ due to technical reason. It should be interesting to consider the case $1 < p \leq 2$.

As indicated in Remark 1.7, the case of wave inequalities, under homogeneous constraints, i.e. $f = g = 0$, is very special. We may have no critical criteria of Fujita type in general. However, the simple example there only works for $a, b \leq 0$. It could be interesting to understand the long term behavior of solutions to (1.1) with $a, b > 0$ and various type of homogeneous constraints with $f = g = 0$.

In the case of homogeneous constraints, another way to avoid the simple example in Remark 1.7 is to add sign condition or nonnegative average constraint on $\partial_t u(0, x)$, $\partial_t v(0, x)$ as in [8, 11]. For example, consider the following problem:

$$\square u \geq |x|^a |u|^p \text{ in } \mathbb{R}_+ \times B_r^c, \quad u \geq 0 \text{ on } \mathbb{R}_+ \times \partial B_r \quad \text{and} \quad \partial_t u(0, x) \geq 0,$$

where $B_r \subset \mathbb{R}^N$, $N \geq 3$, $a > -2$. Laptev [10] showed that the critical exponent for existence of non trivial global solution is $\frac{N+1+a}{N-1}$.

The understanding for wave equation on exterior domains with homogeneous Dirichlet boundary condition is more difficult. Consider (1.9) with $N \geq 2$, $a = 0$ and $u = 0$ on $\partial\Omega$. There are many works who suggest that the critical exponent of Fujita type could be the same as for the whole space, i.e. $p_c(N)$ given by (1.6).

- Let $1 < p \leq p_c(N)$, it's showed that for special choice of $(u_0, u_1) \succ (0, 0)$, the solution to (1.9) with $(u, \partial_t u)|_{t=0} = (\varepsilon u_0, \varepsilon u_1)$ will blow up for any $\varepsilon > 0$, see [9] and the references therein. However, the blow-up result for general (u_0, u_1) seems unknown.
- For $p > p_c(N)$, there exist global existence results for some $p > p_c(N)$ in low dimensions $N \leq 4$ with non trapping obstacle Ω and suitable $u_0, u_1 > 0$. See for instance [2, 14].

As far as we know, it seems that there is no general result for the global existence of wave equation on exterior domains (1.9) with homogeneous Neumann boundary condition.

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REFERENCES

- [1] D. Del Santo, V. Georgiev, E. Mitidieri, Global existence of the solutions and formation of singularities for a class of hyperbolic systems. Geometrical optics and related topics, Progr. Nonlinear Diff. Equ. Appl., 32, Birkhäuser (1997), 117–140.
- [2] Y. Du, J. Metcalfe, C.D. Sogge, Y. Zhou, Concerning the Strauss conjecture and almost global existence for nonlinear Dirichlet-wave equations in 4-dimensions, Comm. Partial Diff. Equ. 33 (2008), 1487–1506.
- [3] R.T. Glassey, Finite-time blow up for solutions of nonlinear wave equations, Math. Z. 177 (1981), 323–340.

- [4] R.T. Glassey, Existence in the large for $\square u = F(u)$ in two space dimensions, *Math. Z.* 178 (1981), 233–261.
- [5] V. Georgiev, H. Lindblad, C.D. Sogge, Weighted Strichartz estimates and global existence for semilinear wave equations, *Amer. J. Math.* 119 (1997), 1291–1319.
- [6] M. Jleli, B. Samet, New blow-up results for nonlinear boundary value problems in exterior domains, *Nonlinear Anal.* 178 (2019), 348–365.
- [7] F. John, Blow-up of solutions of nonlinear wave equations in three space dimensions, *Manuscripta Math.* 28 (1979), 235–268.
- [8] T. Kato, Blow-up of solutions of some nonlinear hyperbolic equations, *Comm. Pure Appl. Math.* 33 (1980), 501–505.
- [9] N. Lai, Y. Zhou, Blow up for initial boundary value problem of critical semilinear wave equation in two space dimensions, *Commun. Pure Appl. Anal.* 17 (2018), 1499–1510.
- [10] G. Laptev, Nonexistence results for higher order evolution partial differential inequalities, *Proc. Amer. Math. Soc.* 131 (2002), 415–423.
- [11] S. Pohozaev, L. Veron, Blow-up results for nonlinear hyperbolic inequalities, *Ann. Scuola Norm. Sup. Pisa Cl. Sci.* 29 (2000), 393–420.
- [12] J. Schaeffer, The equation $\square u = |u|^p$ for the critical value of p , *Proc. Roy. Soc. Edinburgh.* 101A (1985), 31–44.
- [13] T.C. Sideris, Nonexistence of global solutions to semilinear wave equations in high dimensions, *J. Diff. Equ.* 52 (1984), 378–406.
- [14] H.F. Smith, C.D. Sogge, C. Wang, Strichartz estimates for Dirichlet-wave equations in two dimensions with applications, *Trans. Amer. Math. Soc.* 364 (2012), 3329–3347.
- [15] W.A. Strauss, Nonlinear scattering theory at low energy, *J. Funct. Anal.* 41 (1981), 110–133.
- [16] Y. Sun, Nonexistence results for systems of elliptic and parabolic differential inequalities in exterior domains of \mathbb{R}^n , *Pacific J. Math.* 293 (2018), 245–256.
- [17] B. Yordanov, Q. Zhang, Finite time blow up for critical wave equations in high dimensions, *J. Funct. Anal.* 231 (2006), 361–374.
- [18] Q. Zhang, A general blow-up result on nonlinear boundary-value problems on exterior domains, *Proc. Roy. Soc. Edinburgh* 131A (2001), 451–475.
- [19] Y. Zhou, Blow-up of solutions to semilinear wave equations with critical exponent in high dimensions, *Chin. Ann. Math.* 28B (2007), 205–212.

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