

Flue gas injection into depleted tight hydrocarbon reservoirs in the context of global warming mitigation: Computed evolution of some reservoir parameters

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- 2 warming mitigation: computed evolution of some reservoir parameters
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- 15 Keywords: depleted oil reservoir, CO₂ storage, temperature profile, oxy-combustion, CFD
- 16

17 Abstract:

CO₂ storage in depleted tight hydrocarbon reservoir seems to be a promising solution to mitigate greenhouse gas emissions. However, flue gas contains not only CO₂ but also minor gaseous impurities due to CO₂ production or capture processes. In the case of oxy-combustion, the main impurity can be O₂ in concentration up to 7 %. O₂ injection into the reservoir can lead to the oxidation of the residual hydrocarbons and therefore, it is necessary to assess the thermal consequences on the reservoir. 23 COMSOL Multiphysics[®] has been used to model an axisymmetric fractured porous reservoir with its 24 cap and base rocks and containing n-octane as a model compound of residual oil. A global kinetic model 25 for *n*-octane oxidation has been derived from a previous detailed free-radical model, and implemented 26 into the reservoir model. Simulated injections of N_2/O_2 mixture (representing a simplified flue gas) 27 have been performed to compute temperature, pressure and n-octane concentration profiles. The 28 results show that the oxidation exothermicity may have a strong influence on temperature profile, especially when the heat capacity of the rock is rather low. In these conditions, the remaining 29 30 hydrocarbons may be consumed in some months. However, the influence of the thermal conductivity 31 seems negligible. Therefore, it appears safer to select reservoirs whose rocks compositions, in 32 particular the cap and base rocks, have a high heat capacity to promote heat dissipation. It is a criterion 33 to consider when selecting a storage site containing residual hydrocarbons, especially for CO₂ captured 34 after oxy-combustion.

35

36 **1. Introduction**

37 The CO₂ emissions from combustion of fossil fuel is still increasing and reached 33.62 Giga ton/year in 38 2019 (IEA, 2022), leading to dramatic climate warming due to CO₂ greenhouse gas effect (IPCC, 2021). 39 The mitigation of such alarmingly increasing CO₂ emissions mainly considers three combined 40 strategies: 1) switching to a low-carbon economy which relies on renewable and alternative energy sources; 2) increasing the efficiency of the current technologies and decreasing the energy loss; 3) 41 42 using CCS (Carbon Capture and Storage) technologies to reduce CO₂ emissions, since the worldwide 43 dependency on fossil fuels is probably going to last several more decades (Maroto-Valer, 2010). CCS 44 appears to be an efficient strategy to mitigate anthropogenic CO₂ emissions from large sources, by up 45 to 90% (Maroto-Valer, 2010). CCS mainly consists of: 1) capturing and potentially separating CO₂ from 46 the other exhaust gases; 2) carrying CO_2 to its storage location, after compression to supercritical 47 conditions; 3) storing CO_2 , e.g. in several types of geological storage sites (Liu et al., 2018).

48 The geological reservoirs that are the most frequently considered are the deep saline aquifers. They 49 may be subjected to regional pressure build-up due to water movement, which may affect the 50 reservoir integrity (Hamza et al. 2021) and the storage capacities. Injecting CO₂ into coal seam 51 reservoirs is also considered, potentially for EMR (Enhanced Methane Recovery) purpose (Zhang and 52 Ranjith, 2019). Another type of geological reservoir that is the focus of attention, is the depleted 53 hydrocarbon (gas or oil) reservoirs. Their storage capacity is much higher than that of coal seams (Kaldi et al., 2009) and they present several advantages over saline aquifers. Indeed, the oil and gas fields 54 55 have proved their potential containment by the retention of hydrocarbons for millions of years. They 56 have been widely studied (storage capacity, structural and stratigraphic patterns...) (Kaldi et al., 2009). 57 Another advantage is the limited pressure inside the reservoirs due to the former hydrocarbon 58 production, which allows injecting CO_2 at moderate pressure, which is technically and economically 59 favorable.

60 Power plants and other industrial processes that use combustion usually produce flue gases at 61 approximately atmospheric pressure, with CO_2 diluted mostly in N_2 . CO_2 must then be separated from 62 other components in the flue gas (post-combustion capture). Another solution is pre-combustion 63 which consists of the removal of carbon from the fuel prior to combustion. At last, oxy-combustion, i.e. combustion in purified O₂ instead of air, is another way to solve this issue. Oxy-combustion enables 64 65 a much easier capture of CO_2 by the flue gas condensation than post-combustion in which N_2 is the 66 dominant component (Tumsa et al., 2017). Nevertheless, flue gas from oxy-combustion may contain a 67 relatively high level of impurities, such as water, Ar (0-5%), N_2 (0-15%), O_2 (0-7%) and SO_2 (0-1.5%) 68 (Mikunda, 2012).

This study focuses on CO_2 captured after oxy-combustion processes, which contains residual O_2 . It aims to assess some effects of the presence of O_2 in the injected gas on some parameters of the gas/oil depleted reservoir and to assess if it is necessary to remove O_2 before injection since purification is expensive. Previous studies already showed that the impurities reduce the structural trapping capacity 73 for CO_2 by reducing the CO_2 density (Wang et al., 2012). Another consequence of the presence of O_2 is 74 examined here. Indeed our previous studies (Pacini-Petitjean et al., 2015a, 2015b, 2016) showed that 75 the oxidation of the hydrocarbons occurs in the temperature conditions of relatively deep hydrocarbon 76 reservoirs (T > 373 K). Therefore, the oxidation kinetics of the residual hydrocarbons may impact on 77 the pressure and temperature evolution in the reservoir during the injection and storage. That is why 78 a model of one part of a fractured porous reservoir has been setup by COMSOL Multiphysics® (version 79 5.6), which considers heat and mass transfers, flow and kinetics. Several assumptions have been made, 80 in particular the flue gas has been replaced by air to emphasize the oxidation phenomenon and to stay 81 in the same conditions as in our previous studies. The other impurities have not been considered. This 82 simplified representation of a flue gas appears appropriate since the study mainly concerns the 83 thermal effects of O₂. An injection well in the center and perforated over the entire height of the 84 reservoir has been considered. The model aims at computing the consumption of O_2 and the hydrocarbons, the pressure increase and the temperature evolution as a function of time, throughout 85 86 the part of the reservoir considered. Since the temperature appears to be a key parameter for a safe 87 storage, the influence of the thermal properties (heat capacity and thermal conductivity) of the rocks 88 are also addressed. Although a tight reservoir is considered here, this study could also be valuable for EOR (Enhanced Oil Recovery) modelling (Ahmadi et al., 2016) or modelling of the combination of EOR 89 and CO₂ storage (Bender and Akin, 2017). 90

91

92 2. Materials and methods

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94 2.1. Modelling tools
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The Computational Fluid Dynamics modelling has been performed by COMSOL Multiphysics® (version
5.6). The different steps of the modelling are described below. A detailed kinetic model for *n*-octane
oxidation (Pacini-Petitjean et al., 2015b) has been used in order to derive accurate kinetic parameters

to be implemented in COMSOL Multiphysics. The computations of the kinetic model have been
performed by CHEMKIN II using the code SENKIN in a closed isochoric reactor (Kee et al., 1989).

100 2.2. Geological context and rock characteristics

101 The geological structure considered is a deep, isolated, faulted Jurassic horst, overlaid by a 4500 m 102 thick overburden series of turbiditic flysch deposits of Upper Cretaceous (Cenomanian) to Tertiary 103 (Eocene) age, and situated within the Pyrenean foredeep basin. The analysis of cored materials led to 104 the mineralogical composition of the cap-rocks (Table 1) and the reservoir (Table 2) (Girard et al., 105 2013). The composition of the reservoir varies because of the mineralogical heterogeneity of such a 106 fractured reservoir. It has been assumed that the injection of CO₂ does not modify the mineralogical 107 composition of the rocks, which remains constant throughout the study, as well as the porosity and 108 permeability of the rocks. The mineralogical and geochemical impact of gas impurity injection in the 109 same geological environment have been tested by Renard et al. (2014).

111 **Table 1**. Mineralogical composition of the cap-rocks.

Mineral	Mass fraction %
Pyrite	0.3
Ankerite	4.6
Fe Dolomite	0.0
Non-Fe Dolomite	2.1
Calcite	63.2
Quartz	10.5
Illite	8.3
Interstratified clays	6.5

113 **Table 2**. Mineralogical composition of the reservoir.

Mineral	Range mass fraction %	Average mass fraction %
Dolomite	84-98	93.5
Calcite	0-4	0.5
Quartz	1-4	2.5
Pyrite	0-2	0.4
Illite	1-5	2.5
Chlorite	0-1	0.3
Apatite	0-1	0.2
Insoluble Organic Matter	<<1	0.1

114

According to the mineralogical composition, Calcite and Quartz are the major minerals in the cap-rock and the reservoir is mainly made of Dolomite. The other minerals were neglected for the study. Table Sums up the thermal properties of these three minerals (heat capacity and thermal conductivity) (Birch and Clark, 1940; Krupka et al., 1985).

119

Table 3. Heat capacitiy (C_p) and thermal conductivity (K) of the major minerals of cap-rocks and
 reservoir (Birch and Clark, 1940; Krupka et al., 1985).

Mineral	Highest C _p	Т(К)	Lowest C _p	Т(К)	Highest K	Т(К)	Lowest K	Т(К)
	(J.kg ⁻¹ .K ⁻¹)		(J.kg ⁻¹ .K ⁻¹)		(W.m ⁻¹ .K ⁻¹)		(W.m ⁻¹ .K ⁻¹)	
Dolomite	1188	620	933	350	4.98	273	2.55	573
Quartz	1163	925	319	115	6.19	1400	1.34	280
Calcite	1289	1200	807	293	4	273	2.13	600

Since the dolomite is the main rock of the reservoir and since the heat capacity of dolomite does not depend a lot on temperature in the considered temperature range (Supplementary material), the heat capacity of the reservoir has been considered as constant and equal to 933 J.kg⁻¹.K⁻¹ (Table 3). For the cap-rocks, by considering only the main minerals, the mean heat capacity is from the same order of magnitude and does not change a lot with temperature as well. So, to simplify, the heat capacity of the base and cap-rocks has been taken constant and equal to 933 J.kg⁻¹.K⁻¹ too.

For the CFD modelling, the thermal conductivity has to be expressed as a function of temperature,since it depends a lot on temperature (Table 3). An equation as eq (1) is usually used.

131

132
$$K = \sum a_i \times T^{ni}$$
 eq (1)

133 Where T is the temperature (K) and a_i and n_i are coefficients, positive or negative.

134 The equations used in this study are given in the Supplementary Material.

135 The thermal conductivity of the reservoir and cap rocks was estimated by eq (2).

136
$$K (bulk) = \sum_{i=1}^{n} K_i M_i / M_{tot}$$

137

138 Where M_i , K_i are respectively the mass and thermal conductivity of component *i* of the rocks, and 139 M_{tot} is the total mass of the bulk material. According to the previous equations, the thermal 140 conductivity of the cap and base rocks varies between 3.35 and 4.95 W.m⁻¹.K⁻¹, and the thermal 141 conductivity of the reservoir rocks from 3.65 to 5.45 W.m⁻¹.K⁻¹, in the studied conditions, when the 142 temperature increases.

The part of the reservoir considered has been modelled by a cylinder (diameter: 800 m; thickness: 200
m), covered by 100 m of cap rock and lying on 100 m of base rock with the same properties as the caprock (Figure 1). The injection well is supposed to be in the center and perforated over the entire height

eq (2)

of the reservoir. It is surrounded by several injection wells in which the injection flow rates are the same. So, the drainage zone considered is surrounded by the drainage zones of the other wells and it seems separated from the other drainage zones by impermeable boundaries. These boundaries act as tight walls since the composition and temperature are the same on both sides and so no heat or mass transfer is possible. Therefore, the model is radial and axisymmetric and the cylinder has been considered as tight and thermally insulated on the sides.



152

153 Figure 1. Axisymmetric modelling of the geological pile considering homogeneous cap-rock and base-

rock of 100 m, and a heterogeneous part of a reservoir of 200 m of thickness.

155

156 2.3. Modelling of the reservoir fractured medium

157 The reservoir is considered as a fractured porous medium, *i.e.* a medium intersected by a network of

158 interconnected fractures. The fluid flows in the fractures as well as in the matrix blocks. These two

159 superposed continua coexist and interact with each other.

- 160 Another approach to model this system is to use effective parameters for porosity and permeability
- 161 (Table 4) without dividing the porous media into matrix blocks and fractures.

162

163 **Table 4**. Effective porosity and permeability of reservoir rocks (unpublished data).

Conduction properties	Effective
Porosity	3%
Permeability	1 mD

165 Since the fluid flows more rapidly in the fractures than in the matrix, it has been assumed that it does 166 not directly flow from one block to another. First it flows from a block into the fractures, and then into 167 another block or it remains in the fractures (Arbogast et al., 1990). Therefore, the equations that 168 describe the flow in the fracture continuum contain a source term that represents the flow of fluid 169 from the matrix to the fractures. This term is considered as being distributed over the entire medium. 170 Finally, we have assumed that the external sources and sinks interact only with the fracture system, 171 since the flow is much faster in this system than in the matrix blocks. Based on these assumptions, the mass conservation of the flow through the matrix and fractures in a fractured porous medium is given 172 173 by eq (3-4).

174

175
$$\frac{\partial(\rho \phi_m)}{\partial t} + \operatorname{div}(\rho \vec{v}_m) = q_{mf} \qquad \text{eq (3)}$$

176
$$\frac{\partial(\rho \phi_f)}{\partial t} + \operatorname{div}(\rho \vec{v}_f) = -q_{mf} + q_{ext} \qquad \text{eq (4)}$$

177 Where the subscripts *m* and *f* represent the matrix and fracture continua respectively, ϕ stands for 178 the porosity, \vec{v} for the velocity, ρ for the density, q_{mf} denotes the flow from the matrix to the fractures, 179 q_{ext} indicates the external sources and sinks. The velocities \vec{v}_m and \vec{v}_f are determined by Darcy's law. 180 Since in this study the effective parameters were used for the gas multicomponent system behavior in 181 fractured media, eq (5) and (6) have to be considered instead of eq (3) and (4).

183
$$\Phi_{ef} \frac{\partial(\rho C_i)}{\partial t} + \operatorname{div}(\rho C_i \vec{v}) = D^i - S^i \qquad \text{eq (5)}$$

184
$$\vec{v} = -\frac{k_{ef}}{\eta} (\overrightarrow{\text{grad } P} + \rho g)$$
 eq (6)

186 Where ϕ_{ef} , k_{ef} are the effective porosity and permeability respectively, η is the fluid dynamic viscosity, 187 P is the pressure, C_i is the molar concentration of species *i*, ρ is the gas density and D^i , S^i schematically 188 represent diffusion and source terms respectively. These terms are detailed below.

189 Furthermore, we assume that porosity does not change during the simulations, as explained below.

190

191 2.4. Governing equations

192

193 The equations that describe a reacting single-phase gaseous flow in a porous medium are detailed in 194 this sub-section. It is assumed that the gas is ideal.

195 Eq (7) corresponds to the energy equation.

196
$$C_p \frac{\partial T}{\partial t} + \vec{v} C_{pg} \rho_g grad(T) = div (K_T grad(T)) + Q \qquad \text{eq (7)}$$

198
$$C_p = \phi_{ef} C_{pg} \rho_g + (1 - \phi_{ef}) C_{ps} \rho_s$$
, $K_T = \phi_{ef} K_g + (1 - \phi_{ef}) K_s$

where C_{pg} and C_{ps} are specific heat capacity of gas and solid rock respectively, K_g and K_s are thermal conductivity of fluid and rocks, ρ_g is fluid density, ρ_s is rock density. Q stands for the volumetric power released by the fuel oxidation. It is given by eq (8).

$$202 Q = H \times W eq (8)$$

where *H* is the enthalpy of reaction and *W* the rate of the chemical reaction. These terms are detailed
below, in a specific subsection.

The mass conservation equation of the species *i* (*n*-octane, oxygen, nitrogen or products) describing the chemical reaction in single gas phase is written as eq (9).

208
$$\Phi_{ef} \frac{\partial(\rho_g C_i)}{\partial t} + \operatorname{div}(\rho_g C_i \vec{v}) = \Phi_{ef} \operatorname{div}(\rho_g \vec{J}_i) - \rho_g^0 W_i \qquad \text{eq (9)}$$

209

210 Where $\vec{J_1}$ is the molar flux of species *i*, ρ_g^0 -the gas density at initial conditions and W_i the rate of 211 production or consumption of *i* component with consideration of its mass-weighted stoichiometric 212 coefficient (Yucel Akkutlu and Yortsos, 2003).

The velocity of the flow is deduced from the Darcy's law (eq (6)), assuming that the fluid is mostly composed of nitrogen. Nitrogen viscosity is given as a function of temperature in the Supplementary material.

The gradient of pressure in the reservoir is described by the diffusion-type equation eq (10) (Monteiroet al., 2012).

218
$$\Phi_{ef} \frac{\partial \rho}{\partial t} = \operatorname{div}\left(\frac{k_{ef}}{\eta} \rho \, \overline{\operatorname{grad}} P\right)$$
 eq (10)

The system of the equations (7)-(10) should be completed by the equation of state for gas and porous media. The density of gas can be found as follows (eq (11)):

221
$$\rho_g = M_g \frac{P}{z(P)RT}, \qquad \text{eq (11)}$$

where *P* is the pressure, *T* is the temperature, *R* is the universal gas constant, z(P) is the gas compressibility, and M_g the average molecular weight of the gas. We assume that the equation of state for the gas phase is the ideal gas law, thus z=1. The porous medium is considered to be low deformable. In comparison with the gas compressibility, the compressibility of rocks can be neglected. The permeability of porous rocks is also assumed as constant. Moreover, the potential cooling of the injected gas by the Joule-Thomson effect has been considered as negligible since the pores are very small.

229 2.5. Kinetic modelling

It has been assumed that the hydrocarbons remaining in the depleted reservoir can be modelled considering a single hydrocarbon skeleton of an average chain-length (such as *n*-octane), and that the fluid does not contain any water, except the product of the oxidation.

All effects of the injection on the reservoir properties (composition, porosity, ...) and the exothermicity of the oxidation reactions of the reservoir rocks have been neglected, since there was no water initially contained in the reservoir and the production of water is rather low (Pacini-Petitjean et al., 2016). Indeed, all reactions of the rocks with the injected gas requiring water (Renard et al., 2014), they have not been considered in this study.

238 It has been previously demonstrated (Pacini-Petitjean et al. 2015a) that CO₂ does not react with the 239 hydrocarbons in the reservoir conditions. Therefore, CO₂ has been replaced by N₂ in this study, as in 240 the previous studies (Pacini-Petitjean et al. 2015a, 2015b and 2016), though their thermal and flow 241 properties could slightly modify the quantitative results. Although the molar fraction of O₂ remaining 242 after oxy-combustion is usually less than 6%, as for in the previous experimental studies, the air composition was used to simulate the injected gas. Indeed, Pacini-Petitjean (2015) compared the 243 244 experimental results of n-octane oxidation with 3% and 20% of O2, and got the same consumption 245 kinetics of *n*-octane, which is due to the large excess of O_2 in comparison to *n*-octane. So, to be 246 consistent with our previous studies, air has been considered for the injection as a model of flue gas.

In a previous study (Pacini-Petitjean et al., 2015b), a detailed kinetic model for *n*-octane oxidation was
built, which includes more than 3000 free-radical reactions. Due to the size of the model, it is not
possible to implement the full mechanism into COMSOL Multiphysics software. Therefore, a single

250 reaction has been implemented in the software, that is the total oxidation of *n*-octane (eq (12)), and a 251 very simple kinetic law was used in this study. Oxygen is present in large excess in comparison to n-252 octane, since the injection of gas is continuous until the limit value is reached and the molar fraction 253 of *n*-octane drops not only by oxidation, but also by dilution in the injected gas stream. In these 254 conditions, it has been considered that the kinetic law does not depend on oxygen concentration, but 255 only on *n*-octane concentration. It has been observed that a kinetic order of 2 fits *n*-octane 256 consumption satisfactorily (eq (13)) (Figure S1). Using the kinetic model in other conditions than a large 257 excess of O_2 compared to *n*-octane would lead to totally erroneous conclusions.

258
$$C_8 H_{18} + \frac{25}{2} O_2 \rightarrow 8 C O_2 + 9 H_2 O$$
 eq (12)

259
$$W = A \times e^{-E_a/_{RT}} \times C_{octane}^2 \qquad eq (13)$$

260 It must be noted that, according to a previous study (Pacini-Petitjean et al., 2016), the production of 261 water and CO_2 is rather slow and low, and that the intermediary products (ketones and carboxylic 262 acids) are predominant. That is why the potential effect of water and CO_2 on the phase state or the 263 fluid properties has not been considered in this study. It also means that eq (12) is a rough equation of 264 reaction, but what really matters in this study is the kinetics of *n*-octane consumption and the heat 265 release.

Concerning the kinetic parameters, the best match with CHEMKIN data has been obtained with A = $2.3 \times 10^6 \text{ m}^3 \text{.mol}^{-1} \text{.s}^{-1}$ and $\text{E}_a = 100 \text{ kJ/mol}$, where A is the frequency factor and E_a the activation energy (Figure S1).

The enthalpy of reaction is of the utmost importance in this study, in order to be able to estimate the temperature evolution in the reservoir. First, it has been assumed that the heat release is only related to the enthalpy of total *n*-octane oxidation (eq (12)), and that it is constant and equal to - 5112.6 kJ/mol (Turns, 2012, value at 298 K). The influence of temperature on the evolution of the enthalpy of reaction has been neglected. Nevertheless, the reaction is much more complex than a single reaction and the enthalpy of reaction varies according to the evolution of main reaction pathways. Therefore, the
volumetric power released (eq (8)) was directly computed as a function of *n*-octane concentration,
using CHEMKIN. Then, the curve obtained was fitted by a polynomial function in order to be used
directly in COMSOL Multiphysics (Supplementary material).

278

279 2.6. Initial and boundary conditions

In the model, the initial temperature of the top of the reservoir was set to 423.15 K, which corresponds to a rather hot and deep reservoir, and the initial Bottom Hole Pressure (BHP) to 30 bar, characteristic of a depleted reservoir. The temperature changes with depth are considered by a geothermal gradient of 30°C/km. The initial molar fraction of *n*-octane was taken equal to 0.015 in the entire reservoir, the rest being nitrogen.

In the well, the BHP was set to 100 bar and maintained constant during the entire injection. That means that the mass flow rate of the injected gas decreases during the process. The molar fractions of oxygen and nitrogen were taken equal to 0.22 and 0.78 respectively, as in our previous studies (Pacini-Petitjean et al. 2015a, 2015b and 2016). It has been considered that the injection lasts until the pressure reaches 100 bar in the entire reservoir.

The boundary conditions set in this study correspond to the Neumann conditions, which means that there is no heat flux or fluid flow across the boundaries (eq (14)).

292
$$\begin{cases} \vec{n} \cdot (K_{ef} \vec{\nabla} T) = 0 \\ \vec{n} \cdot (D_i \vec{\nabla} C_i) = 0 \end{cases}$$
 eq (14)

293 Where \vec{n} is a vector normal to the wall.

The thermal reservoir models account for heat exchange with the surroundings, typically through heat conduction to the cap and base rocks. The insulated walls are shown in Figure 2. As explained before, the sides are considered as insulated walls, because the drainage zone is surrounded by other drainage 297 zones leading to symmetrical temperature profiles and so to the same temperature on both sides. On

the other hand, it has been assumed that the fluid does not penetrate into the cap or the base rocks.



300

Figure 2. Walls in blue are thermal insulated, and walls in red prevent the fluid flux. Dashed green line
 corresponds to the well location.

- 303 3. Results and discussion
- 304

305 *3.1. Simulations at constant enthalpy of reaction*

306 The temperature and *n*-octane concentration in the entire reservoir have been computed up to about 307 10 years. For an example, Figure 3 shows the results after 1500 days (about 4 years) along a cut-line. 308 Near the well, in the first 20 m, almost all *n*-octane is consumed and the consumption of *n*-octane has 309 begun in the entire reservoir. The temperature increases in less than half of the reservoir's radius, but 310 the increase is less than 10 K. The maximum temperature is reached after 2500 days (almost 7 years -311 end of injection) and it is equal to 443 K. The temperature profile is represented in Figure 4. The 312 maximum temperature is approximately located at 200 m of the injection well (radial position) and 313 half height. The temperature of the cap and base rocks have also increased by thermal conduction. 314 After 2500 days, the pressure reaches 100 bar and so no more gas is injected. It should be noted that 315 the auto-ignition temperature of *n*-octane is 493 K (De Soete and Feugier, 1976), which means that 316 there is apparently no auto-ignition in the reservoir in the studied conditions.



Figure 3. Temperature (a) and *n*-octane conversion (b) as a function of the distance from injection well, after 1500 days (cross-sections) – constant enthalpy of reaction (end of injection: 2500 days – $C_{p \text{ rocks}} =$ 933 J.kg⁻¹.K⁻¹).



Figure 4. Temperature profile in the reservoir (between -100 and -300 m), in the cap rock (between 0 and -100 m) and in the base rock (between -300 and -400 m), after 2500 days (end of injection – $C_{p \text{ rocks}}$ 331 = 933 J.kg⁻¹.K⁻¹).

332 3.2. Simulations with the heat release estimated from CHEMKIN results

16

333 In this subsection, the reservoir temperature has been computed using the volumetric power release 334 estimated from CHEMKIN as a function of remaining *n*-octane concentration. The temperature 335 increases much faster than in the previous computation (Figures 5a and b) and reaches 533 K in the 336 center of the reservoir after only 60 days, which is above the auto-ignition temperature. Figure 6 shows 337 the temperature evolution as a function of time and radial position in a cross section and so the 338 reaction wave can be observed. The heat transfer by the rocks is much slower than by the fluid, so the 339 cap and base rocks warm up very slowly. According to Figure 7a, n-octane is reduced to half of the 340 initial value of the reservoir after 50 days. Figure 7b shows the progression of oxygen into the reservoir. 341 After the total consumption of *n*-octane, the dissipation of heat in the base and cap rocks can be 342 observed (Kihm et al., 2019) and the temperature slowly decreases and is almost uniform in the 343 reservoir and equal to 480 K after 500 days (Figure 5c).

344





Figure 5. a) Average temperature in the reservoir as a function of time until 50 days; b) Temperature
 profile in the reservoir after 50 days; c) Temperature profile in the reservoir after 500 days (end of
 injection: 300 days- C_{p rocks} = 933 J.kg⁻¹.K⁻¹).



Figure 6. Temperature as a function of the distance from injection well (cross-sections after 10, 20, 30,
40 and 50 days) – heat release from CHEMKIN results (end of injection: 300 days– C_{p rocks} = 933 J.kg⁻¹.K⁻

¹).



Figure 7. Molar fractions (gas phase) of the reactants (a, *n*-octane and b, O_2) as a function of the distance from injection well (cross-sections after 10, 20, 30, 40 and 50 days - end of injection: 300 days- $C_{p \text{ rocks}} = 933 \text{ J.kg}^{-1}$.K⁻¹).

Figure 8a shows the pressure profile after 50 days. Pressure reaches 80 bar near the injection well. The average pressure reaches 100 bar after almost one year (300 days - Figure 8b), which corresponds to the end of the gas injection into the well in the simulations. Indeed, the temperature increase is greater and faster with the heat release estimated from CHEMKIN results than with a constant enthalpy of reaction, and so the pressure limit is reached sooner than by the first simulations. In the meantime, the oxygen concentration slowly stabilizes.



Figure 8. a) Pressure (bar) profile in the reservoir 50 days after the beginning of the injection; b) Average pressure in the entire reservoir as a function of time (end of injection: 300 days– $C_{p \text{ rocks}} = 933$ J.kg⁻¹.K⁻¹).

According to this study, which considers the heat released in a more realistic way than only the enthalpy of total oxidation, the gas injection into the reservoir would strongly affect the hydrocarbon composition and the temperature over a few months after the beginning of the injection. Nevertheless, the reactivity of *n*-octane against oxidation was estimated in homogeneous phase. The reactivity in a porous medium could be lowered by the recombination of the radicals onto the walls (Garo et al. 1984). The heat release and the reaction rate would be lower than estimated here. Moreover, the temperature increase depends on the competition between the kinetics of oxidation and the thermal properties of the rocks, which is studied in the next sub-section.

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3.3. Influence of thermal properties anisotropy on heat dissipation into the reservoir

386 The evolution of temperature in the reservoir is governed by the competition between the 387 exothermicity of reaction, the dynamics of the gas injected and the thermal properties of the reservoir, 388 the cap and the base rocks. In order to assess the influence of anisotropy in thermal conductivity on 389 heat dissipation, first, a horizontal thermal conductivity 10 times higher than the vertical one was set 390 in the simulations (Figure 9). The opposite situation was also considered (Figure 9). The maximum 391 temperature is reached after 80 days and is about 540 K in both cases. No major difference is observed, 392 but the slope after the maximum temperature point in the case of a higher vertical conductivity than 393 the horizontal one. It could be explained by the higher vertical conductivity, which leads to more 394 thermal diffusion into the cap and base rocks. Nevertheless, the thermal conductivity of the rocks has 395 only a small influence in comparison with the dynamics of the gas injected and the heat released by 396 the oxidation.



Figure 9. Temperature as a function of time in case of anisotropy (red circle markers: horizontal thermal conductivity 10 times higher than the vertical one; blue square markers: opposite situation end of injection: 300 days- $C_{p \text{ rocks}} = 933 \text{ J.kg}^{-1}$.

Secondly, the influence of the heat capacity of the reservoir rock was considered. The heat capacity of usual reservoir rocks ranges from 500 to 1500 J.kg⁻¹.K⁻¹, depending on mineral composition. These two limit cases were successively investigated. First, the heat capacity of the reservoir, cap and base rocks was set to 500 J.kg⁻¹.K⁻¹. In this case, the rocks warm up very quickly (Figure 10). The reservoir temperature reaches its maximum (about 570 K) after 40 days, and then the temperature slightly decreases. After 100 days, the temperature in about three quarters of the reservoir is greater than 530 K.



Figure 10. Temperature as a function of time and distance from the injection well (heat capacity of the
reservoir, cap and base rocks set to 500 J.kg⁻¹.K⁻¹- end of injection: 300 days).

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414 Secondly, the heat capacity was set to 1500 J.kg⁻¹.K⁻¹ (Figure 11). In this situation, since the rocks absorb 415 most of the energy released by the exothermicity of the reactions, the temperature only increases by 416 3 K after 100 days, which is completely negligible. Figure S2 compares the evolution of temperature 417 computed for the three heat capacities tested in order to better highlight their influence. As expected, 418 the consumption of the reactants is all the faster than the heat capacity is lower and the temperature 419 higher, since the kinetics is temperature dependent (Figure 12). Figures S3 and S4 compare the 420 evolution of respectively *n*-octane and O_2 molar fractions computed for the three heat capacities. The 421 difference between the results obtained with 500 and 933 J.kg⁻¹.K⁻¹ is more pronounced than between 422 933 and 1500 J.kg⁻¹.K⁻¹ since the increase of temperature with 500 J.kg⁻¹.K⁻¹ is very important. It appears that there is no major difference between 933 and 1500 J.kg⁻¹.K⁻¹ except at the very beginning of the 423 424 injection. In the case of O₂, since it is progressively injected into the reservoir, the concentration increases faster when the heat capacity is high. As O2 is less consumed at low temperature, which 425 426 corresponds to a high heat capacity, it penetrates the reservoir faster.

Contrary to the thermal conductivity, the heat capacity has a major influence on temperature in the reservoir. In order to choose the location of the storage site, which ensures the maximum thermal stability of the reservoir, it is particularly important to pay attention to the heat capacity of the reservoir, base and cap rocks and to choose rocks with the highest heat capacity. Furthermore, for a given depleted reservoir, it is important to adapt to the thermal characteristics of the reservoir and to the remaining hydrocarbons, the composition of the injected gas, its flow rate, and the temperature and the pressure in the well, in order to limit the risk.

435



437 **Figure 11.** Temperature as a function of time and distance from the injection well (heat capacity of the

438 reservoir, cap and base rocks set to 1500 J.kg⁻¹.K⁻¹ - end of injection: 300 days).

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Figure 12. Reactant molar fractions as a function of time and distance from the injection well: a) *n*octane (reactant present in the reservoir) and heat capacity of the reservoir, cap and base rocks set to 500 J.kg⁻¹.K⁻¹; b) *n*-octane and heat capacity set to 1 500 J.kg⁻¹.K⁻¹; c) O₂ (injected into the reservoir) and heat capacity set to 500 J.kg⁻¹.K⁻¹; d) O₂ (injected into the reservoir) and heat capacity set to 1 500 J.kg⁻¹.K⁻¹ 1 .K⁻¹ (end of injection: 300 days).

448 **4.** Conclusions

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The paper aimed at studying the influence of oxygen contained in the flue gas, for an example captured after oxy-combustion, on the oxidation kinetics of the residual fuel in the depleted reservoir and its influence on some of the reservoir parameters or properties. More precisely, the paper focused on the influence of the oxidation exothermicity and the thermal properties of the rocks on the pressure and the temperature evolution, and on the kinetics of the reaction inside a depleted oil reservoir.

Despite some assumptions and a quite high initial concentration of O₂, it clearly appears that the system is not inert with respect to the hydrocarbons reactivity. The residual oil is probably consumed in less than one year in relatively small reservoirs, leading to the temperature increase throughout the reservoir. After the end of the reaction, the temperature decreases slowly by heat dissipation in the cap and base rocks.

460 The extent of increase in the reservoir temperature is clearly a function of the thermal properties of 461 the rocks. It has been demonstrated that the thermal conductivity of the rocks and their potential 462 anisotropy have only a small influence on the temperature evolution. Nevertheless, the temperature 463 evolution is highly dependent on the heat capacity and quite contrasting behaviors can be observed 464 depending on low or high heat capacities of the rocks. A high heat capacity of the rock leads to a 465 negligible increase in the temperature of the reservoir. Therefore, in order to limit the temperature 466 increase and the kinetics of the reaction, it appears important to select reservoirs whose rocks as well 467 as the rocks of the cap and the base have a high heat capacity (greater than 1000 J.kg⁻¹.K⁻¹), especially 468 when O₂ concentration in the injected gas is high (at least greater than 3%). In this case, the 469 temperature increase seems negligible and the oil oxidation is very slow, which appears to be an 470 essential condition for a safe injection operation. Consequently, for a depleted reservoir with rocks of 471 low heat capacity, it is important to decrease the O_2 concentration in the injected gas, its flow rate,

- and the temperature and the pressure in the well, in order to limit the risk of significantly rapid rise in
- 473 the reservoir temperature and its associated thermal instability issues.
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