



HAL
open science

Laves phase crystal analysis (LaCA): Atomistic identification of lattice defects in C14 and C15 topologically close-packed phases

Zhuocheng Xie, Dimitri Chauraud, Erik Bitzek, Sandra Korte-Kerzel, Julien Guéno le

► To cite this version:

Zhuocheng Xie, Dimitri Chauraud, Erik Bitzek, Sandra Korte-Kerzel, Julien Gu no le. Laves phase crystal analysis (LaCA): Atomistic identification of lattice defects in C14 and C15 topologically close-packed phases. *Journal of Materials Research*, 2021, 10.1557/s43578-021-00237-y . hal-03178464

HAL Id: hal-03178464

<https://hal.univ-lorraine.fr/hal-03178464>

Submitted on 18 May 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destin e au d p t et   la diffusion de documents scientifiques de niveau recherche, publi s ou non,  manant des  tablissements d'enseignement et de recherche fran ais ou  trangers, des laboratoires publics ou priv s.



Distributed under a Creative Commons Attribution 4.0 International License



Laves phase crystal analysis (LaCA): Atomistic identification of lattice defects in C14 and C15 topologically close-packed phases

Zhuocheng Xie¹, Dimitri Chauraud², Erik Bitzek², Sandra Korte-Kerzel¹, Julien Guénoles^{1,3,4,a)} 

¹Institute of Physical Metallurgy and Materials Physics, RWTH Aachen University, 52056 Aachen, Germany

²Department of Materials Science & Engineering, Institute I: General Materials Properties, Friedrich-Alexander-Universität Erlangen-Nürnberg, 91058 Erlangen, Germany

³Université de Lorraine, CNRS, Arts et Métiers ParisTech, LEM3, 57070 Metz, France

⁴Labex Damas, Université de Lorraine, 57070 Metz, France

a) Address all correspondence to this author. e-mail: julien.guenole@univ-lorraine.fr

Received: 6 March 2021; accepted: 3 May 2021

The identification of defects in crystal structures is crucial for the analysis of atomistic simulations. Many methods to characterize defects that are based on the classification of local atomic arrangement are available for simple crystalline structures. However, there is currently no method to identify both, the crystal structures and internal defects of topologically close-packed (TCP) phases such as Laves phases. We propose a new method, Laves phase crystal analysis (LaCA), to characterize the atomic arrangement in Laves crystals by interweaving existing structural analysis algorithms. The new method can identify the polytypes C14 and C15 of Laves phases, typical crystallographic defects in these phases, and common deformation mechanisms such as synchroshear and non-basal dislocations. Defects in the C36 Laves phase are detectable through deviations from the periodic arrangement of the C14 and C15 structures that make up this phase. LaCA is robust and extendable to other TCP phases.

Introduction

Atomistic simulations are routinely used to study the deformation mechanisms of materials at the atomic scale [1–3]. Fundamental to the analysis of such simulations is the ability to identify defects in crystalline structures [4–7]. Many structure characterization methods, such as centrosymmetry parameter, CSP [8] and common neighbor analysis, CNA [9, 10], are widely used for highly symmetric crystalline structures such as face-centered cubic (FCC), body-centered cubic (BCC) and hexagonal close-packed (HCP) crystals. However, to the best of our knowledge, there is currently no method dedicated to the characterization of the local structural environments of topologically close-packed (TCP) phases, in particular of Laves phases.

Laves phases with ideal chemical composition AB_2 are TCP structures with three common polytypes: cubic $MgCu_2$ (C15), hexagonal $MgZn_2$ (C14) and $MgNi_2$ (C36) [11, 12]. The larger A-type atoms are ordered as in cubic and hexagonal diamond structures and the smaller B-type atoms form tetrahedra around

the A-type atoms. The ideal atomic size ratio for Laves phase formation is 1.225 but encompasses a range from 1.05 to 1.68 [13, 14]. Laves phases have a layered structure which consists of quadruple atomic layers along the basal and {111} planes for the hexagonal and cubic Laves phases, respectively (see Fig. 1a). The quadruple atomic layers consist of a single layer of B-type atoms forming a Kagomé net and a triple layer with an A–B–A structure [15, 16]. The atomic distances within the triple layer are much smaller than the distance between the single and triple layers. The different compositions and stacking orders of the quadruple layers lead to the various polytypes of Laves phases.

Laves phases form in many alloys and have a large impact on their mechanical properties [12, 17–19]. The understanding of the underlying deformation mechanisms is thus crucial for tailoring material properties. The synchroshear mechanism of the basal slip was predicted by Hazzledine and Pirouz [15] for Laves phases, which was confirmed by their experimental observation of synchro-Shockley dislocation in the C14 $HfCr_2$

Laves phase [16] and by the ab initio calculations performed by Vedmedenko et al. [20]. After synchroshearing a C14 Laves phase on the alternate triple layer, a strip of C15 structure is formed as a stacking fault. Additionally, Schulze and Paufler et al. [21–23] first reported non-basal slip in the C14 MgZn₂ Laves phase. Zehnder et al. [24] also observed the activation of non-basal slip and measured the corresponding critical resolved shear stresses in the C14 Mg₂Ca Laves phase. Zhang et al. [25, 26] revealed slip mechanisms for prismatic and pyramidal dislocations in a C14 Laves phase ($M_2\text{Nb}$, $M = \text{Cr}$, Ni, and Al).

In addition to those experimental studies and ab initio calculations, the deformation of Laves phases was also investigated using classical atomistic simulations. This was in part enabled by the recently developed modified embedded atom method (MEAM) potential by Kim et al. [27] that accurately describes the mechanical properties of the C14 Mg₂Ca Laves phase. Using this potential, synchroshear was identified as the most favorable mechanism for basal slip in the C14 Mg₂Ca Laves phase [28]. This potential was also used to study the mechanisms of slip transmissions from an Mg matrix to the C14 Mg₂Ca Laves phase at high temperatures using molecular dynamics (MD) simulations [29]. However, in these studies, the defects had to be identified using the von Mises shear strain invariant based on atomic-level strain tensors [30], as commonly used structural characterization methods such as CNA [9, 10], CSP [8] and Voronoi index [31, 32], are not able to identify Laves phase crystal structures or to detect their crystallographic defects. Hammerschmidt et al. [33] recently developed a method to characterize the crystal structures of TCP phases, including Laves phases, by computing the moments of the electronic density-of-states with analytic bond-order potentials on density-functional theory relaxed structures. Their approach is, however, not targeted toward identifying crystalline defects and is limited to systems with few hundreds of atoms.

With the aim of providing a better access to the deformation mechanisms of TCP phases via atomistic simulations, we chose Laves phases as archetype for TCP phases to develop a new structural analysis method named Laves phase crystal analysis (LaCA). By classifying the atoms according to the neighbor list obtained from the combination of existing structural characterization methods, the different Laves phase crystal structures and their typical crystallographic and chemical defects, e.g., anti-site atoms, stacking faults and dislocations, can be identified. In “Laves phase crystal analysis” section, we outline the workflow of the new analysis method. In “Assessments” section, we assess the method with dedicated atomistic simulations of Laves phases and comparisons with existing analysis methods. The details of the implementation and the simulation methods are introduced in the “Methodology” section.

Laves phase crystal analysis

Identification of atomic clusters

To characterize the local structural environment of Laves phases, the CNA is used to identify the characteristic atomic clusters. The CNA (also known as the Honeycutt–Andersen bond type [9]) describes the number of shared neighboring atoms and the bonding relationships among the shared neighbors of a pair of atoms. The local environment of a pair of atoms is classified by three indices, $ijkl$, where j indicates the number of neighboring atoms the paired atoms have in common, k indicates the number of bonds among the shared neighboring atoms, and l indicates the number of bonds in the longest bond chain formed by the k bonds among the shared neighbors. For instance, the only bonded pairs in a FCC crystal are $12 \times$ type 421, and a HCP crystal has equal numbers ($6 \times$) of type 421 and type 422. This has been widely used to identify FCC [34, 35], HCP [36, 37] and BCC [38, 39] crystal structures. Laves crystal structures (C14, C15 and C36) consist of Z12 and Z16 Frank–Kasper (FK) clusters [40] as shown in Fig. 2, in which the Z12 FK (icosahedral) cluster has $12 \times$ type 555 and the Z16 FK cluster has $12 \times$ type 555 and $4 \times$ type 666. A- and B-type atoms are the central atoms of Z16 and Z12 FK clusters in AB₂ Laves phases, respectively. We modified the implementation of the CNA in the Open Visualization Tool (OVITO) [41] to be able to identify Z12 and Z16 FK clusters in Laves phases. Details of the implementation of the modified CNA can be found in the “Methodology” section.

Alternatively, the Voronoi index can be used to determine the local atomic environment by identifying the Voronoi polyhedron enclosing an atom. A Voronoi polyhedron can be expressed using a vector of Voronoi indices ($n_1, n_2, n_3, n_4, n_5, n_6, \dots$), where n_i is the number of polyhedron faces with i edges. The Voronoi cell of a central atom of the Z12 FK cluster has 12 faces with 5 edges each, its index vector is (0, 0, 0, 0, 12, 0). For a central atom of the Z16 FK cluster, the Voronoi cell has 12 faces with 5 edges each and 4 faces with 6 edges each, and the corresponding Voronoi index vector is (0, 0, 0, 0, 12, 4).

Centrosymmetry

The spatial and chemical distributions of neighboring atoms around the central atom of the Z16 FK cluster are the same in these three Laves phases. However, the local structural environment around the central atom of the Z12 FK cluster in C15 is different than in C14 and C36 Laves phase crystal structures. As shown in Fig. 2, C14 and C36 Laves phases have two types of Z12 FK clusters with different chemical distributions of neighboring atoms around the central atoms. These two Z12 FK clusters are classified as B1- and B2-centered clusters (with central atoms B1 and B2). The C15 Laves phase only has B1-centered Z12 FK clusters. To differentiate B1- and B2-centered clusters

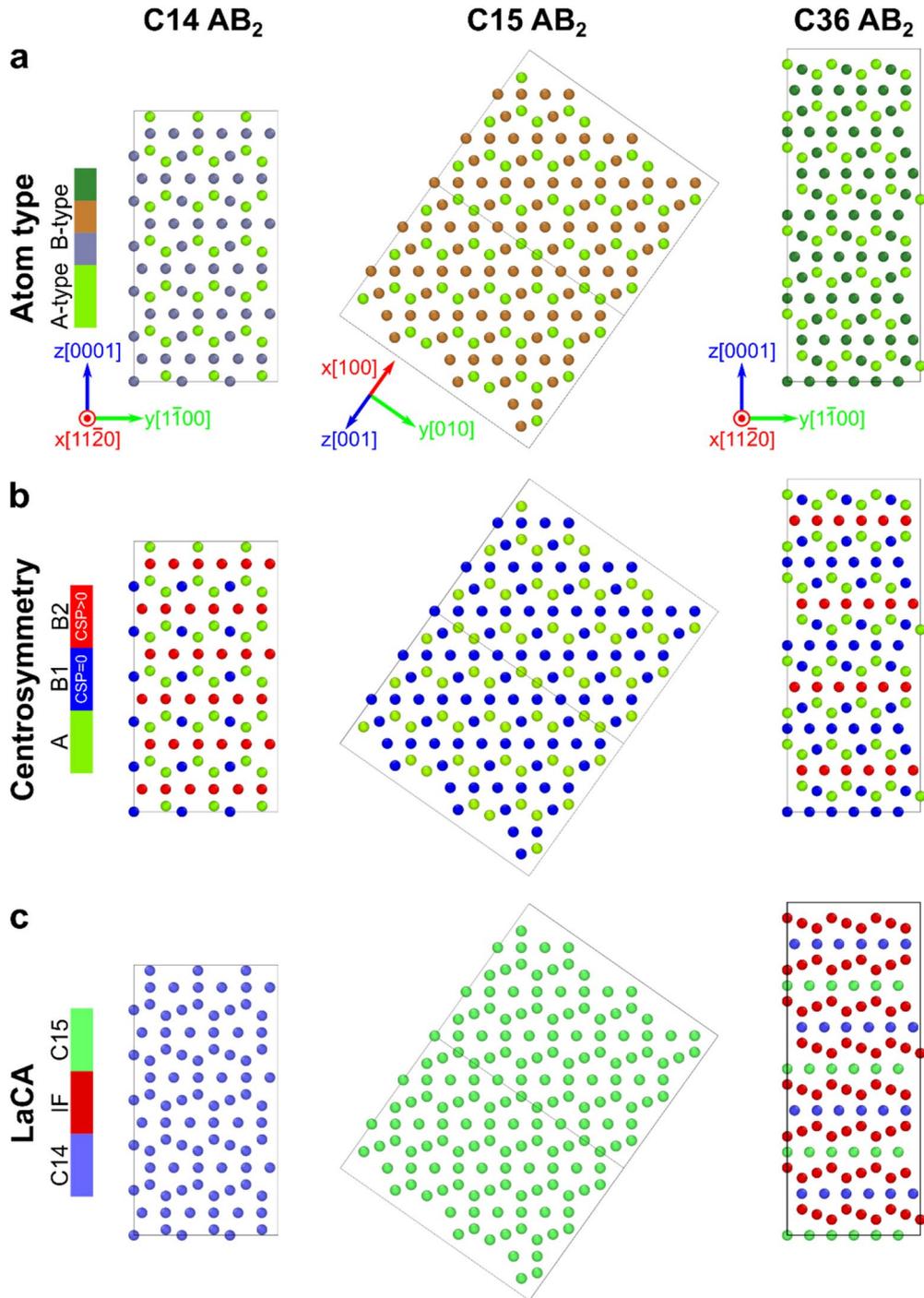


Figure 1: Laves phase crystal structures C14 ($3 \times 3 \times 3$ unit cells), C15 ($3 \times 3 \times 3$ unit cells) and C36 ($3 \times 3 \times 2$ unit cells). (a) Colored by particle type. A-type and B-type atoms are colored bright green and different subdued shades, respectively. (b) Colored by centrosymmetry parameter (CSP) which is calculated on the central atoms of Z12 Frank–Kasper (FK) clusters. B1-type and B2-type atoms are colored blue and red, respectively. (c) Colored by Laves phase crystal analysis (LaCA) method. Atoms in C14, interface (IF) between C14 and C15, and C15 Laves crystal structures are colored blue, red, and green, respectively.

mathematically, the centrosymmetry parameter (CSP) is used. By providing the number of nearest neighbors (N), the CSP is calculated by:

$$CSP = \sum_{i=1}^{N/2} |r_i + r_{i+N/2}|^2, \quad (1)$$

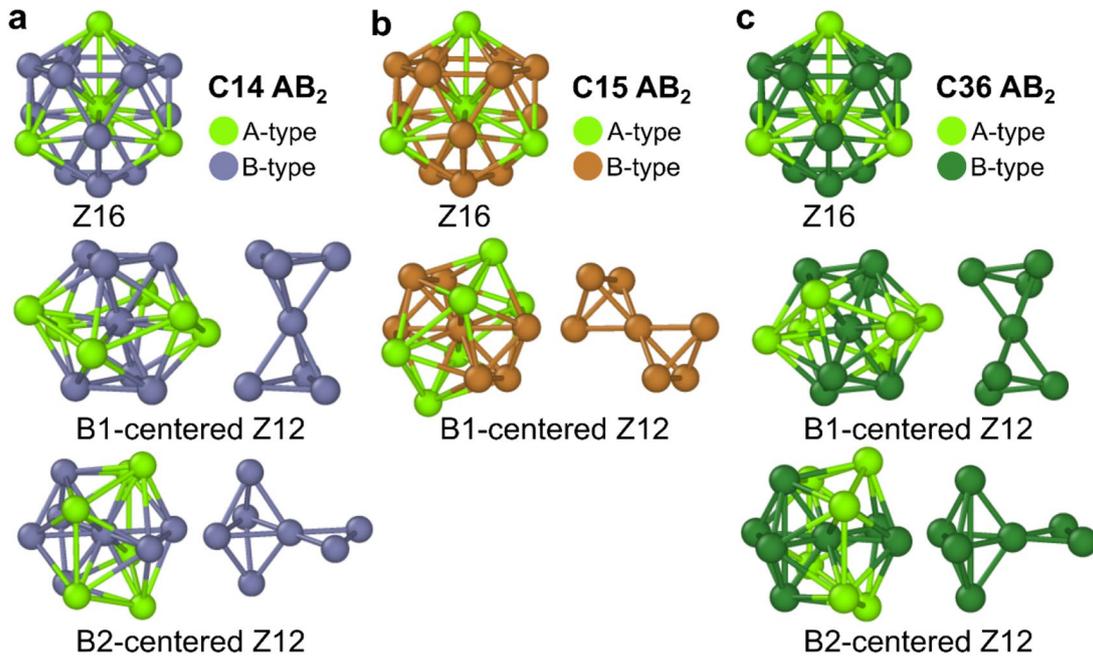


Figure 2: Characteristic atomic clusters in Laves phase crystal structures (a) C14, (b) C15 and (c) C36. B-type atoms in the Z12 icosahedral clusters are shown on the right side of the individual Z12 icosahedral clusters.

where r_i and $r_{i+N/2}$ are vectors corresponding to a pair of opposite nearest neighbors [8]. The ideal centrosymmetric crystal structure has a CSP value of zero since the contributions of all nearest neighboring pairs are canceled out. The CSP is calculated on B-type atoms with $N=6$ nearest neighbors (only B-type atoms are considered). B1-type atoms have a CSP value of zero due to the centrosymmetric structure and, in contrast, B2-type atoms have a CSP value larger than zero. As shown in Fig. 1b, C14, C15 and C36 Laves crystal structures can be differentiated by color coding according to the CSP values of B1- and B2-type atoms.

Neighbor list

To characterize the Laves phase crystal structures quantitatively, neighbor lists of atoms in the characteristic structures are built by considering the type of atomic cluster and the CSP value (see Table 1). The neighbor list is constructed by considering the number of nearest neighbors as defined by the modified CNA classification, e.g., for central atoms of the Z12 and Z16 FK clusters, 12 and 16 nearest neighbors are visited, respectively. The classifications include atoms in C14 (A-, B1- and B2-types), C15 (A- and B1-types), interface (IF) between C14 and C15 (A1-, A2- and B1-types), Other Laves (OL) and Other. For example, if a central atom (A-type) of one Z16 FK cluster has a neighbor list of $4 \times$ A-type, $3 \times$ B1-type, and $9 \times$ B2-type neighboring atoms, this A-type atom is classified as C14 Laves phase. Figure 1c shows C14, C15 and C36 Laves phases characterized by neighbor listing. C14 and C15 Laves phases can be directly recognized, whereas

the C36 Laves phase consists of a periodic succession of C14 and C15 structures.

In the classification of the neighbor list, small perturbations due to thermal fluctuation and elastic deformation are allowed. The neighbor list is treated as a vector where the numbers of A-, B1- and B2-type neighbors are the first, second and third components, respectively. The Euclidean distance between the referenced vector and the vector of interest is calculated:

$$d_{\text{Eucl}}(\mathbf{I}, \mathbf{R}) = \sqrt{\sum_{i=1}^N (\mathbf{R}_i - \mathbf{I}_i)^2}, \quad (2)$$

where \mathbf{R} is the referenced vector and \mathbf{I} is the vector of interest. If the Euclidean norm is below the threshold of similarity (see Table 1), the atom is classified as the referenced Laves crystal structure, otherwise the atom is classified as OL.

Other Laves (OL) indicates an atom identified as the central atom of Z12 or Z16 FK clusters that does not match the referenced neighbor lists. Other indicates an atom not identified as the central atom of Z12 or Z16 FK clusters. The Euclidean distance between the referenced vector and the vector of interest is calculated. If the Euclidean distance is below the threshold of similarity, the atoms belong to the characteristic Laves crystal structures.

Dislocation analysis

The Dislocation Extraction Algorithm, DXA [42] is applied to atomistic configurations to extract dislocation lines and determine their Burgers vectors. In brief, the algorithm tries to build

TABLE 1: Neighbor lists of atoms in Laves phase crystal structures: C14, C15 and interface (IF) between C14 and C15.

Type	Atomic cluster	CSP of B-type atom	Referenced neighbor list	Referenced vector	Threshold of similarity
C14, A	Z16	–	4×A, 3×B1, 9×B2	(4, 3, 9)	$d_{\text{Eucl}} \leq \sqrt{2}$
C14, B1	Z12	≈ 0	6×A, 0×B1, 6×B2	(6, 0, 6)	$d_{\text{Eucl}} \leq \sqrt{2}$
C14, B2	Z12	$\gg 0$	6×A, 2×B1, 4×B2	(6, 2, 4)	$d_{\text{Eucl}} = 0$
C15, A	Z16	–	4×A, 12×B1, 0×B2	(4, 12, 0)	$d_{\text{Eucl}} \leq \sqrt{2}$
C15, B1	Z12	≈ 0	6×A, 6×B1, 0×B2	(6, 6, 0)	$d_{\text{Eucl}} \leq \sqrt{2}$
IF, A1	Z16	–	4×A, 6×B1, 6×B2	(4, 6, 6)	$d_{\text{Eucl}} \leq \sqrt{2}$
IF, A2	Z16	–	4×A, 9×B1, 3×B2	(4, 9, 3)	$d_{\text{Eucl}} \leq \sqrt{2}$
IF, B1	Z12	≈ 0	6×A, 3×B1, 3×B3	(6, 3, 3)	$d_{\text{Eucl}} = 0$
OL	Z12 or Z16	–	Else	–	–
Other	Else	–	–	–	–

Burgers circuits around clusters of atoms that have been identified as crystallographic defects by the CNA method, assuming they could form dislocation cores. The implementation of the DXA in the OVITO can extract dislocations in FCC, HCP, BCC, diamond cubic, and diamond hexagonal crystal structures. For Laves phases, the central atoms of Z16 FK clusters (A-type atoms) are ordered as in diamond structures (diamond hexagonal for C14 Laves phase and diamond cubic for C15 Laves phase). Moreover, the crystallographic information of Laves phases can be solely defined based on the A-type atoms. Therefore, the dislocations and their Burgers vectors in Laves crystal structures can be extracted by performing the DXA on the central atoms of Z16 FK clusters which are classified using the modified CNA presented in this work.

Transferability of LaCA

The C36 hexagonal unit cell is made of a periodic stacking along the *c* axis of one C14 unit cell and one C15 unit cell. The identification of the C36 phase thus requires a larger cutoff able to include this periodic stacking when constructing the neighbor list. Any defects formed in the C36 phase will either break this stacking or directly create a defect in one of the two sub-phases. It remains however unclear how to precisely identify defects in the C36 phase, in particular when considering planar defects such as stacking faults or twins. For example, by considering the glide of a synchroshear partial dislocation in the C36 phase, that transforms a C14 layer into a C15 layer: is the resulting structure a C36 phase with a stacking fault, or a dual-phase composed of a C36 and C15 phase? A similar question can be raised with the glide of a synchroshear partial dislocation in the C14 phase: is the resulting defected C14 phase embedding a layer of C15 phase or a layer of C36 phase formed by the C15 layer and an adjacent C14? To our knowledge, no generally agreed-upon definition currently exists. A solution could be to adapt the results of LaCA based on the user input, whether the C36 phase has to be found in the sample or not. Such an approach would be similar

to experimental analysis, where growth and deformation defects cannot always be distinguished based on a post-mortem analysis only, and as such requires prior knowledge on the sample. This is however beyond the development of the LaCA analysis method and thus out of the scope of the current work.

Besides the C36 Laves phase, LaCA can be extended to identify other TCP phases. Indeed, the crystallographic structures of any TCP phases can be classified into coordination polyhedra (denoted by FK clusters) by means of CNA or Voronoi index analysis. Then, the spatial and chemical distribution of the polyhedral units can be identified using the CSP and the neighbor listing. For each new TCP phase to be identified, the relationships between the atom structure types and the FK clusters have to be found, as in Table 1 in the case of Laves phases. Additionally, these parameters need to be adapted on a case-by-case basis:

- the ratio between the CNA adaptive cutoff and the first neighboring distance,
- the number of nearest neighbors when calculating the CSP value,
- the threshold of the CSP value to evaluate the centrosymmetry,
- the referenced neighbor list,
- the threshold of similarity.

Assessments

Identification of typical crystallographic defects

Typical crystallographic defects in Laves phases are summarized in Fig. 3, as identified using the percentage deviation of potential energy to bulk (ΔE_p) and our method LaCA. Details of the simulation method employed to generate these atomistic defects can be found in the “Methodology” section. On the left side of each pair of images in Fig. 3, coloring is based on ΔE_p , while the right side of each pair is colored using LaCA. Point defects such as anti-site atoms and vacancies are widely observed

experimentally in Laves phases, especially in off-stoichiometric compositions and at elevated temperatures [12]. As shown in Fig. 3a, the values of ΔE_p around the vacancies are below 4.1% (~ 0.1 eV) in C14 Mg₂Ca. By applying LaCA, the atoms in the first two nearest neighbors of vacancies are identified as Other. The anti-site defects in C14 Mg₂Ca are identified as shown in Fig. 3b. For perfect Laves crystals, A-type atoms are the center atoms of Z16 FK clusters, and B-type atoms are the center atoms of Z12 FK clusters. As LaCA combines the results of the modified CNA and the particle type information from the atomistic configuration, chemical defects such as anti-site atoms and anti-phase boundaries can be identified.

In Laves phases, plastic deformation occurs via dislocation-mediated processes. Dislocation motion on the basal plane in the hexagonal and the {111} plane in the cubic Laves phases occurs by synchroshear, characterized by two synchro-Shockley dislocations moving on adjacent planes of the triple layer [16, 20, 28]. Figure 3c shows the core structures of two typical synchro-Shockley dislocations with mixed characters (60° and 30°) in C14 Mg₂Ca. Atoms at the dislocation cores show higher values of ΔE_p up to 17.6% (~ 0.5 eV) than the atoms in stacking faults, and are identified as Other using LaCA. The stacking faults (C15 structure), which are bounded by the synchro-Shockley partial dislocations, are identified using LaCA, as well as the interface between the C15 stacking fault and the C14 matrix. Due to the low stacking fault energy of C14 Mg₂Ca (14 mJ/m²), the ΔE_p of the atoms in stacking faults is less than 0.3% (~ 0.006 eV), see Fig. 3d. The stacking fault and twin boundary in C15 Mg₂Ca are also identified using LaCA. Moreover, atoms at the free surfaces are classified as Other with LaCA (Fig. 3e).

Similarly, point defects, dislocation cores and free surfaces in C15 and C36 Laves phases can be identified by LaCA, since the local atomic packings near these defects are not classified as Z12 or Z16 FK clusters using the modified CNA.

Identification of basal slip in the hexagonal C14 structure: the synchroshear mechanism

The synchroshear mechanism is the most favorable mechanism for basal slip in the C14 Laves phase. It has been confirmed by our previous study using nudged-elastic band (NEB) calculations [28] and experimentally in the Laves phase layer of the μ -phase crystal structure [43, 44]. It is likely to occur in the same way on the basal plane of the C36 phase and the {111} plane of the C15 phase. The minimum energy path (MEP) of a full slip with a Burgers vector of $1/3(11\bar{2}0)$ is associated with two energy maxima corresponding to two separated synchro-Shockley dislocations with Burgers vectors of $1/3(10\bar{1}0)$. Details of the simulations can be found elsewhere [28]. Figure 4 shows the propagation of the leading

and the trailing synchro-Shockley dislocations, and the formation and annihilation of a stacking fault during the slip in a C14 Mg₂Ca pillar are visualized using atomic shear strain, modified CNA, LaCA, and the combination of LaCA and DXA. The atomic shear strain characterizes the deformation gradient of the slip event without further information on the dislocation and stacking fault [Fig. 4a-c(ii)]. By applying the modified CNA, atoms belonging to dislocation cores and free surfaces are separated from the central atoms in Z12 and Z16 FK clusters [Fig. 4a-c(iii)]. The stacking fault with C15 Laves crystal structure, produced by the passage of the leading synchro-Shockley dislocation, is identified using LaCA, see Fig. 4a-c(iv). By combining LaCA and DXA, the synchroshear mechanism is revealed with a dislocation line and a stacking fault [Fig. 4a-c(v)]. In addition, the Burgers vectors of the leading and trailing synchro-Shockley dislocations are correctly characterized as $1/3(10\bar{1}0)$.

Identification of non-basal slip in the C14 structure: slip transmission at Laves phase/metal interfaces

Slip transmission at Laves phase interfaces is important to understand the strengthening effect of Laves precipitates in a metallic matrix. In our previous study, we simulated the co-deformation of Mg₂Ca C14 Laves phase with an a-Mg matrix by means of atomistic simulations in order to explore the transmission of dislocations at the interface [29]. The uniaxial compressive deformation was applied parallel to the matrix-Laves phase interface (along the y -axis in Fig. 5) at low temperature within the microcanonical (NVE) ensemble (initial temperature at 0 K). More details on the simulations can be found elsewhere [29]. These atomistic configurations are of particular interest to assess the LaCA method, as they contain two phases, one interface and the nucleation of a non-basal dislocation in the Laves phase on a prismatic plane. Here, this prismatic slip in the Laves phase is analyzed using atomic shear strain, modified CNA, LaCA, and the combination of LaCA and DXA (Fig. 5). The deformation gradient of the slip transfer is measured using atomic shear strain (Fig. 5b). Mg matrix, Laves phase interface, and C14 Mg₂Ca Laves phase are classified using the modified CNA. The stacking faults and twins in the Mg matrix and the potential dislocation cores in the C14 Laves phase are identified (Fig. 5c). The crystal structure remains C14 after the dislocation glide according to LaCA, see Fig. 5d, which indicates that the dislocation is a perfect dislocation. DXA on the configurations analyzed by the modified CNA with only the central atoms of Z16 FK clusters reveals the dislocation to possess a Burgers vector of $1/3(1\bar{2}10)$, corresponding to a perfect $\langle\alpha\rangle$ -dislocation on the prismatic slip plane (Fig. 5e).

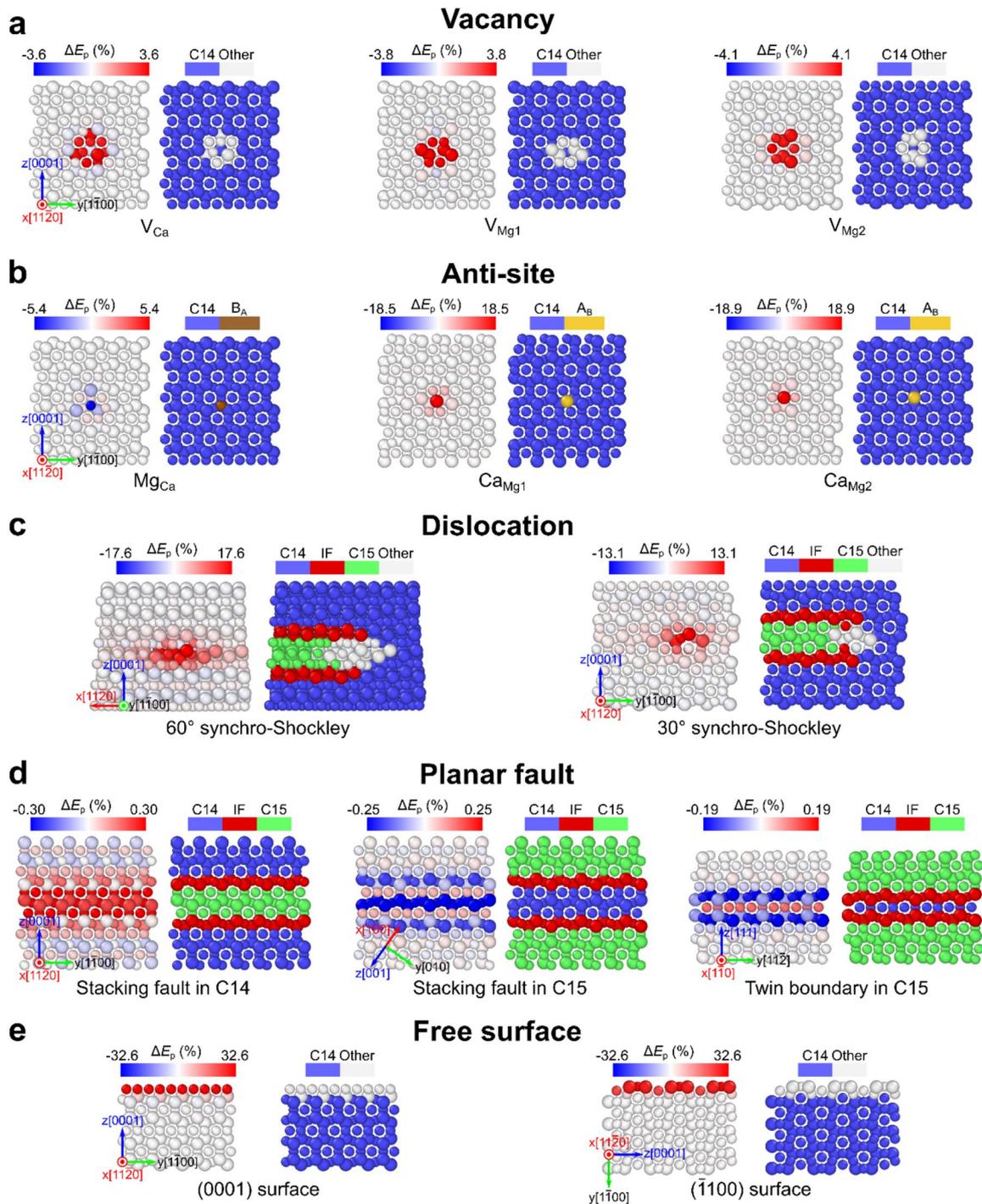


Figure 3: Overview of typical crystallographic defects in Laves phases. Left: colored by percentage deviation of potential energy to bulk (ΔE_p). Right: colored by LaCA. Large and small atoms are A- and B-type atoms, respectively. The atomic size ratio is 1.231 for Mg–Ca systems. (a) Vacancy defects (denoted, e.g., as V_A for vacancy on an A site) in C14 Mg_2Ca . (b) Anti-site defects (denoted, e.g., as A_B for an A atom on a B site) in C14 Mg_2Ca . (c) Synchro-Shockley dislocations in C14 Mg_2Ca . (d) Stacking faults and twin boundary in C14 and C15 Mg_2Ca . (e) Free surfaces in C14 Mg_2Ca .

Identification of unknown defects: nanoindentation

Nanoindentation tests by MD simulations were performed on a Mg_2Ca C14 Laves phase to explore activated plastic events with strain concentration, see Fig. 6a. Details of the

simulation method can be found in the “Methodology” section. The atomistic configuration of the Mg_2Ca C14 Laves phase after nanoindentation is analyzed using atomic shear strain, relative potential energy to bulk ΔE_p , LaCA, and the

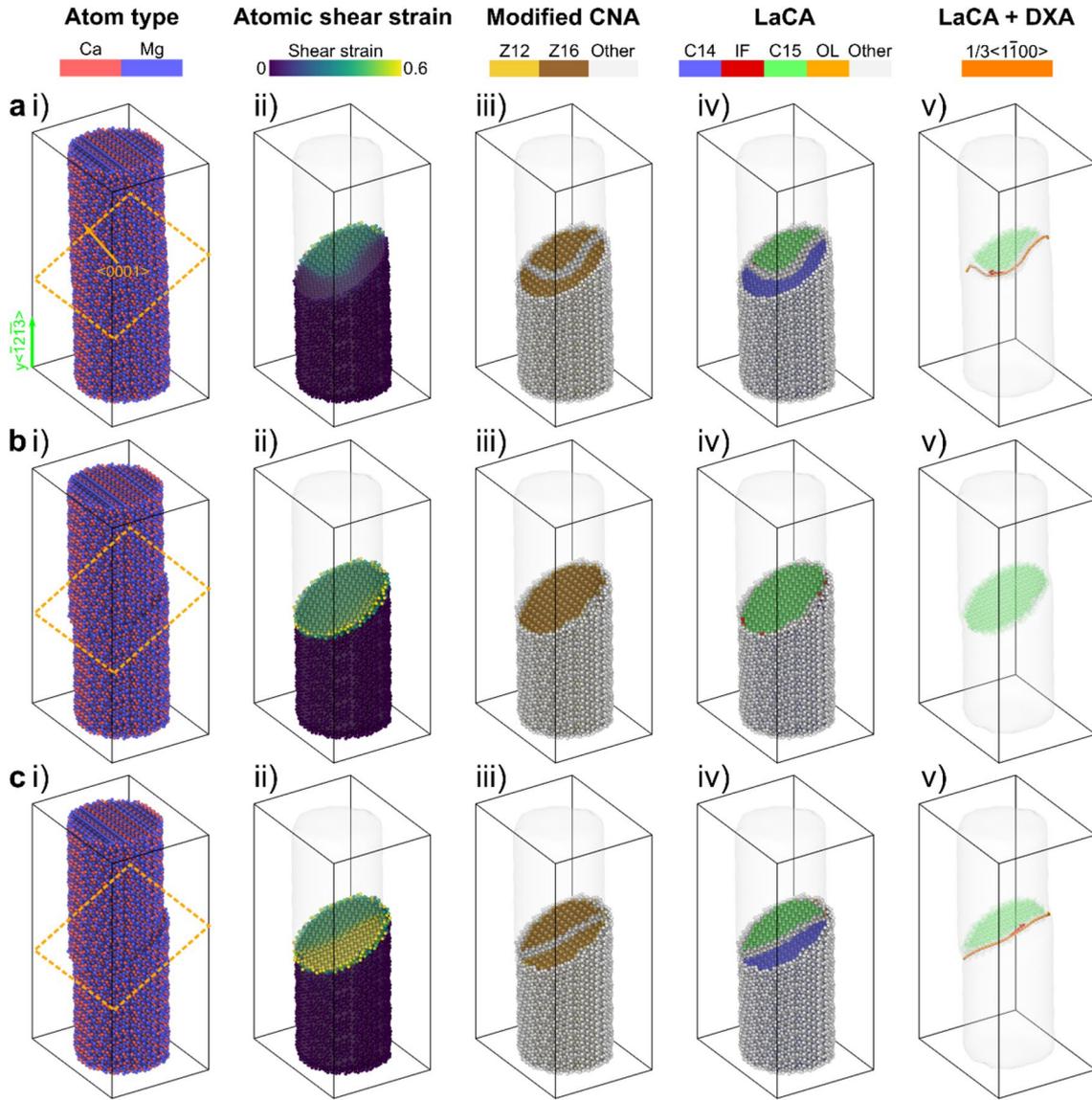


Figure 4: Synchroshear mechanism for basal slip in C14 Mg₂Ca Laves phase as computed by nudged-elastic band (NEB) method. (a) Glide of the leading synchro-Shockley dislocation. (b) Formation of a stacking fault (C15 Laves crystal structure) in the C14 Laves phase after the glide of the leading synchro-Shockley dislocation. (c) Glide of the trailing synchro-Shockley dislocation. (i) Colored by particle type, the dashed box indicates the slice along the basal plane in (ii)–(v). (ii) Colored by shear strain. (iii) Colored by the modified CNA. (iv) Colored by LaCA. (v) DXA was performed on the Laves phase crystal structures as identified by the modified CNA, by considering only the central atoms of the Z16 atomic clusters. Only atoms in stacking faults (C15 Laves crystal structure and Other) are shown in half-transparent. Burgers vector is colored red. The surface mesh of the pillar is shown in half-transparent.

combination of LaCA and DXA, see Fig. 6b–d. As opposed to the post-processed results from shear strain and ΔE_p filtering methods, LaCA successfully captures the overall plastic events under the localized strain. All high-energy atoms were correctly identified by LaCA as being part of a defect configuration as shown in Fig. S1. Moreover, dislocation cores, stacking faults, vacancies and anti-site defects in the deformed sample are recognized using LaCA. Dislocations, such as $\langle a \rangle$ dislocations on basal and prismatic slip planes and synchro-Shockley

dislocations, are characterized using DXA on the configuration identified with the modified CNA.

Identification of defects at high temperatures

The sensitivity of the new method LaCA to thermal fluctuations was tested by analyzing the C14 Mg₂Ca atomistic configurations with defects at 500 K, which is approximately 50% of the melting temperature and the highest temperature considered

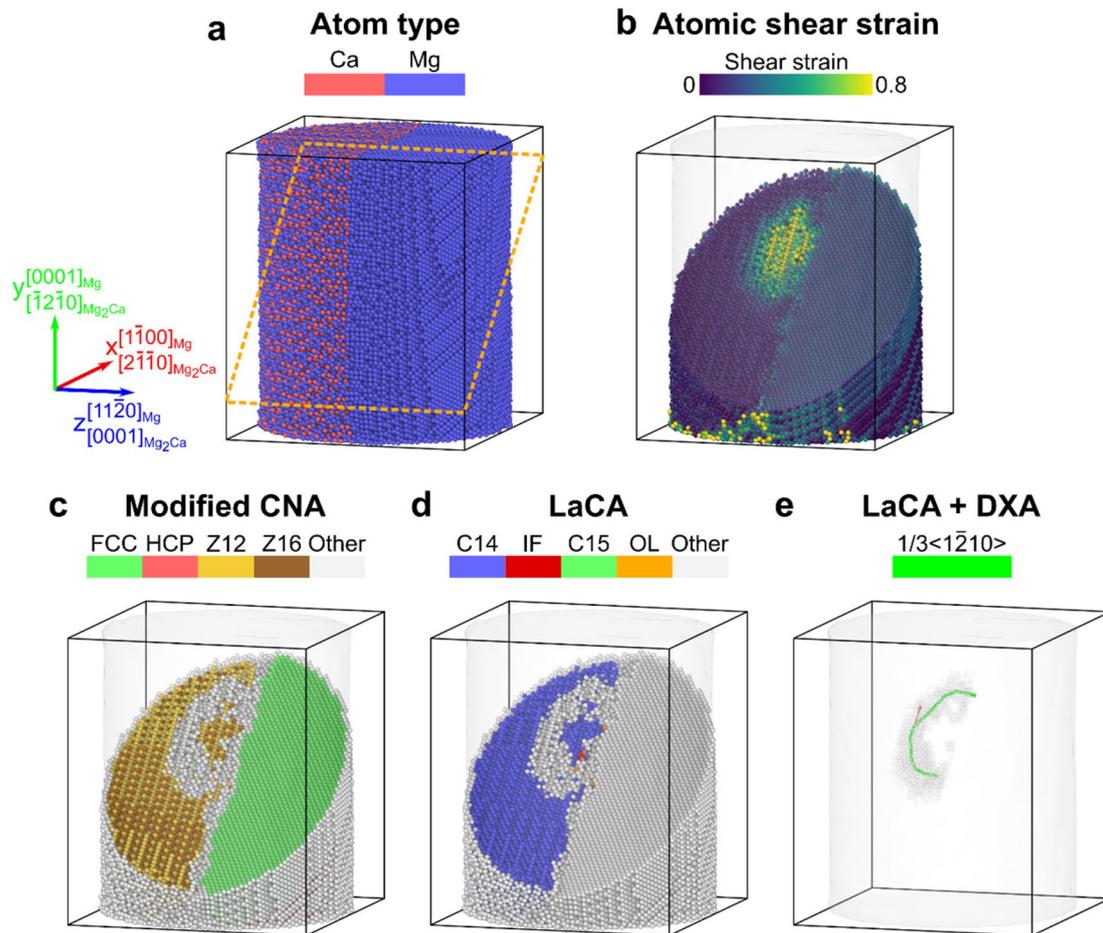


Figure 5: Slip transmission from the basal plane of the Mg matrix to the prismatic plane of the C14 Mg₂Ca Laves phase during a compression test at 50 K using MD simulations. (a) Colored by particle type, the dashed box indicates the slice along the parallel basal (Mg matrix) and prismatic (C14 Mg₂Ca) planes in (b)–(e). (b) Colored by shear strain. (c) Colored by modified CNA. (d) Colored by LaCA. (e) DXA was performed on the post-processed Laves crystal structures using modified CNA, where only central atoms of Z16 atomic clusters were used. Only atoms in dislocation core (atoms identified as Other by LaCA) are shown in half-transparent. Burgers vector is colored red. The surface mesh of the sample is shown in half-transparent.

experimentally for this alloy [19, 45]. Details of the simulation method can be found in the “Methodology” section. The recognition of characteristic atomic clusters as the first step of the algorithm of LaCA can be carried out by the modified CNA or the Voronoi index. Figure 7a, c shows the defect atoms recognized by the CNA-based LaCA in the simulated samples at 500 K. Besides the atoms in the shell of the vacancy (see Fig. 7a) and the stacking fault (see Fig. 7c), a few atoms were also identified as Other due to the thermal noise. In contrast, the Voronoi-based LaCA is less sensitive to the thermal perturbations of the atomic coordinates. Only atoms in the nearest neighbors of the vacancy (Fig. 7b) and the stacking fault (Fig. 7d) were identified as defects.

Although the Voronoi index has certain limitations for the structural analysis of highly symmetric crystalline packings such as FCC and HCP [5, 6], it shows better recognition rates of the Z12 and Z16 FK clusters than the modified CNA at elevated

temperatures. The Voronoi index shows similar outcomes as the modified CNA at low temperatures (see Table S1), but it has a much higher computational cost (one order of magnitude slower) than the CNA. This makes the Voronoi-based LaCA unfavorable for large-scale configurations with millions of atoms. Therefore, we recommend to only employ the Voronoi index to analyze Laves phase structures at high temperatures and when it is not possible nor desired to reduce the thermal noise by other means such as time averaging.

Conclusions

In this work, we present a new analysis method, Laves phase crystal analysis (LaCA), which can be used in atomistic simulations to identify the polytypes of C14 and C15 Laves phases and reveal common crystallographic and chemical defects of

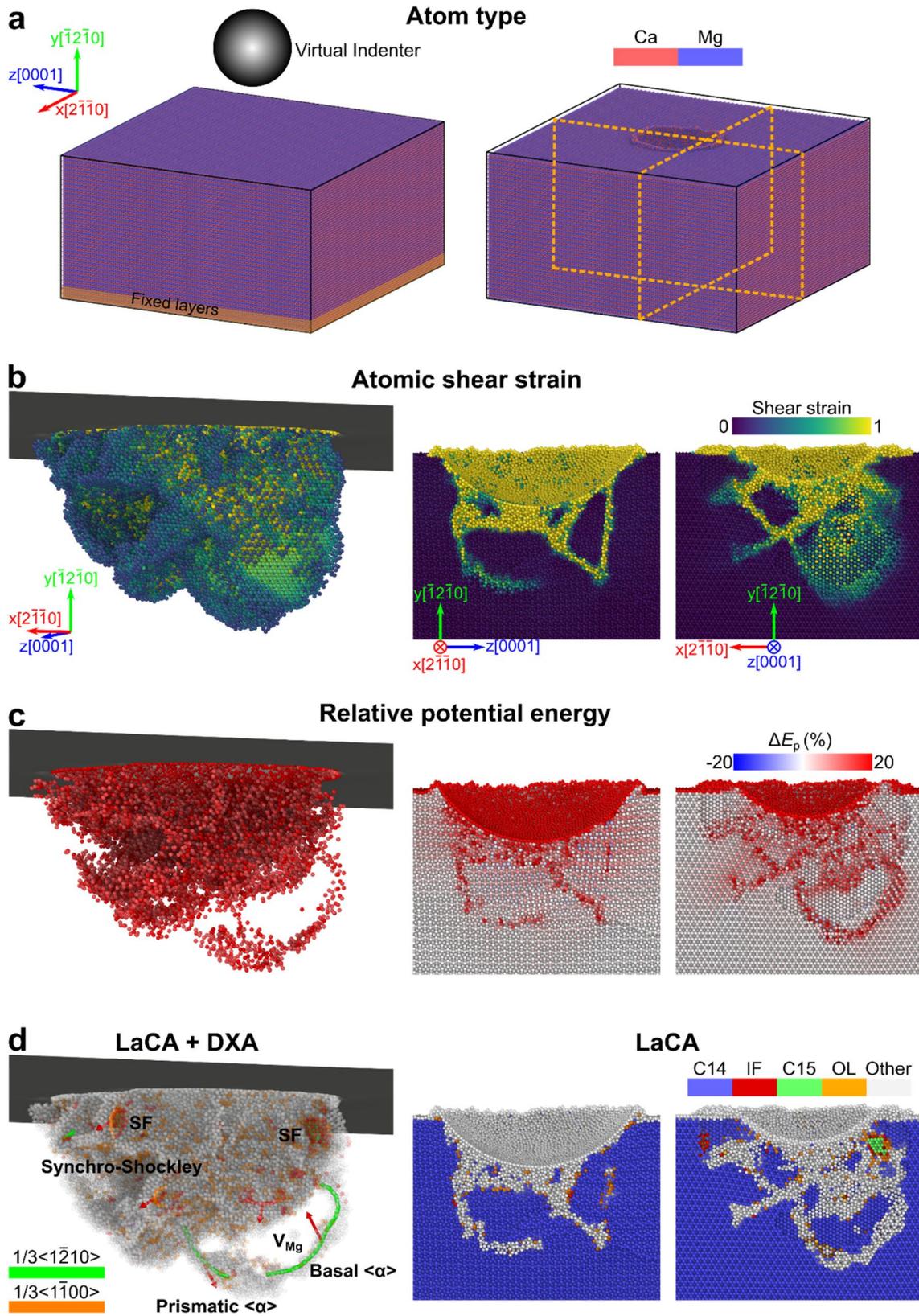


Figure 6: Nanoindentation test on the C14 Mg₂Ca Laves phase. (a) Atomistic sample before and after nanoindentation (indent depth = 5 nm) colored by particle type. The dashed boxes indicate the slices perpendicular to the x and z axes in the right insets of (b)–(d). (b)–(d) Plastic deformation zone beneath the indent. Colored by (b) shear strain (left: only atoms with shear strain > 0.2 are shown), (c) ΔE_p (left: only atoms with $|\Delta E_p| > 12\%$ are shown), and (d) LaCA (left: atoms in C14 structure are not shown) and DXA (Burgers vector is colored red). The right insets show the indent sliced as illustrated in (a).

those. The atomistic structure is classified according to the neighbor list obtained from the close combination of a modified CNA and the CSP. By combining the modified CNA and DXA, the dislocation information can be extracted from plastically deformed Laves phase crystals. The new LaCA method now allows the automatic identification of

- hexagonal C14 and cubic C15 Laves phase crystal structures,
- point, linear and planar defects in the C14 and C15 Laves phases, including
 - basal slip via partial synchroshear dislocations,
 - non-basal slip via perfect dislocations,
 - basal stacking fault formed during synchroshear,
- mixtures of these defects from nanomechanical deformation tests of Laves phases.

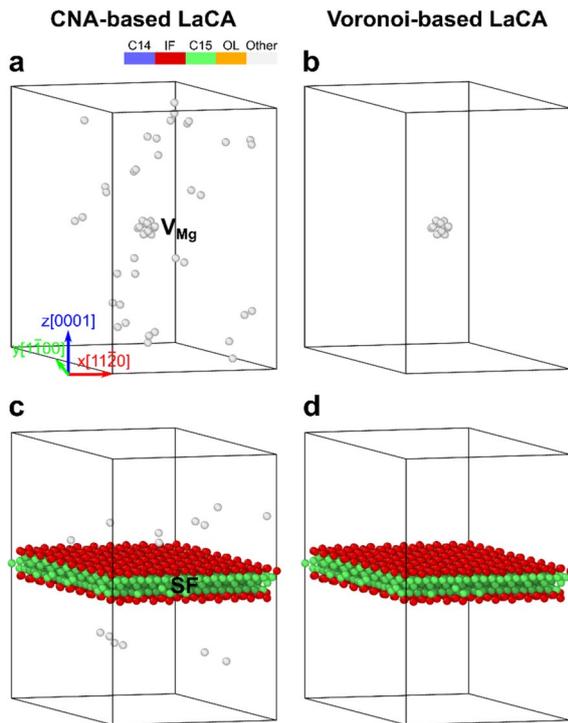


Figure 7: Defect structures of C14 Mg₂Ca atomistic configurations with (a, b) a Mg vacancy and (c, d) a stacking fault at 500 K. Characteristic atomic clusters were identified using (a, c) the modified CNA and (b, d) the Voronoi index. Colored by LaCA.

Defects in the C36 Laves phase are currently not automatically identified, but detectable by eye through deviations from the periodic arrangement of the C14 and C15 structures that make up this phase. The implementation of the method, with a variable cutoff for the modified CNA and the permissible deviation in neighbor listing, makes LaCA robust against small perturbations of atomic coordinates. For atomistic structures with large atomic perturbations, i.e., at elevated temperatures, the Voronoi-based LaCA achieves a higher recognition rate of the local atomic environment than the CNA-based LaCA. Unlike the common structural identification methods, LaCA relies on the unique combination of FK clusters classification, centrosymmetry estimation and chemical order detection. By this, LaCA can not only identify complex crystal structures and structural defects, but also compositional defects like anti-site or anti-phase defects that are not captured by methods like CNA or CSP. Our novel method therefore paves the way for future atomistic studies aimed at understanding the deformation mechanisms of Laves and other TCP phases.

Methodology

Implementation

The new approach in the analysis of crystal structure and lattice defects in Laves phases, as one of the most abundant intermetallic structures, is currently available as a Python program invoking a modified version of the open-source code OVITO.

Variable cutoff of CNA

The modified CNA is implemented in OVITO with adaptive cutoff, in which the optimal cutoff radius is determined automatically [5]. The adaptive cutoff r_{adapt} is set to be halfway between the first and second neighbor shells:

$$r_{\text{adapt}} = \frac{1}{2} (r^{1\text{st}} + r^{2\text{nd}}) \quad (3)$$

and the ratio between r_{adapt} and $r^{1\text{st}}$ is:

$$R = 0.5 + 0.5 \frac{r^{2\text{nd}}}{r^{1\text{st}}} \quad (4)$$

The adaptive cutoff has to be modified to identify Z12 and Z16 FK clusters in Laves phase crystal structures. Figure 8 shows the radial distribution function of the prototype Laves phases and the C14 CaMg₂ Laves phase introduced in “Assessments” section. For a B-centered Z12 FK cluster, the first neighbor shell consists of the first nearest B and A-type atoms within the cutoffs of $r_{\text{B-B}}^{1\text{st}}$ and $r_{\text{A-B}}^{1\text{st}}$, respectively (see Fig. 8a). The second neighbor shell of the B-centered Z12 FK cluster consists of the second nearest B and A-type atoms within the cutoffs of $r_{\text{B-B}}^{2\text{nd}}$ and $r_{\text{A-B}}^{2\text{nd}}$, respectively. In our implementation, the cutoffs of the first and

second neighbor shells of the Z12 FK cluster take the average values of $r_{B-B}^{1st/2nd}$ and $r_{A-B}^{1st/2nd}$. Therefore, the ratio between r_{adapt} and r^{1st} of the Z12 FK cluster is:

$$R_{Z12} = 0.5 + 0.5 \frac{r_{B-B}^{2nd} + r_{A-B}^{2nd}}{r_{B-B}^{1st} + r_{A-B}^{1st}} \quad (5)$$

Similarly, for an A-centered Z16 FK cluster:

$$R_{Z16} = 0.5 + 0.5 \frac{r_{A-A}^{2nd} + r_{A-B}^{2nd}}{r_{A-A}^{1st} + r_{A-B}^{1st}} \quad (6)$$

Taking the radial distribution function of the C14 CaMg_2 atomic sample ($a=6.262 \text{ \AA}$ and $c=9.991 \text{ \AA}$) modeled using the interatomic potential [27] and the prototype Laves phases C14 MgZn_2 ($a=5.221 \text{ \AA}$ and $c=8.567 \text{ \AA}$) [46], C15 MgZn_2 ($a=7.060 \text{ \AA}$) [47] and C36 MgNi_2 ($a=4.824 \text{ \AA}$ and $c=15.826 \text{ \AA}$) [46], $R_{Z12}=1.32$ and $R_{Z16}=1.30$. Based on the calculation, we set the default R values in our implementation as 1.32 for Z12 FK cluster and 1.30 for Z16 FK cluster. It is worth noticing that the choice of cutoff is somewhat arbitrary in the original algorithm [5] and a more recent work proposed using two-third of the distance instead of halfway [48].

CSP threshold

To classify B1- and B2-type atoms, an appropriate threshold should allow for small perturbations due to thermal noises and elastic strains. According to the distribution of probability density of the CSP values of all central atoms of Z12 FK clusters in C14 Mg_2Ca Laves phase crystals as well as in the prototype Laves phases, we set the threshold at 5. With a CSP value greater than 5, the atom is identified as B2-type, otherwise, the atom is identified as B1-type.

Simulation methods

Atomistic simulations were carried out using the software package LAMMPS [49]. The code AtomsK [50] was used to construct the atomistic configurations of the Laves phases. The interatomic interactions were modeled using the modified embedded atom method (MEAM) potential from Kim et al. [27] published in 2015. Energy minimization was performed using a combination of the conjugate gradient (CG) algorithm with box relaxation and the FIRE algorithm [51, 52] until the force norm was below 10^{-8} eV/\AA .

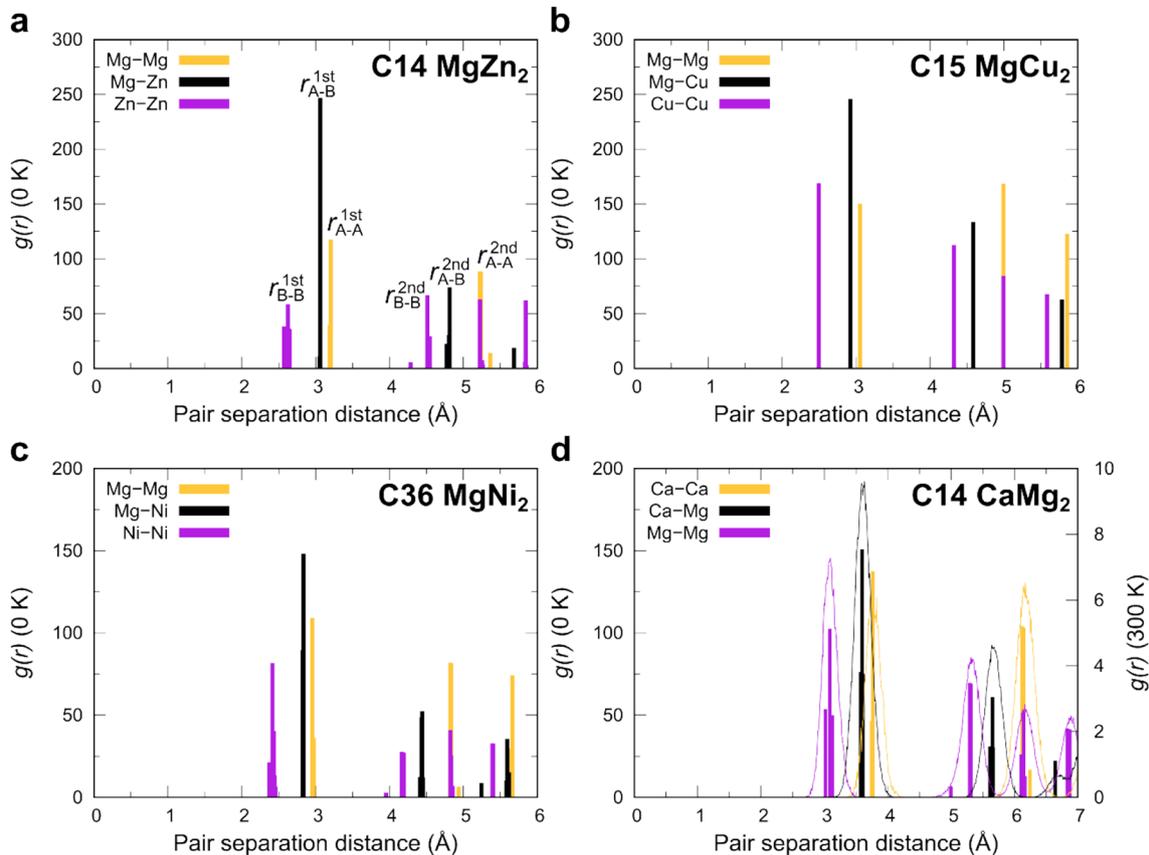


Figure 8: Radial distribution function of the prototype (a) C14 MgZn_2 , (b) C15 MgCu_2 and (c) C36 MgNi_2 Laves phase atomistic structures. (d) Radial distribution function of the C14 CaMg_2 atomistic sample. The curves indicate the distribution function at 300 K.

Typical crystallographic defects

A C14 (C15) Mg₂Ca sample with dimensions of 6.3 × 5.4 × 10.0 nm³ (6.4 × 4.9 × 12.0 nm³) containing 12,000 (13,824) atoms was constructed. Energy minimization was performed after sample generation with box relaxation in all directions.

The point defects were introduced by removing an atom (vacancy) or swapping atom types (anti-site) in the C14 sample with periodic boundary conditions (PBCs) in all directions. Energy minimization was performed on the atomic samples with box relaxation in all directions.

The stacking faults were introduced by shifting one half of the crystal along one partial Burgers vector across the slip planes. The twin boundary in C15 was constructed by merging two mirror symmetric crystals along the [111] plane. PBCs were removed in the direction perpendicular to planar faults. The samples were allowed to relax in the directions perpendicular to the planar faults.

The free surfaces were created by removing PBCs in the surface normal direction. Energy minimization was performed on the atomistic samples with box relaxation in the directions orthogonal to the surface normal direction.

The synchro-Shockley dislocations were obtained from our previous study on the synchroshear mechanism using nudged-elastic band (NEB), where details of the simulations can be found [28].

Nanoindentation

The dimensions of the atomic sample are 61.5 × 31.8 × 59.9 nm³ with around 4.3 Mio atoms. PBCs were applied along x -[2 $\bar{1}\bar{1}$ 0] and z -[0001] directions. A layer of atoms with thickness of 1 nm at the bottom of the sample was rigidly fixed. A spherical indenter exerted a force of magnitude $F(r) = -K(r - R)^2$ on each atom, where K is the specified force constant ($K = 10$ eV/Å³), r is the distance from the center of the indenter to the atom, and R is the radius of the indenter ($R = 10$ nm). The indenter was displaced along the y -[$\bar{1}$ 2 $\bar{1}$ 0] direction toward the sample with a velocity of 10 m/s and a maximum indent depth of 5 nm. The microcanonical (NVE) ensemble was applied during the nanoindentation tests with a time step of 1 fs. Energy minimization was performed on the indented configuration, while keeping the indenter at the maximum depth.

At elevated temperatures

The procedures of sample generation and energy minimization were conducted in a similar way as described above. The atomic samples were heated up to 500 K with a heating rate of 10¹² K/s and stress relaxation in periodic directions within the isothermal–isobaric (NPT) ensemble. Then the samples were annealed at 500 K for 50 ps. Nosé–Hoover thermostat and barostat were used to control the temperature and pressure [53]. A similar approach was applied to heat up the defect-free bulk C14 Mg₂Ca sample to 300 K.

Acknowledgments

The authors acknowledge financial support by the Deutsche Forschungsgemeinschaft (DFG) through the Projects A02, A05 and C02 of the SFB1394 Structural and Chemical Atomic Complexity: From Defect Phase Diagrams to Material Properties, Project ID 409476157. This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 Research and Innovation Programme (Grant Agreement No. 852096 FunBlocks). EB gratefully acknowledges support from the German Research Foundation (DFG) through Projects C3 of the Collaborative Research Centre SFB/TR 103. Simulations were performed with computing resources granted by RWTH Aachen University under Project (rwth0591), by the EXPLOR Center of the Université de Lorraine and by the GENCI-TGCC (Grant 2020-A0080911390). Z.X. wants to thank Dr.-Ing. Wei Luo (RWTH Aachen) for fruitful discussions.

Funding

Open Access funding enabled and organized by Projekt DEAL.

Data availability

We have implemented our new method LaCA for use by other researchers and made our code and relevant atomic configurations available online at <https://doi.org/10.5281/zenodo.4560051> [54].

Declarations

Conflict of interest The authors declare no conflict of interests.

Supplementary Information

The online version contains supplementary material available at <https://doi.org/10.1557/s43578-021-00237-y>.

Open Access

This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit <http://creativecommons.org/licenses/by/4.0/>.

References

1. S. Yip, *Handbook of Materials Modeling* (Springer, Cham, 2007). <https://doi.org/10.1007/978-1-4020-3286-8>
2. C.R. Weinberger, G.J. Tucker, *Multiscale Materials Modeling for Nanomechanics* (Springer, Cham, 2016). <https://doi.org/10.1007/978-3-319-33480-6>
3. H. Huang, H. Van Swygenhoven, Atomistic simulations of mechanics of nanostructures. *MRS Bull.* **34**, 160–166 (2009). <https://doi.org/10.1557/mrs2009.46>
4. J. Li, Atomistic visualization. In: *Handbook of Materials Modeling* (Springer, Dordrecht, 2005), pp. 1051–1068. https://doi.org/10.1007/978-1-4020-3286-8_52
5. A. Stukowski, Structure identification methods for atomistic simulations of crystalline materials. *Model. Simul. Mater. Sci. Eng.* (2012). <https://doi.org/10.1088/0965-0393/20/4/045021>
6. E.A. Lazar, J. Han, D.J. Srolovitz, Topological framework for local structure analysis in condensed matter. *Proc. Natl Acad. Sci. USA* **112**, E5769–E5776 (2015). <https://doi.org/10.1073/pnas.1505788112>
7. A. Stukowski, Visualization and analysis strategies for atomistic simulations. In: *Multiscale Materials Modeling for Nanomechanics* (Springer, Cham, 2016), pp. 317–336. https://doi.org/10.1007/978-3-319-33480-6_10
8. C.L. Kelchner, S.J. Plimpton, J.C. Hamilton, Dislocation nucleation and defect structure during surface indentation. *Phys. Rev. B* **58**, 58 (1998). <https://doi.org/10.1103/PhysRevB.58.11085>
9. J.D. Honeycutt, H.C. Andersen, Molecular dynamics study of melting and freezing of small Lennard-Jones clusters. *J. Phys. Chem.* **91**, 4950–4963 (1987). <https://doi.org/10.1021/j100303a014>
10. D. Faken, H. Jónsson, Systematic analysis of local atomic structure combined with 3D computer graphics. *Comput. Mater. Sci.* **2**, 279–286 (1994). [https://doi.org/10.1016/0927-0256\(94\)90109-0](https://doi.org/10.1016/0927-0256(94)90109-0)
11. A.K. Sinha, Topologically close-packed structures of transition metal alloys. *Prog. Mater. Sci.* **15**, 81–185 (1972). [https://doi.org/10.1016/0079-6425\(72\)90002-3](https://doi.org/10.1016/0079-6425(72)90002-3)
12. F. Stein, A. Leineweber, Laves phases: a review of their functional and structural applications and an improved fundamental understanding of stability and properties. *J. Mater. Sci.* **56**, 5321–5427 (2020). <https://doi.org/10.1007/s10853-020-05509-2>
13. D.J. Thoma, J.H. Perepezko, A geometric analysis of solubility ranges in Laves phases. *J. Alloys Compd.* **224**, 330–341 (1995). [https://doi.org/10.1016/0925-8388\(95\)01557-4](https://doi.org/10.1016/0925-8388(95)01557-4)
14. J.H. Zhu, C.T. Liu, L.M. Pike, P.K. Liaw, A thermodynamic interpretation of the size-ratio limits for Laves phase formation. *Metall. Mater. Trans. A* **30**, 1449–1452 (1999). <https://doi.org/10.1007/s11661-999-0292-5>
15. P.M. Hazzledine, P. Pirouz, Synchroshear transformations in Laves phases. *Scr. Metall. Mater.* **28**, 1277–1282 (1993). [https://doi.org/10.1016/0956-716X\(93\)90468-8](https://doi.org/10.1016/0956-716X(93)90468-8)
16. M.F. Chisholm, S. Kumar, P. Hazzledine, Dislocations in complex materials. *Science* (80-) **307**, 701–703 (2005). <https://doi.org/10.1126/science.11105962>
17. J.D. Livingston, Laves-phase superalloys? *Phys. Status Solidi* **415**, 415–423 (1992). <https://doi.org/10.1002/pssa.2211310215>
18. T.M. Pollock, Weight loss with magnesium alloys. *Science* (80-) **328**, 986–987 (2010). <https://doi.org/10.1126/science.1182848>
19. M. Zubair, S. Sandlöbes, M.A. Wollenweber, C.F. Kusche, W. Hildebrandt, C. Broeckmann, S. Korte-Kerzel, On the role of Laves phases on the mechanical properties of Mg–Al–Ca alloys. *Mater. Sci. Eng. A* **756**, 272–283 (2019). <https://doi.org/10.1016/j.msea.2019.04.048>
20. O. Vedmedenko, F. Rösch, C. Elsässer, First-principles density functional theory study of phase transformations in NbCr₂ and TaCr₂. *Acta Mater.* **56**, 4984–4992 (2008). <https://doi.org/10.1016/j.actamat.2008.06.014>
21. H. Kubsch, P. Paufler, G.E.R. Schulze, The mobility of grown-in dislocations in the intermetallic compound MgZn₂ during prismatic slip. *Phys. Status Solidi* **25**, 269–275 (1974). <https://doi.org/10.1002/pssa.2210250125>
22. T.H. Müller, P. Paufler, Yield strength of the monocrystalline intermetallic compound MgZn₂. *Phys. Status Solidi* **40**, 471–477 (1977). <https://doi.org/10.1002/pssa.2210400213>
23. P. Paufler, Early work on Laves phases in East Germany. *Intermetallics* **19**, 599–612 (2011). <https://doi.org/10.1016/j.intermet.2010.11.032>
24. C. Zehnder, K. Czerwinski, K.D. Molodov, S. Sandlöbes-Haut, J.S.K.L. Gibson, S. Korte-Kerzel, Plastic deformation of single crystalline C14 Mg₂Ca Laves phase at room temperature. *Mater. Sci. Eng. A* **759**, 754–761 (2019). <https://doi.org/10.1016/j.msea.2019.05.092>
25. Y. Zhang, K. Du, W. Zhang, B. Du, D. Qi, W. Li, M. Song, L. Sheng, H. Ye, Shear deformation determined by short-range configuration of atoms in topologically close-packed crystal. *Acta Mater.* **179**, 396–405 (2019). <https://doi.org/10.1016/j.actamat.2019.08.056>
26. Y. Zhang, W. Zhang, B. Du, W. Li, L. Sheng, H. Ye, K. Du, Shuffle and glide mechanisms of prismatic dislocations in a hexagonal C14-type Laves-phase intermetallic compound. *Phys. Rev. B* **102**, 1–9 (2020). <https://doi.org/10.1103/PhysRevB.102.134117>
27. K.H. Kim, J.B. Jeon, B.J. Lee, Modified embedded-atom method interatomic potentials for Mg–X (X=Y, Sn, Ca) binary systems. *CALPHAD Comput. Coupling Phase Diagr. Thermochem.* **48**, 27–34 (2015). <https://doi.org/10.1016/j.calphad.2014.10.001>
28. J. Guérolé, F.Z. Mouhib, L. Huber, B. Grabowski, S. Korte-Kerzel, Basal slip in Laves phases: the synchroshear dislocation. *Scr. Mater.* **166**, 134–138 (2019). <https://doi.org/10.1016/j.scriptamat.2019.03.016>
29. J. Guérolé, M. Zubair, S. Roy, Z. Xie, M. Lipińska-Chwałek, S. Sandlöbes-Haut, S. Korte-Kerzel, Exploring the transfer of plasticity across Laves phase interfaces in a dual phase magnesium

- alloy. *Mater. Des.* (2021). <https://doi.org/10.1016/j.matdes.2021.109572>
30. F. Shimizu, S. Ogata, J. Li, Theory of shear banding in metallic glasses and molecular dynamics calculations. *Mater. Trans.* **48**, 2923–2927 (2007). <https://doi.org/10.2320/matertrans.MJ200769>
 31. G. Voronoi, Nouvelles applications des paramètres continus à la théorie des formes quadratiques. Deuxième mémoire. Recherches sur les paralléloèdres primitifs. *J. Für Die Reine Angew. Math. (Crelles J.)* **1908**, 198–287 (1908). <https://doi.org/10.1515/crll.1908.134.198>
 32. G. Voronoi, Nouvelles applications des paramètres continus à la théorie des formes quadratiques. Premier mémoire. Sur quelques propriétés des formes quadratiques positives parfaites. *J. Für Die Reine Angew. Math.* **1908**, 97–102 (1908). <https://doi.org/10.1515/crll.1908.133.97>
 33. T. Hammerschmidt, A.N. Ladines, J. Koßmann, R. Drautz, Crystal-structure analysis with moments of the density-of-states: application to intermetallic topologically close-packed phases. *Curr. Comput. Aided Drug Des.* (2016). <https://doi.org/10.3390/cryst6020018>
 34. Z. Xie, J. Shin, J. Renner, A. Prakash, D.S. Gianola, E. Bitzek, Origins of strengthening and failure in twinned Au nanowires: insights from in-situ experiments and atomistic simulations. *Acta Mater.* **187**, 166–175 (2020). <https://doi.org/10.1016/j.actamat.2020.01.038>
 35. N.J. Schrenker, Z. Xie, P. Schweizer, M. Moninger, F. Werner, N. Karpstein, M. Mačković, G.D. Spyropoulos, M. Göbel, S. Christiansen, Microscopic deformation modes and impact of network anisotropy on the mechanical and electrical performance of five-fold twinned silver nanowire electrodes. *ACS Nano* (2020). <https://doi.org/10.1021/acsnano.0c06480>
 36. A. Vaid, J. Guérolé, A. Prakash, S. Korte-Kerzel, E. Bitzek, Atomistic simulations of basal dislocations in Mg interacting with Mg₁₇Al₁₂ precipitates. *Materialia* **7**, 100355 (2019). <https://doi.org/10.1016/j.mtla.2019.100355>
 37. Z. Wu, W.A. Curtin, Mechanism and energetics of <c+a> dislocation cross-slip in hcp metals. *Proc. Natl Acad. Sci. USA* **113**, 11137–11142 (2016). <https://doi.org/10.1073/pnas.1603966113>
 38. J.J. Möller, M. Mrovec, I. Bleskov, J. Neugebauer, T. Hammerschmidt, R. Drautz, C. Elsässer, T. Hickel, E. Bitzek, 110 Planar faults in strained BCC metals: origins and implications of a commonly observed artifact of classical potentials. *Phys. Rev. Mater.* **2**, 1–16 (2018). <https://doi.org/10.1103/PhysRevMaterials.2.093606>
 39. R.F. Zhang, J. Wang, I.J. Beyerlein, T.C. Germann, Twinning in BCC metals under shock loading: a challenge to empirical potentials. *Philos. Mag. Lett.* **91**, 731–740 (2011). <https://doi.org/10.1080/09500839.2011.615348>
 40. F.C. Frank, J.S. Kasper, Complex alloy structures regarded as sphere packings. I. Definitions and basic principles. *Acta Crystallogr.* **11**, 184–190 (1958). <https://doi.org/10.1107/S0365110X58000487>
 41. A. Stukowski, Visualization and analysis of atomistic simulation data with OVITO—the Open Visualization Tool. *Model. Simul. Mater. Sci. Eng.* (2010). <https://doi.org/10.1088/0965-0393/18/1/015012>
 42. A. Stukowski, K. Albe, Extracting dislocations and non-dislocation crystal defects from atomistic simulation data. *Model. Simul. Mater. Sci. Eng.* (2010). <https://doi.org/10.1088/0965-0393/18/8/085001>
 43. S. Schröders, S. Sandlöbes, C. Birke, M. Loeck, L. Peters, C. Tromas, S. Korte-Kerzel, Room temperature deformation in the Fe₇Mo₆ μ-phase. *Int. J. Plast.* **108**, 125–143 (2018). <https://doi.org/10.1016/j.ijplas.2018.05.002>
 44. S. Schröders, S. Sandlöbes, B. Berkels, S. Korte-Kerzel, On the structure of defects in the Fe₇Mo₆ μ-phase. *Acta Mater.* **167**, 257–266 (2019). <https://doi.org/10.1016/j.actamat.2019.01.045>
 45. H. Okamoto, T.B. Massalski, *Binary Alloy Phase Diagrams* (ASM International, Materials Park, 1990). <https://doi.org/10.31399/asm.hb.v03.a0006247>
 46. Y. Komura, K. Tokunaga, Structural studies of stacking variants in Mg-base Friauf-Laves phases. *Acta Crystallogr. B* **36**, 1548–1554 (1980). <https://doi.org/10.1107/S0567740880006565>
 47. M.H. Braga, J.J.A. Ferreira, J. Siewenie, T. Proffen, S.C. Vogel, L.L. Daemen, Neutron powder diffraction and first-principles computational studies of CuLi_xMg_{2-x} (x ≅ 0.08), CuMg₂, and Cu₂Mg. *J. Solid State Chem.* **183**, 10–19 (2010). <https://doi.org/10.1016/j.jssc.2009.09.010>
 48. P.M. Larsen, Revisiting the Common Neighbour Analysis and the Centrosymmetry Parameter (2020). ArXiv Prepr. ArXiv2003.08879
 49. S. Plimpton, Fast parallel algorithms for short-range molecular dynamics. *J. Comput. Phys.* **117**, 1–19 (1995). <https://doi.org/10.1006/jcph.1995.1039>
 50. P. Hirel, Atomska: a tool for manipulating and converting atomic data files. *Comput. Phys. Commun.* **197**, 212–219 (2015). <https://doi.org/10.1016/j.cpc.2015.07.012>
 51. E. Bitzek, P. Koskinen, F. Gähler, M. Moseler, P. Gumbsch, Structural relaxation made simple. *Phys. Rev. Lett.* **97**, 170201 (2006). <https://doi.org/10.1103/PhysRevLett.97.170201>
 52. J. Guérolé, W.G. Nöhring, A. Vaid, F. Houllé, Z. Xie, A. Prakash, E. Bitzek, Assessment and optimization of the fast inertial relaxation engine (fire) for energy minimization in atomistic simulations and its implementation in LAMMPS. *Comput. Mater. Sci.* **175**, 109584 (2020). <https://doi.org/10.1016/j.commatsci.2020.109584>
 53. W.G. Hoover, Canonical dynamics: equilibrium phase-space distributions. *Phys. Rev. A* **31**, 1695–1697 (1985). <https://doi.org/10.1103/PhysRevA.31.1695>
 54. Z. Xie, D. Chauraud, E. Bitzek, S. Korte-Kerzel, J. Guérolé, Laves phase Crystal Analysis (LaCA) (2021). <https://doi.org/10.5281/zenodo.4560051>