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# Generation of spin waves via spin-phonon interaction in a buried dielectric thin film

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In this paper, we investigate the magnetic, optical, and lattice responses of a Pt/Cu/Bi<sub>1</sub>Y<sub>2</sub>Fe<sub>5</sub>O<sub>12</sub>/Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub> heterostructure to femtosecond laser excitation of the opaque Pt/Cu metallic bilayer. The electronic excitation generates coherent and incoherent phonons, which trigger high-frequency standing spin waves (SWs) in the dielectric Bi<sub>1</sub>Y<sub>2</sub>Fe<sub>5</sub>O<sub>12</sub> layer via a phonon-induced change of magnetic anisotropy. We find that the incoherent phonons (heat) can induce a fast (<1ps) and slow (>1000ps) decrease of the magnetic order by different spin-phonon interaction scenarios. These results open perspectives for generating high-frequency SWs in buried magnetic garnets.

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## I. INTRODUCTION

Controlling the spin dynamics at the fastest speed and with the highest efficiency has always been one of the ultimate goals of modern magnetism, as it constitutes a fundamentally important challenge with a potential high impact on data storage and processing technologies. Since the discovery of the subpicosecond demagnetization in nickel [1], laser-induced ultrafast magnetization dynamics have attracted a lot of attention and have become the rapidly developing field of femtomagnetism [2,3]. Indeed, over the past 25 years, it was shown that femtosecond laser pulses can trigger many important magnetic phenomena such as all-optical magnetization switching [4,5], ultrafast magnetic phase transition [6–8], and the generation of collective spin oscillations [9,10]. These results have provided important insight about fundamental mechanisms behind the excitation of spins with light. It was demonstrated that thermal effects play an important role in the majority of ultrafast magnetic processes induced in conducting [1,11,12] and semiconducting [13–16] magnets, while in magnetic insulators, the ultrafast control of spin dynamics can occur via nonthermal mechanisms like the inverse Faraday effect [10,17,18], the Cotton-Mouton effect [19], and the photo-induced magnetic anisotropy [20–23].

These results have motivated researchers to explore alternative possibilities to control the magnetization using other ultrashort stimuli like picosecond strain pulses and hot-electron pulses. For instance, the propagation of an ultrashort hot-electron pulse in conducting magnets can induce an ultrafast demagnetization [24–26] and a full magnetization reversal [25,26] via thermal effects. Picosecond acoustic pulses can trigger a homogeneous collective spin oscillation

( $k = 0$ , i.e., ferromagnetic resonance mode) via inverse magnetostriction in conducting [27–29], semiconducting [30], and insulating [31] magnetic materials. Very recently, the geometry for exciting spin dynamics indirectly via pumping a Pt/Cu bilayer that protects a magnetic layer from optical excitation [24,25] was investigated by ultrafast x-ray diffraction (UXRD) [32]. These experiments quantify the arrival of hot electrons, strain waves, and vibrational heat in a magnetic Ni layer.

The collective spin oscillations, known also as spin waves (SWs) or magnons, represent one of the most promising approaches toward the development of more energy-efficient and faster data transport and processing technology than the current charge-based semiconductor-based technology [33]. This is due to their very useful physical properties at room temperature, such as the low-energy dissipation, the high-frequency spectrum from GHz to THz, and the tunable wavelength down to nanoscales [34]. A key requirement for high-speed applications is to generate SWs that combine low damping with high frequency. Nanoscale thin films of yttrium iron garnet (Y<sub>3</sub>Fe<sub>5</sub>O<sub>12</sub>, YIG) with its intrinsic low magnetic damping constant [35] are particularly interesting for these applications. The frequency of exchange standing SW (SSW) modes increases quadratically with the mode number  $n$  and with decreasing thickness  $d$  of the nanolayer. A moderate exchange stiffness  $D_{\text{ex}}$  can thus support the high frequency  $\sim D_{\text{ex}}(\pi n/d)^2$  required by applications [36,37].

Generating SSWs requires nonuniform excitation across the film thickness [38], which is very challenging to induce using conventional microwave antennae [38–40]. On the other hand, for around 15 years, the excitation of SSWs with femtosecond laser pulses has been limited to conducting [9,41,42] and semiconducting [13,15,16] magnets. In our recent papers, we demonstrated that a femtosecond laser pulse can excite SSWs in insulating magnetic films of bismuth-substituted

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YIG with high in-plane magnetic anisotropy [43,44]. In this context, an important fundamental study is to investigate the possibility to take advantage of femtosecond laser pulses to trigger high-frequency SSWs in bismuth-substituted YIG films buried below a thick metallic structure. This configuration is of utmost importance since it allows exploring the effects of picosecond strain excitation and heat transfer across a normal metal-magnetic insulator interface on the magnetic properties.

Here, we report an experimental investigation of the ultrafast magnetization dynamics in a dielectric film of  $\text{Bi}_1\text{Y}_2\text{Fe}_5\text{O}_{12}$  (Bi-YIG) with high in-plane magnetic anisotropy and buried below a thick Cu/Pt metallic bilayer. We demonstrate that the indirect excitation of the Bi-YIG layer via exclusive absorption of light in the metal triggers high-frequency SSWs in the insulating Bi-YIG. Two mechanisms mediate this indirect excitation of SSWs: (i) Inverse magnetostriction changes the magnetic anisotropy proportional to lattice constant changes in the form of propagating strain waves. (ii) The lattice-heating additionally changes the magnetic anisotropy via the temperature dependence of the anisotropy constants. The strain pulse (coherent phonon) launched in the metal, traversing the Bi-YIG layer within 50 ps, and the rapid propagation of heat (incoherent phonons) can both excite SSWs. We suggest that phonons created at the Cu/Bi-YIG interface via hot electrons reflected at the metal-insulator interface additionally induce a fast ( $<1$  ps) change of the magnetic order. The slower timescale ( $>1$  ns) on which the magnetic order is further reduced can be attributed to a combination of slow phonon-spin coupling and heat transport into the Bi-YIG layer. Our model is based on the experimental results from time-resolved magneto-optical Kerr effect (TR-MOKE) and UXRD.

## II. SAMPLE CHARACTERIZATION AND EXPERIMENTAL METHODS

The sample studied in this paper was based on a 135-nm-thick  $\text{Bi}_1\text{Y}_2\text{Fe}_5\text{O}_{12}$  single crystalline film grown by pulsed laser deposition on a gadolinium gallium garnet (GGG) (100) substrate [45]. To address the effect of a phononic excitation on ultrafast magnetization dynamics that excludes optical excitation, we deposited a thick Pt (5 nm)/Cu (100 nm) non-magnetic metallic bilayer on top of the Bi-YIG film using direct current magnetron sputtering [Fig. 2(b)]. The top Pt layer was added due to its high absorption at 800 nm compared with Cu, which also ensured an efficient generation of a hot-electron pulse [24,25,32]. Experimentally, the transmission of the pump light was less than  $10^{-3}$ . The Cu was chosen for its high hot-electron lifetime [24,25,46]. Thus, the 100-nm-thick Cu layer protected the Bi-YIG film from the direct light excitation, while the hot-electron pulses propagated until reaching the back side of the Cu layer [24,25,32].

Figure 1 shows the UXRD study of the heterostructure, performed with femtosecond x-ray diffraction using a tabletop laser-driven plasma source [47]. In full analogy to the recent publication on a Pt/Cu/Ni heterostructure [32], we were able to model the time-resolved strain in Pt, Cu, and Bi-YIG with a two-temperature model in combination with a one-dimensional masses-and-springs model, using the

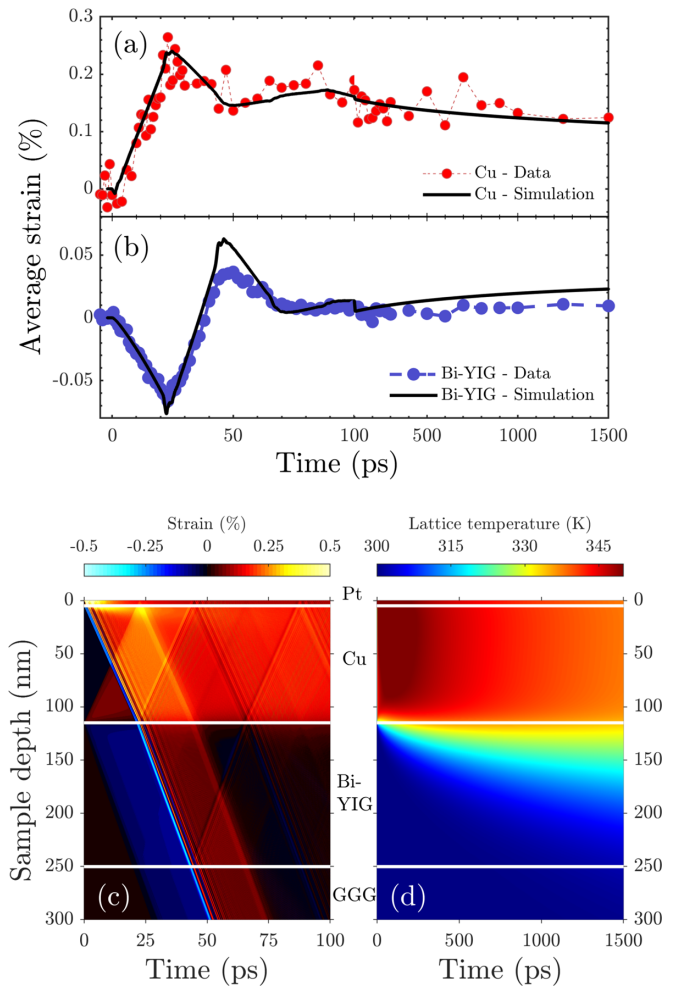


FIG. 1. Laser-induced lattice dynamics in the Pt/Cu/ $\text{Bi}_1\text{Y}_2\text{Fe}_5\text{O}_{12}$ /Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub> heterostructure. Average strain in (a) the Cu layer and (b) Bi-YIG layer as measured by ultrafast x-ray diffraction (UXRD). Calculated spatiotemporal maps of (c) the strain and (d) temperature in the heterostructure. The solid lines in (a) and (b) are results of our modeling. The UXRD measurements were performed at room temperature with a pump wavelength of 800 nm, a pulse duration of 50 fs, and an incident pump energy density of  $5 \text{ mJ/cm}^2$ .

software package *udkm1Dsim-toolbox* [48]. The two-temperature model includes an enhanced electronic heat conductivity  $\kappa_e = \kappa_0 \times T_e/T_{ph}$  compared with the equilibrium value  $\kappa_0$ . This enhancement is due to the reduction of the electron-phonon scattering channel if the lattice of Cu is cold compared with its electronic system [32]. The main difference to the previous paper is that the magnetic layer now is the insulating Bi-YIG, which cannot accept the thermal energy in the form of hot electrons. Instead, the hot electrons in Cu can excite optical phonons in Bi-YIG when they scatter at the interface, which quickly decay. In addition, heat can be transported into the Bi-YIG film via incoherent acoustic phonons from the Cu. Figures 1(a) and 1(b) show the measured time-resolved average strains  $\bar{\eta}_{\text{Cu}}(t)$  of the Cu layer and  $\bar{\eta}_{\text{Bi-YIG}}(t)$  of the Bi-YIG layer. As a first approximation, we disregard the direct transfer of energy from Cu electrons to Bi-YIG phonons and calculate the spatiotemporal lattice strain  $\eta(z, t)$  from the

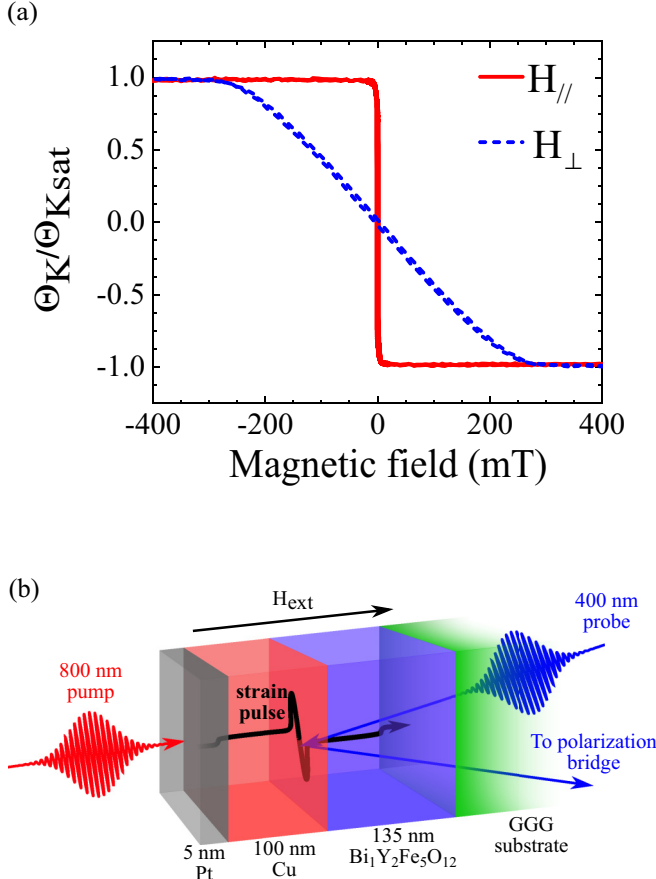


FIG. 2. Static room temperature magnetic properties of the garnet sample and the experimental configuration. (a) Normalized hysteresis loop measured in polar (dashed line) and longitudinal (solid line) configurations. (b) Sketch of the pump-probe time-resolved magneto-optical configuration that allows studying the ultrafast magnetization dynamics induced by coherent and incoherent phonons. The sketch also shows a snapshot of the bipolar strain pulse that traverses the Bi-YIG layer as modeled in Fig. 1(c). It contains a bipolar pulse with a broad, low-amplitude contribution from Cu excitation and a sharp high-amplitude contribution from Pt excitation.

temperature  $T(z, t)$  in the heterostructure [see Figs. 1(c) and 1(d)]. In our simulation, we used the same parameters for Pt and Cu in Ref. [32] except that the Cu thickness was 110 nm to match the timings of the coherent strain signals. For Bi-YIG, we used a longitudinal sound velocity  $V_{\text{Bi-YIG}}^L = 6300$  m/s [49,50], a phonon heat capacity  $C_{\text{Bi-YIG}}^{ph} = 2.9$  J/m<sup>3</sup>K [51], a phonon heat conductivity  $\kappa_{\text{Bi-YIG}}^{ph} = 6$  W/mK [51], and a Grüneisen constant  $\Gamma_{\text{Bi-YIG}}^s = 1$  [50]. The modeling of the average lattice strain is plotted as solid lines in Figs. 1(a) and 1(b), showing a reasonable agreement with the experimental data. We can use the transient temperature and strain which was calibrated to the UXRD experiment to discuss the excitation channels of the SWs. In Fig. 2(b), we display the calculated shape of the strain pulse that traverses the Bi-YIG layer, i.e., the time-dependent local strain that modulates each unit cell. It is a vertical cut through Fig. 1(c) at 25 ps. It contains a low-amplitude feature originating from the expan-

sion of the Cu layer [Fig. 1(a)], which induces the immediate decrease of the strain in Fig. 1(b), starting within less than 1 ps. At this time, the heat energy transported by the hot electrons through Cu arrived at the Bi-YIG interface. The sharper high-amplitude feature originated from the expansion of the thin Pt transducer. Both features of the strain pulse together with the thermal expansion are shown in the strain map of Fig. 1(c), which highlights the spatiotemporal dynamics of the lattice strain in the heterostructure. The heat diffusion is illustrated by the phonon temperature profile [Fig. 1(d)]. This phonon heat energy density was the main component of the stress that drove the strain [32,48,50]. More importantly, Fig. 1(d) shows that the heating of the Bi-YIG layer is strongly inhomogeneous. Although the diffusion reached the back end of the Bi-YIG layer only at  $\sim 2$  ns, the temperature rose by more than 40 K within 100 ps in the first 10 nm behind the Cu interface [see Fig. 1(d)].

The static magnetic properties of the sample are characterized using the MOKE. The normalized polar and longitudinal Kerr hysteresis loops are shown in Fig. 2(a). The out-of-plane normalized remanence ( $M_r/M_s$ ) was about 0.02, which is very small compared with the in-plane one of 0.85. The out-of-plane saturation field ( $\sim 260$  mT) was about 50 times larger than the in-plane one ( $\sim 5$  mT). These results reveal a strong in-plane magnetic anisotropy. The origin of this in-plane easy axis of magnetization is a large growth-induced anisotropy field [52,53], which has the value of  $-83$  mT in the case of our film. The TR-MOKE and reflectivity measurements were investigated with the all-optical pump-probe configuration sketched in Fig. 2(b). We have employed a femtosecond laser pulse issued from an amplified Ti-Sapphire laser system operating at a 5-kHz repetition rate and delivering 35-fs pulses at 800 nm to generate the pump and the probe beams. The pump beam was kept at the fundamental wavelength of the amplifier at 800 nm and excited the sample at normal incidence from the Pt side. The probe beam was frequency doubled to 400 nm using a beta-barium borate (BBO) crystal and incident with a small angle of  $6^\circ$  onto the GGG substrate. Both beams were linearly polarized and focused onto the sample in spot diameters of  $\sim 260 \mu\text{m}$  for the pump and  $\sim 60 \mu\text{m}$  for the probe. The probe wavelength was well above the optical absorption edge of the GGG, which allowed the probe pulses to penetrate the substrate and reach the Bi-YIG layer. After interacting with the Bi-YIG, the reflected probe pulses allowed measuring the differential changes of the magneto-optical polar Kerr rotation  $\Delta\Theta_K(t)$  and reflectivity  $\Delta R(t)$  induced by the excitation pulse as a function of the time delay  $t$  between the pump and probe pulses using a lock-in detection scheme. The external magnetic field  $H_{\text{ext}}$  was applied perpendicular to the plane of the film. All measurements were performed at room temperature.

### III. RESULTS AND DISCUSSION

Figure 3(a) shows the TR-MOKE induced by an incident pump energy density of  $E_{\text{pump}} = 9.4$  mJcm<sup>-2</sup> for  $H_{\text{ext}} = 164$  mT. Three important features can be clearly distinguished. First, the  $\Delta\Theta_K$  signal shows a small ultrafast decrease within the first picosecond after the arrival of hot electrons at the Cu/Bi-YIG interface, followed by a recovery toward the initial state in the next few picoseconds, as



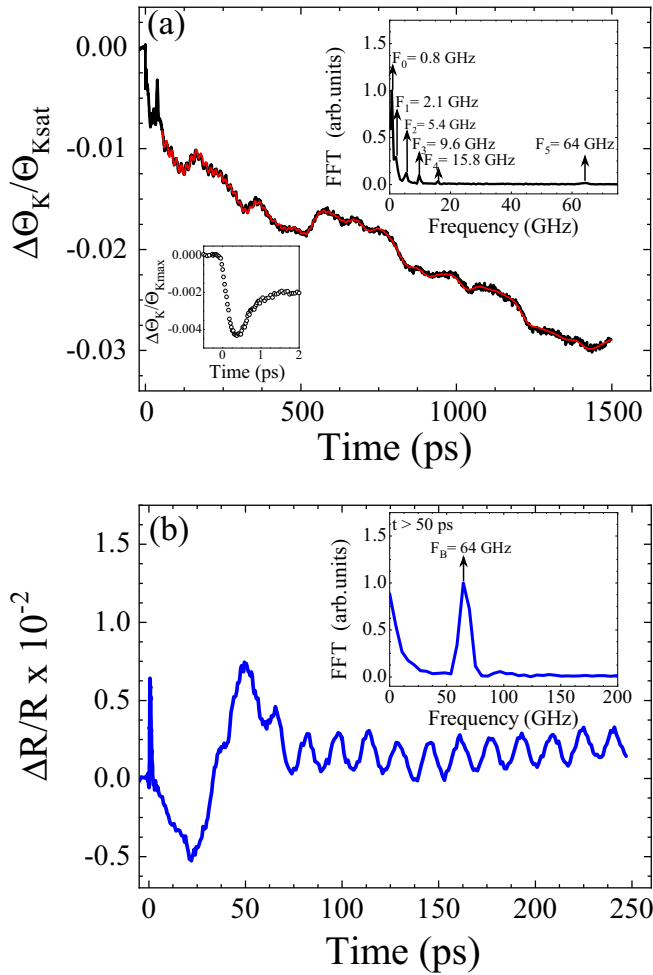


FIG. 3. Dynamics of spin and reflectivity in the  $\text{Bi}_1\text{Y}_2\text{Fe}_5\text{O}_{12}/\text{GGG}(100)$  buried below thick Pt/Cu bilayers.  $\Delta\Theta_K/\Theta_K$  and  $\Delta R/R$  induced by a pump energy density of  $E_{\text{pump}} = 9.4 \text{ mJcm}^{-2}$  for  $H_{\text{ext}} = 164 \text{ mT}$ . The solid red line in (a) is the fit using Eq. (1). Insets: (a) Fourier transform spectrum of the  $\Delta\Theta_K/\Theta_K$  data for the time delay  $t \geq 50 \text{ ps}$  (top) and a magnified view of the  $\Delta\Theta_K/\Theta_K$  signal (bottom) around the zero time delay. (b) Fourier transform spectrum of the  $\Delta R/R$  data for the time delay  $t \geq 50 \text{ ps}$  (top).

highlighted in the bottom inset of Fig. 3(a). Second, a gradual slow decrease of the  $\Delta\Theta_K$  signal started at  $\sim 5 \text{ ps}$  and reached a large value at  $1500 \text{ ps}$ . Third, very complex oscillations are superimposed on the slow exponential change of the  $\Delta\Theta_K$  signal. The characteristic times of the fast and slow decrease of the TR-MOKE are  $\tau_{\text{fast}} = 0.4 \pm 0.09 \text{ ps}$  and  $\tau_{\text{slow}} = 1220 \pm 60 \text{ ps}$ . On the other hand, the complex oscillations are formed by six modes with different amplitudes in the frequency range between  $0.8$  and  $64 \text{ GHz}$ , as seen in the fast Fourier transform (FFT) spectrum in the top inset of Fig. 3(a). To study these modes in more detail, the TR-MOKE measurements are fitted with the following time damped oscillators:

$$\Delta\Theta_K(t > 41 \text{ ps}) = \sum_{i=0}^5 A_i e^{-t/\tau_i} \sin(2\pi f_i t - \phi_i) + B e^{-t/\tau_{\text{slow}}} + C \quad (1)$$

Where  $A_i$ ,  $f_i$ ,  $\phi_i$ , and  $\tau_i$  are, respectively, the amplitude, frequency, initial phase, and decay time characterizing the oscillation of the mode  $i$  ( $i = 0, 1, 2, 3, 4$ , and  $5$ ). The term  $B e^{-t/\tau_{\text{slow}}}$  represents the slow exponential change with the characteristic time  $\tau_{\text{slow}}$  and  $C$  is an offset. The corresponding fitting with Eq. (1) is plotted in Fig. 3(a) with a red solid line, showing a good agreement with the experimental data. It yields the frequencies  $f_0 = 0.8 \text{ GHz}$ ,  $f_1 = 2.2 \text{ GHz}$ ,  $f_2 = 5.4 \text{ GHz}$ ,  $f_3 = 9.6 \text{ GHz}$ ,  $f_4 = 15.5 \text{ GHz}$ , and  $f_5 = 64 \text{ GHz}$ , which are similar to the ones obtained by the FFT analysis.

To have more information about these modes, we measured the pump-induced change in the reflectivity signal [see Fig. 3(b)]. Clear oscillations with a frequency of  $\sim 64 \text{ GHz}$  are observed from the time delay  $t = 41 \text{ ps}$ , as seen in the FFT spectrum displayed in the inset of Fig. 3(b). These oscillations can be related to the time-domain Brillouin scattering, which originates from the interference of the probe beam reflected from the Bi-YIG surface with the reflection from the longitudinal strain pulse propagating into the GGG. Indeed, the observed frequency perfectly matches the value  $f_{\text{GGG}} = 2V_{\text{GGG}}^L \sqrt{n^2 - \sin^2\theta}/\lambda = 63.9 \text{ GHz}$ , calculated from the probe wavelength  $\lambda = 400 \text{ nm}$ , the longitudinal sound velocity  $V_{\text{GGG}}^L = 6400 \text{ m/s}$  of GGG [54], its refractive index  $n \approx 2$  [55], and the angle of incidence  $\theta = 6^\circ$ . Based on these observations, we can conclude that the highest frequency mode in the TR-MOKE is due to the propagation of the longitudinal acoustic pulse in the GGG substrate, since it has the same frequency as the Brillouin oscillation. We note that the observation of Brillouin oscillations in a TR-MOKE signal was previously reported in the literature [29–31,56]. It can be understood considering a small difference in the reflection of right ( $\sigma^+$ ) and left ( $\sigma^-$ ) helicity of light induced by the acoustic pulse during its propagation in the sample [30,56].

The frequencies of the low-frequency modes are in the range of SW frequencies in Bi-YIG films [43,44,57,58]. To clarify the nature of these modes, we investigated their frequencies as a function of the mode number. Interestingly, the precession frequencies of these modes have a quadratic dependence on the mode number. Such a variation is well described by the dispersion relation of SSW [36,37], which is written in the case of our experimental configuration as [36,37]:

$$\omega(k) = \omega_0 + \gamma D_{\text{ex}} k^2, \quad (2)$$

where  $\omega$  is the angular precession frequency,  $\omega_0$  is the angular frequency of the fundamental mode which depends on  $H_{\text{ext}}$  and the effective magnetic anisotropy field,  $\gamma$  is the gyro-magnetic ratio,  $D_{\text{ex}}$  is the exchange stiffness, and  $k = n\pi/d$  represents the wave-vector of the standing SWs given by  $d = 135 \text{ nm}$ . Indeed, a fit of experimental data with Eq. (2) using  $(\gamma/2\pi) = 28 \text{ GHz/T}$  yields an excellent agreement with the experimental results (Fig. 4). We obtained an exchange stiffness  $D_{\text{ex}} = (5.9 \pm 0.2) \cdot 10^{-17} \text{ T m}^2$  in good agreement with those reported in the literature for YIG and Bi-YIG [49,59,60]. This clearly demonstrates the SSW nature of the observed magnetic resonance modes. It also shows that we can excite SSWs with a frequency up to  $15 \text{ GHz}$ , which is

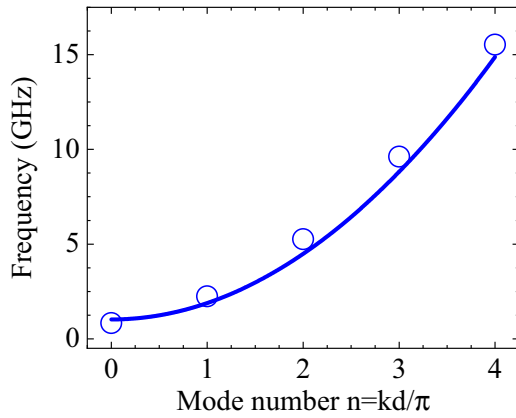


FIG. 4. Spin precession frequency as a function of the mode number. The solid line is the fit obtained using Eq. (2).

15 times higher than the one associated with the ferromagnetic resonance mode (i.e.,  $n = 0$ ). The general excitation mechanism of the SSWs can be related to an ultrafast change of magnetic anisotropy via coherent and incoherent phonons. The coherent phonons correspond to the strain wave shown in Fig. 1(c), and it changes the magnetic anisotropy via the inverse magnetostriction. The incoherent phonons correspond to the phonon-temperature change shown in Fig. 1(d). They change the magnetic anisotropy due to the temperature dependence of magnetic anisotropy constants. We discuss below that the physics behind the generation of incoherent phonons should combine the heat energy directly deposited by the hot-electron pulse at the Bi-YIG surface together with the heat energy transferred from Cu to Bi-YIG.

Let us now discuss the fast and slow decrease in the  $\Delta\Theta_K$  signal. To give a comprehensive discussion of possible mechanisms behind these phenomena, it is worth mentioning first that the fast decrease and recovery in the  $\Delta\Theta_K$  signal are similar to the ultrafast demagnetization process usually induced in metallic magnetic films by ultrashort laser [1] and hot-electron [24] excitations. The exact mechanism of this ultrafast demagnetization is still under debate, but different models have been proposed with the most used being the so-called three-temperature model [1]. This model is purely phenomenological and based on the heat transfer between three interacting subsystems representing electrons, lattice (i.e., phonon), and spins. Such an energy transfer leads to an ultrafast increase of spin temperature, which results in an ultrafast decrease of the magnetic order according to the dependence of the magnetization on the temperature. In magnetic insulators where no free electrons are present, the heating of the spin occurs mainly via the phonon-magnon coupling, which is very weak in the case of iron garnets [61–63]. Therefore, this mechanism cannot be the origin of the subpicosecond decrease of the magnetic order. However, the weak phonon-magnon coupling characterizing iron garnets is in good agreement with the slow decrease in the  $\Delta\Theta_K$  signal, which can therefore be related to a slow demagnetization process. These slow demagnetization dynamics are consistent with the generation and propagation of incoherent phonons (heat) into the Bi-YIG layer [see Fig. 1(d)]. We note that Fig. 1(d) does not include the process of generating incoherent

phonons directly deposited by the hot-electron pulse at the Bi-YIG surface. Including this process in the simulations is a complex task and will essentially increase the temperature in the first nanometer of Bi-YIG, which could compensate a reduced interface conductance for acoustic phonons at the metal-insulator contact. Due to the absence of free electrons in insulators, heat transport in Bi-YIG is identical to incoherent phonon transport. The decrease of the magnetic order is expected to be inhomogeneous across the Bi-YIG thickness and follow the phonon temperature gradient [Fig. 1(d)] with much larger demagnetization close to the Cu interface.

Apart from the phenomenological three-temperature model, the ultrafast demagnetization dynamics are explained from a microscopic point of view using processes involving fundamental interactions at the origin of magnetism: the spin-orbit interaction and the exchange interaction [64–67]. The spin-orbit interaction should only make a small contribution to the ultrafast loss of the magnetic order observed in our sample. Similar ultrafast decrease in the  $\Delta\Theta_K/\Theta_{Kmax}$  signal, i.e., amplitude and timescale, have been measured in heterostructures based on a  $Bi_1Y_2Fe_5O_{12}$  or a  $Bi_3Fe_5O_{12}$  (not shown), in which different spin-orbit couplings are present due to various Bi concentrations. The large difference in spin-orbit coupling has been experimentally verified in  $Bi_xY_{3-x}Fe_5O_{12}$ , in which a huge increase of the MO properties was reported when increasing the Bi concentration [68,69]. We therefore suggest that the ultrafast change of the magnetic order is induced by an ultrafast modulation of the exchange interaction [67]. This mechanism has been recently suggested to explain subpicosecond magnetization dynamics induced in YIG and Bi/Ga-substituted YIG by resonant lattice excitation with an ultrashort terahertz pulse and theoretically supported via atomistic spin-dynamics simulations [67]. The scenario behind the modulation of the exchange interaction can be illustrated as follows. The excited incoherent phonons modulate the distance of adjacent atoms, in particular oxygen atoms. Since the superexchange interaction mediated by these atoms is distance dependent, this perturbation leads to an ultrafast modulation of the superexchange interaction between the two iron sublattices, which result in subpicosecond demagnetization dynamics [67].

To investigate further the ultrafast magnetization dynamics, we studied the fast and slow decrease of the magnetic order as a function of the pump energy density and the external magnetic field (see Fig. 5). The amplitude of both processes linearly increases with increasing  $E_{pump}$ . This is in good agreement with the expected dependence of the heat energy density of hot electrons on  $E_{pump}$ , since it is the main origin of the heat energy generated in our structure. The small nonlinearity in the heat transport originating from the temperature-dependent electronic heat transport is not relevant here, as confirmed by our UXRD studies [32]. On the other hand, we show that the magnetic field dependence of these amplitudes is the same as the static magnetization of the sample [see Fig. 2(a)]. Indeed, they saturate at  $H_{ext} \sim 260$  mT, which is the magnetization saturation field characterizing the Bi-YIG film. These results clearly confirm our prediction about the magnetic origin of these phenomena. It is worth mentioning here that the amplitude of the fast demagnetization is about one order of magnitude smaller than the slow

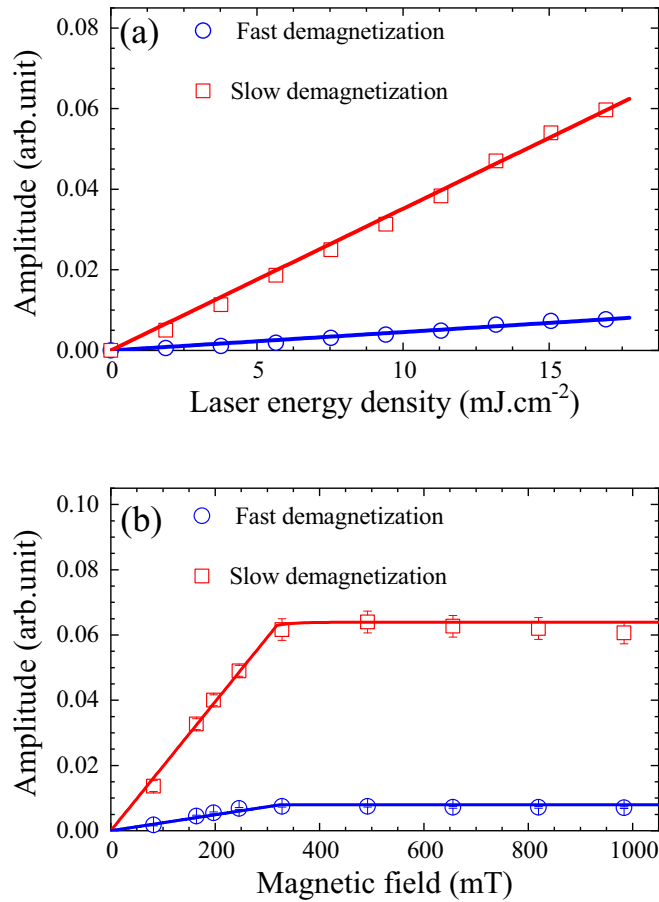


FIG. 5. Amplitude of the fast (circle) and slow (square) demagnetization processes as a function of (a) the pump energy density and (b) the external magnetic field. The solid lines are guides to the eye.

demagnetization for all measured values of  $H_{\text{ext}}$  and  $E_{\text{pump}}$ . This is consistent with a very small amplitude of the motion of oxygen atoms, which is in good agreement with the high stiffness and stability characterizing the garnet structure [70,71].

#### IV. CONCLUSIONS

We have studied the magnetic, optical, and lattice responses of a Pt/Cu/Bi<sub>1</sub>Y<sub>2</sub>Fe<sub>5</sub>O<sub>12</sub>/Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub> heterostructure upon femtosecond laser excitation of the opaque Pt/Cu bilayers. We find that coherent and incoherent phonons generated by the laser-excited metal electrons trigger high-frequency SSWs in the Bi<sub>1</sub>Y<sub>2</sub>Fe<sub>5</sub>O<sub>12</sub> layer via phonon-induced change of the magnetic anisotropy. The incoherent phonons can induce a fast (<1 ps) and slow (>1000 ps) decrease of the magnetic order. We explain these two timescales by considering scenarios based on spin-phonon interaction. Our findings highlight the possibility of triggering high-frequency SWs in magnetic garnets embedded in complex heterostructure devices.

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