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1 **Flue gas injection into depleted tight hydrocarbon reservoirs in the context of global**
2 **warming mitigation: computed evolution of some reservoir parameters**

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14

15 **Keywords:** depleted oil reservoir, CO₂ storage, temperature profile, oxy-combustion, CFD

16

17 **Abstract:**

18 CO₂ storage in depleted tight hydrocarbon reservoir seems to be a promising solution to mitigate
19 greenhouse gas emissions. However, flue gas contains not only CO₂ but also minor gaseous impurities
20 due to CO₂ production or capture processes. In the case of oxy-combustion, the main impurity can be
21 O₂ in concentration up to 7 %. O₂ injection into the reservoir can lead to the oxidation of the residual
22 hydrocarbons and therefore, it is necessary to assess the thermal consequences on the reservoir.

23 COMSOL Multiphysics® has been used to model an axisymmetric fractured porous reservoir with its
24 cap and base rocks and containing *n*-octane as a model compound of residual oil. A global kinetic model
25 for *n*-octane oxidation has been derived from a previous detailed free-radical model, and implemented
26 into the reservoir model. Simulated injections of N₂/O₂ mixture (representing a simplified flue gas)
27 have been performed to compute temperature, pressure and *n*-octane concentration profiles. The
28 results show that the oxidation exothermicity may have a strong influence on temperature profile,
29 especially when the heat capacity of the rock is rather low. In these conditions, the remaining
30 hydrocarbons may be consumed in some months. However, the influence of the thermal conductivity
31 seems negligible. Therefore, it appears safer to select reservoirs whose rocks compositions, in
32 particular the cap and base rocks, have a high heat capacity to promote heat dissipation. It is a criterion
33 to consider when selecting a storage site containing residual hydrocarbons, especially for CO₂ captured
34 after oxy-combustion.

35

36 **1. Introduction**

37 The CO₂ emissions from combustion of fossil fuel is still increasing and reached 33.62 Giga ton/year in
38 2019 (IEA, 2022), leading to dramatic climate warming due to CO₂ greenhouse gas effect (IPCC, 2021).

39 The mitigation of such alarmingly increasing CO₂ emissions mainly considers three combined
40 strategies: 1) switching to a low-carbon economy which relies on renewable and alternative energy
41 sources; 2) increasing the efficiency of the current technologies and decreasing the energy loss; 3)
42 using CCS (Carbon Capture and Storage) technologies to reduce CO₂ emissions, since the worldwide
43 dependency on fossil fuels is probably going to last several more decades (Maroto-Valer, 2010). CCS
44 appears to be an efficient strategy to mitigate anthropogenic CO₂ emissions from large sources, by up
45 to 90% (Maroto-Valer, 2010). CCS mainly consists of: 1) capturing and potentially separating CO₂ from
46 the other exhaust gases; 2) carrying CO₂ to its storage location, after compression to supercritical
47 conditions; 3) storing CO₂, e.g. in several types of geological storage sites (Liu et al., 2018).

48 The geological reservoirs that are the most frequently considered are the deep saline aquifers. They
49 may be subjected to regional pressure build-up due to water movement, which may affect the
50 reservoir integrity (Hamza et al. 2021) and the storage capacities. Injecting CO₂ into coal seam
51 reservoirs is also considered, potentially for EMR (Enhanced Methane Recovery) purpose (Zhang and
52 Ranjith, 2019). Another type of geological reservoir that is the focus of attention, is the depleted
53 hydrocarbon (gas or oil) reservoirs. Their storage capacity is much higher than that of coal seams (Kaldi
54 et al., 2009) and they present several advantages over saline aquifers. Indeed, the oil and gas fields
55 have proved their potential containment by the retention of hydrocarbons for millions of years. They
56 have been widely studied (storage capacity, structural and stratigraphic patterns...) (Kaldi et al., 2009).
57 Another advantage is the limited pressure inside the reservoirs due to the former hydrocarbon
58 production, which allows injecting CO₂ at moderate pressure, which is technically and economically
59 favorable.

60 Power plants and other industrial processes that use combustion usually produce flue gases at
61 approximately atmospheric pressure, with CO₂ diluted mostly in N₂. CO₂ must then be separated from
62 other components in the flue gas (post-combustion capture). Another solution is pre-combustion
63 which consists of the removal of carbon from the fuel prior to combustion. At last, oxy-combustion,
64 i.e. combustion in purified O₂ instead of air, is another way to solve this issue. Oxy-combustion enables
65 a much easier capture of CO₂ by the flue gas condensation than post-combustion in which N₂ is the
66 dominant component (Tumsa et al., 2017). Nevertheless, flue gas from oxy-combustion may contain a
67 relatively high level of impurities, such as water, Ar (0-5%), N₂ (0-15%), O₂ (0-7%) and SO₂ (0-1.5%)
68 (Mikunda, 2012).

69 This study focuses on CO₂ captured after oxy-combustion processes, which contains residual O₂. It aims
70 to assess some effects of the presence of O₂ in the injected gas on some parameters of the gas/oil
71 depleted reservoir and to assess if it is necessary to remove O₂ before injection since purification is
72 expensive. Previous studies already showed that the impurities reduce the structural trapping capacity

73 for CO₂ by reducing the CO₂ density (Wang et al., 2012). Another consequence of the presence of O₂ is
74 examined here. Indeed our previous studies (Pacini-Petitjean et al., 2015a, 2015b, 2016) showed that
75 the oxidation of the hydrocarbons occurs in the temperature conditions of relatively deep hydrocarbon
76 reservoirs ($T > 373$ K). Therefore, the oxidation kinetics of the residual hydrocarbons may impact on
77 the pressure and temperature evolution in the reservoir during the injection and storage. That is why
78 a model of one part of a fractured porous reservoir has been setup by COMSOL Multiphysics® (version
79 5.6), which considers heat and mass transfers, flow and kinetics. Several assumptions have been made,
80 in particular the flue gas has been replaced by air to emphasize the oxidation phenomenon and to stay
81 in the same conditions as in our previous studies. The other impurities have not been considered. This
82 simplified representation of a flue gas appears appropriate since the study mainly concerns the
83 thermal effects of O₂. An injection well in the center and perforated over the entire height of the
84 reservoir has been considered. The model aims at computing the consumption of O₂ and the
85 hydrocarbons, the pressure increase and the temperature evolution as a function of time, throughout
86 the part of the reservoir considered. Since the temperature appears to be a key parameter for a safe
87 storage, the influence of the thermal properties (heat capacity and thermal conductivity) of the rocks
88 are also addressed. Although a tight reservoir is considered here, this study could also be valuable for
89 EOR (Enhanced Oil Recovery) modelling (Ahmadi et al., 2016) or modelling of the combination of EOR
90 and CO₂ storage (Bender and Akin, 2017).

91

92 **2. Materials and methods**

93

94 *2.1. Modelling tools*

95 The Computational Fluid Dynamics modelling has been performed by COMSOL Multiphysics® (version
96 5.6). The different steps of the modelling are described below. A detailed kinetic model for *n*-octane
97 oxidation (Pacini-Petitjean et al., 2015b) has been used in order to derive accurate kinetic parameters

98 to be implemented in COMSOL Multiphysics. The computations of the kinetic model have been
99 performed by CHEMKIN II using the code SENKIN in a closed isochoric reactor (Kee et al., 1989).

100 2.2. Geological context and rock characteristics

101 The geological structure considered is a deep, isolated, faulted Jurassic horst, overlaid by a 4500 m
102 thick overburden series of turbiditic flysch deposits of Upper Cretaceous (Cenomanian) to Tertiary
103 (Eocene) age, and situated within the Pyrenean foredeep basin. The analysis of cored materials led to
104 the mineralogical composition of the cap-rocks (Table 1) and the reservoir (Table 2) (Girard et al.,
105 2013). The composition of the reservoir varies because of the mineralogical heterogeneity of such a
106 fractured reservoir. It has been assumed that the injection of CO₂ does not modify the mineralogical
107 composition of the rocks, which remains constant throughout the study, as well as the porosity and
108 permeability of the rocks. The mineralogical and geochemical impact of gas impurity injection in the
109 same geological environment have been tested by Renard et al. (2014).

110

111 **Table 1.** Mineralogical composition of the cap-rocks.

Mineral	Mass fraction %
Pyrite	0.3
Ankerite	4.6
Fe Dolomite	0.0
Non-Fe Dolomite	2.1
Calcite	63.2
Quartz	10.5
Illite	8.3
Interstratified clays	6.5

112

113 **Table 2.** Mineralogical composition of the reservoir.

Mineral	Range mass fraction %	Average mass fraction %
Dolomite	84-98	93.5
Calcite	0-4	0.5
Quartz	1-4	2.5
Pyrite	0-2	0.4
Illite	1-5	2.5
Chlorite	0-1	0.3
Apatite	0-1	0.2
Insoluble Organic Matter	<<1	0.1

114

115 According to the mineralogical composition, Calcite and Quartz are the major minerals in the cap-rock
 116 and the reservoir is mainly made of Dolomite. The other minerals were neglected for the study. Table
 117 3 sums up the thermal properties of these three minerals (heat capacity and thermal conductivity)
 118 (Birch and Clark, 1940; Krupka et al., 1985).

119

120 **Table 3.** Heat capacity (C_p) and thermal conductivity (K) of the major minerals of cap-rocks and
 121 reservoir (Birch and Clark, 1940; Krupka et al., 1985).

Mineral	Highest C_p ($J.kg^{-1}.K^{-1}$)	T(K)	Lowest C_p ($J.kg^{-1}.K^{-1}$)	T(K)	Highest K ($W.m^{-1}.K^{-1}$)	T(K)	Lowest K ($W.m^{-1}.K^{-1}$)	T(K)
Dolomite	1188	620	933	350	4.98	273	2.55	573
Quartz	1163	925	319	115	6.19	1400	1.34	280
Calcite	1289	1200	807	293	4	273	2.13	600

122

123 Since the dolomite is the main rock of the reservoir and since the heat capacity of dolomite does not
124 depend a lot on temperature in the considered temperature range (Supplementary material), the heat
125 capacity of the reservoir has been considered as constant and equal to $933 \text{ J.kg}^{-1}.\text{K}^{-1}$ (Table 3). For the
126 cap-rocks, by considering only the main minerals, the mean heat capacity is from the same order of
127 magnitude and does not change a lot with temperature as well. So, to simplify, the heat capacity of
128 the base and cap-rocks has been taken constant and equal to $933 \text{ J.kg}^{-1}.\text{K}^{-1}$ too.

129 For the CFD modelling, the thermal conductivity has to be expressed as a function of temperature,
130 since it depends a lot on temperature (Table 3). An equation as eq (1) is usually used.

131

$$132 \quad K = \sum a_i \times T^{n_i} \quad \text{eq (1)}$$

133 Where T is the temperature (K) and a_i and n_i are coefficients, positive or negative.

134 The equations used in this study are given in the Supplementary Material.

135 The thermal conductivity of the reservoir and cap rocks was estimated by eq (2).

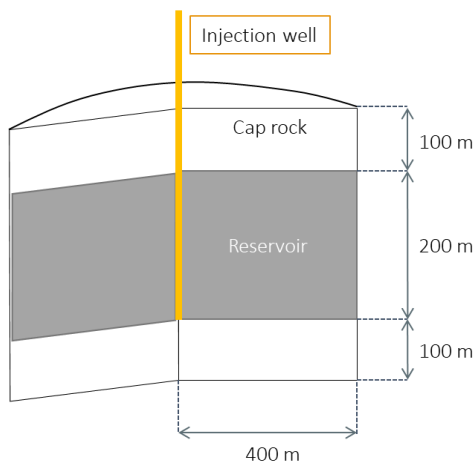
$$136 \quad K (\text{bulk}) = \sum_{i=1}^n K_i M_i / M_{tot} \quad \text{eq (2)}$$

137

138 Where M_i , K_i are respectively the mass and thermal conductivity of component i of the rocks, and
139 M_{tot} is the total mass of the bulk material. According to the previous equations, the thermal
140 conductivity of the cap and base rocks varies between 3.35 and $4.95 \text{ W.m}^{-1}.\text{K}^{-1}$, and the thermal
141 conductivity of the reservoir rocks from 3.65 to $5.45 \text{ W.m}^{-1}.\text{K}^{-1}$, in the studied conditions, when the
142 temperature increases.

143 The part of the reservoir considered has been modelled by a cylinder (diameter: 800 m ; thickness: 200
144 m), covered by 100 m of cap rock and lying on 100 m of base rock with the same properties as the cap-
145 rock (Figure 1). The injection well is supposed to be in the center and perforated over the entire height

146 of the reservoir. It is surrounded by several injection wells in which the injection flow rates are the
147 same. So, the drainage zone considered is surrounded by the drainage zones of the other wells and it
148 seems separated from the other drainage zones by impermeable boundaries. These boundaries act as
149 tight walls since the composition and temperature are the same on both sides and so no heat or mass
150 transfer is possible. Therefore, the model is radial and axisymmetric and the cylinder has been
151 considered as tight and thermally insulated on the sides.



152

153 **Figure 1.** Axisymmetric modelling of the geological pile considering homogeneous cap-rock and base-
154 rock of 100 m, and a heterogeneous part of a reservoir of 200 m of thickness.

155

156 *2.3. Modelling of the reservoir fractured medium*

157 The reservoir is considered as a fractured porous medium, *i.e.* a medium intersected by a network of
158 interconnected fractures. The fluid flows in the fractures as well as in the matrix blocks. These two
159 superposed continua coexist and interact with each other.

160 Another approach to model this system is to use effective parameters for porosity and permeability
161 (Table 4) without dividing the porous media into matrix blocks and fractures.

162

163 **Table 4.** Effective porosity and permeability of reservoir rocks (unpublished data).

Conduction properties	Effective
Porosity	3%
Permeability	1 mD

164

165 Since the fluid flows more rapidly in the fractures than in the matrix, it has been assumed that it does
166 not directly flow from one block to another. First it flows from a block into the fractures, and then into
167 another block or it remains in the fractures (Arbogast et al., 1990). Therefore, the equations that
168 describe the flow in the fracture continuum contain a source term that represents the flow of fluid
169 from the matrix to the fractures. This term is considered as being distributed over the entire medium.
170 Finally, we have assumed that the external sources and sinks interact only with the fracture system,
171 since the flow is much faster in this system than in the matrix blocks. Based on these assumptions, the
172 mass conservation of the flow through the matrix and fractures in a fractured porous medium is given
173 by eq (3-4).

174

$$175 \quad \frac{\partial(\rho \phi_m)}{\partial t} + \text{div}(\rho \vec{v}_m) = q_{mf} \quad \text{eq (3)}$$

$$176 \quad \frac{\partial(\rho \phi_f)}{\partial t} + \text{div}(\rho \vec{v}_f) = -q_{mf} + q_{ext} \quad \text{eq (4)}$$

177 Where the subscripts m and f represent the matrix and fracture continua respectively, ϕ stands for
178 the porosity, \vec{v} for the velocity, ρ for the density, q_{mf} denotes the flow from the matrix to the fractures,
179 q_{ext} indicates the external sources and sinks. The velocities \vec{v}_m and \vec{v}_f are determined by Darcy's law.
180 Since in this study the effective parameters were used for the gas multicomponent system behavior in
181 fractured media, eq (5) and (6) have to be considered instead of eq (3) and (4).

182

183
$$\Phi_{ef} \frac{\partial(\rho C_i)}{\partial t} + \text{div}(\rho C_i \vec{v}) = D^i - S^i \quad \text{eq (5)}$$

184
$$\vec{v} = -\frac{k_{ef}}{\eta} (\text{grad } P + \rho g) \quad \text{eq (6)}$$

185

186 Where Φ_{ef} , k_{ef} are the effective porosity and permeability respectively, η is the fluid dynamic viscosity,
 187 P is the pressure, C_i is the molar concentration of species i , ρ is the gas density and D^i , S^i schematically
 188 represent diffusion and source terms respectively. These terms are detailed below.

189 Furthermore, we assume that porosity does not change during the simulations, as explained below.

190

191 *2.4. Governing equations*

192

193 The equations that describe a reacting single-phase gaseous flow in a porous medium are detailed in
 194 this sub-section. It is assumed that the gas is ideal.

195 Eq (7) corresponds to the energy equation.

196
$$C_p \frac{\partial T}{\partial t} + \vec{v} C_{pg} \rho_g \text{grad}(T) = \text{div}(K_T \text{grad}(T)) + Q \quad \text{eq (7)}$$

197 with

198
$$C_p = \Phi_{ef} C_{pg} \rho_g + (1 - \Phi_{ef}) C_{ps} \rho_s, \quad K_T = \Phi_{ef} K_g + (1 - \Phi_{ef}) K_s$$

199 where C_{pg} and C_{ps} are specific heat capacity of gas and solid rock respectively, K_g and K_s are thermal
 200 conductivity of fluid and rocks, ρ_g is fluid density, ρ_s is rock density. Q stands for the volumetric power
 201 released by the fuel oxidation. It is given by eq (8).

202
$$Q = H \times W \quad \text{eq (8)}$$

203

204 where H is the enthalpy of reaction and W the rate of the chemical reaction. These terms are detailed
205 below, in a specific subsection.

206 The mass conservation equation of the species i (n -octane, oxygen, nitrogen or products) describing
207 the chemical reaction in single gas phase is written as eq (9).

$$208 \quad \Phi_{ef} \frac{\partial(\rho_g C_i)}{\partial t} + \text{div}(\rho_g C_i \vec{v}) = \Phi_{ef} \text{div}(\rho_g \vec{J}_i) - \rho_g^0 W_i \quad \text{eq (9)}$$

209

210 Where \vec{J}_i is the molar flux of species i , ρ_g^0 —the gas density at initial conditions and W_i the rate of
211 production or consumption of i component with consideration of its mass-weighted stoichiometric
212 coefficient (Yucel Akkutlu and Yortsos, 2003).

213 The velocity of the flow is deduced from the Darcy's law (eq (6)), assuming that the fluid is mostly
214 composed of nitrogen. Nitrogen viscosity is given as a function of temperature in the Supplementary
215 material.

216 The gradient of pressure in the reservoir is described by the diffusion-type equation eq (10) (Monteiro
217 et al., 2012).

$$218 \quad \Phi_{ef} \frac{\partial \rho}{\partial t} = \text{div} \left(\frac{k_{ef}}{\eta} \rho \overrightarrow{\text{grad}P} \right) \quad \text{eq (10)}$$

219 The system of the equations (7)-(10) should be completed by the equation of state for gas and
220 porous media. The density of gas can be found as follows (eq (11)):

$$221 \quad \rho_g = M_g \frac{P}{z(P)RT}, \quad \text{eq (11)}$$

222 where P is the pressure, T is the temperature, R is the universal gas constant, $z(P)$ is the gas
223 compressibility, and M_g the average molecular weight of the gas. We assume that the equation of
224 state for the gas phase is the ideal gas law, thus $z=1$.

225 The porous medium is considered to be low deformable. In comparison with the gas compressibility,
226 the compressibility of rocks can be neglected. The permeability of porous rocks is also assumed as
227 constant. Moreover, the potential cooling of the injected gas by the Joule-Thomson effect has been
228 considered as negligible since the pores are very small.

229 *2.5. Kinetic modelling*

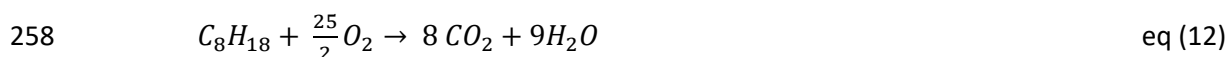
230 It has been assumed that the hydrocarbons remaining in the depleted reservoir can be modelled
231 considering a single hydrocarbon skeleton of an average chain-length (such as *n*-octane), and that the
232 fluid does not contain any water, except the product of the oxidation.

233 All effects of the injection on the reservoir properties (composition, porosity, ...) and the exothermicity
234 of the oxidation reactions of the reservoir rocks have been neglected, since there was no water initially
235 contained in the reservoir and the production of water is rather low (Pacini-Petitjean et al., 2016).
236 Indeed, all reactions of the rocks with the injected gas requiring water (Renard et al., 2014), they have
237 not been considered in this study.

238 It has been previously demonstrated (Pacini-Petitjean et al. 2015a) that CO₂ does not react with the
239 hydrocarbons in the reservoir conditions. Therefore, CO₂ has been replaced by N₂ in this study, as in
240 the previous studies (Pacini-Petitjean et al. 2015a, 2015b and 2016), though their thermal and flow
241 properties could slightly modify the quantitative results. Although the molar fraction of O₂ remaining
242 after oxy-combustion is usually less than 6%, as for in the previous experimental studies, the air
243 composition was used to simulate the injected gas. Indeed, Pacini-Petitjean (2015) compared the
244 experimental results of *n*-octane oxidation with 3% and 20% of O₂, and got the same consumption
245 kinetics of *n*-octane, which is due to the large excess of O₂ in comparison to *n*-octane. So, to be
246 consistent with our previous studies, air has been considered for the injection as a model of flue gas.

247 In a previous study (Pacini-Petitjean et al., 2015b), a detailed kinetic model for *n*-octane oxidation was
248 built, which includes more than 3000 free-radical reactions. Due to the size of the model, it is not
249 possible to implement the full mechanism into COMSOL Multiphysics software. Therefore, a single

250 reaction has been implemented in the software, that is the total oxidation of *n*-octane (eq (12)), and a
251 very simple kinetic law was used in this study. Oxygen is present in large excess in comparison to *n*-
252 octane, since the injection of gas is continuous until the limit value is reached and the molar fraction
253 of *n*-octane drops not only by oxidation, but also by dilution in the injected gas stream. In these
254 conditions, it has been considered that the kinetic law does not depend on oxygen concentration, but
255 only on *n*-octane concentration. It has been observed that a kinetic order of 2 fits *n*-octane
256 consumption satisfactorily (eq (13)) (Figure S1). Using the kinetic model in other conditions than a large
257 excess of O₂ compared to *n*-octane would lead to totally erroneous conclusions.



259
$$W = A \times e^{-E_a/RT} \times C_{octane}^2$$
 eq (13)

260 It must be noted that, according to a previous study (Pacini-Petitjean et al., 2016), the production of
261 water and CO₂ is rather slow and low, and that the intermediary products (ketones and carboxylic
262 acids) are predominant. That is why the potential effect of water and CO₂ on the phase state or the
263 fluid properties has not been considered in this study. It also means that eq (12) is a rough equation of
264 reaction, but what really matters in this study is the kinetics of *n*-octane consumption and the heat
265 release.

266 Concerning the kinetic parameters, the best match with CHEMKIN data has been obtained with A =
267 $2.3 \times 10^6 \text{ m}^3 \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$ and $E_a = 100 \text{ kJ/mol}$, where A is the frequency factor and E_a the activation energy
268 (Figure S1).

269 The enthalpy of reaction is of the utmost importance in this study, in order to be able to estimate the
270 temperature evolution in the reservoir. First, it has been assumed that the heat release is only related
271 to the enthalpy of total *n*-octane oxidation (eq (12)), and that it is constant and equal to - 5112.6 kJ/mol
272 (Turns, 2012, value at 298 K). The influence of temperature on the evolution of the enthalpy of reaction
273 has been neglected. Nevertheless, the reaction is much more complex than a single reaction and the

274 enthalpy of reaction varies according to the evolution of main reaction pathways. Therefore, the
275 volumetric power released (eq (8)) was directly computed as a function of *n*-octane concentration,
276 using CHEMKIN. Then, the curve obtained was fitted by a polynomial function in order to be used
277 directly in COMSOL Multiphysics (Supplementary material).

278

279 *2.6. Initial and boundary conditions*

280 In the model, the initial temperature of the top of the reservoir was set to 423.15 K, which corresponds
281 to a rather hot and deep reservoir, and the initial Bottom Hole Pressure (BHP) to 30 bar, characteristic
282 of a depleted reservoir. The temperature changes with depth are considered by a geothermal gradient
283 of 30°C/km. The initial molar fraction of *n*-octane was taken equal to 0.015 in the entire reservoir, the
284 rest being nitrogen.

285 In the well, the BHP was set to 100 bar and maintained constant during the entire injection. That means
286 that the mass flow rate of the injected gas decreases during the process. The molar fractions of oxygen
287 and nitrogen were taken equal to 0.22 and 0.78 respectively, as in our previous studies (Pacini-
288 Petitjean et al. 2015a, 2015b and 2016). It has been considered that the injection lasts until the
289 pressure reaches 100 bar in the entire reservoir.

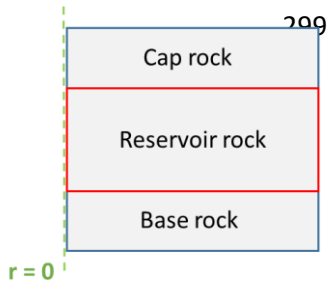
290 The boundary conditions set in this study correspond to the Neumann conditions, which means that
291 there is no heat flux or fluid flow across the boundaries (eq (14)).

$$292 \begin{cases} \vec{n} \cdot (K_{ef} \vec{\nabla} T) = 0 \\ \vec{n} \cdot (D_i \vec{\nabla} C_i) = 0 \end{cases} \quad \text{eq (14)}$$

293 Where \vec{n} is a vector normal to the wall.

294 The thermal reservoir models account for heat exchange with the surroundings, typically through heat
295 conduction to the cap and base rocks. The insulated walls are shown in Figure 2. As explained before,
296 the sides are considered as insulated walls, because the drainage zone is surrounded by other drainage

297 zones leading to symmetrical temperature profiles and so to the same temperature on both sides. On
298 the other hand, it has been assumed that the fluid does not penetrate into the cap or the base rocks.



300

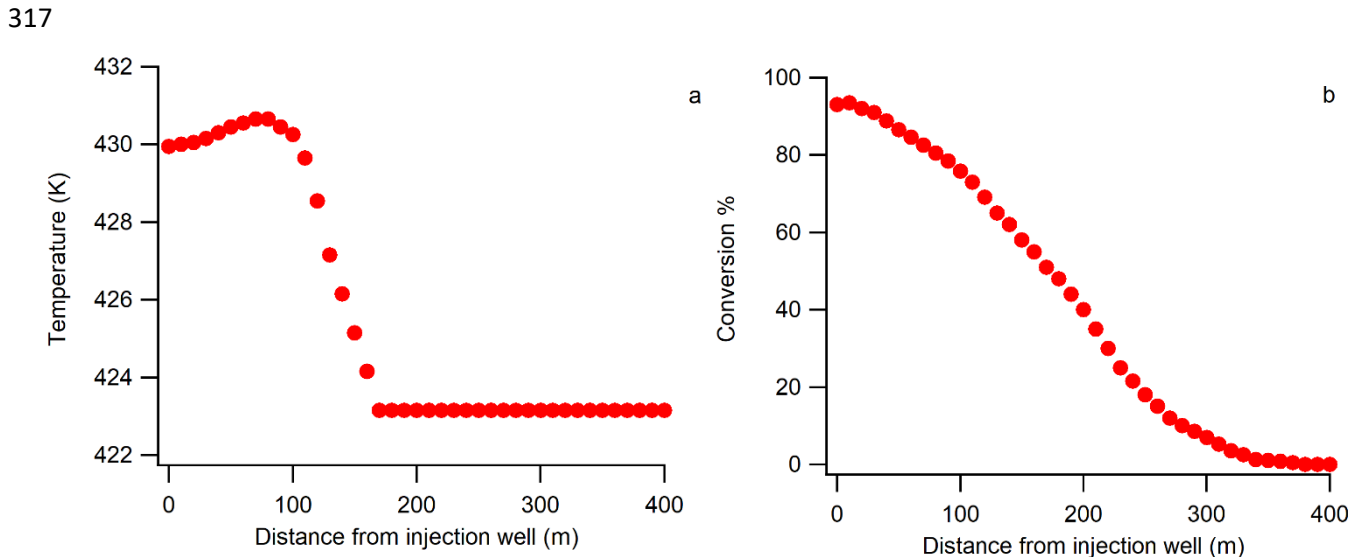
301 **Figure 2.** Walls in blue are thermal insulated, and walls in red prevent the fluid flux. Dashed green line
302 corresponds to the well location.

303 3. Results and discussion

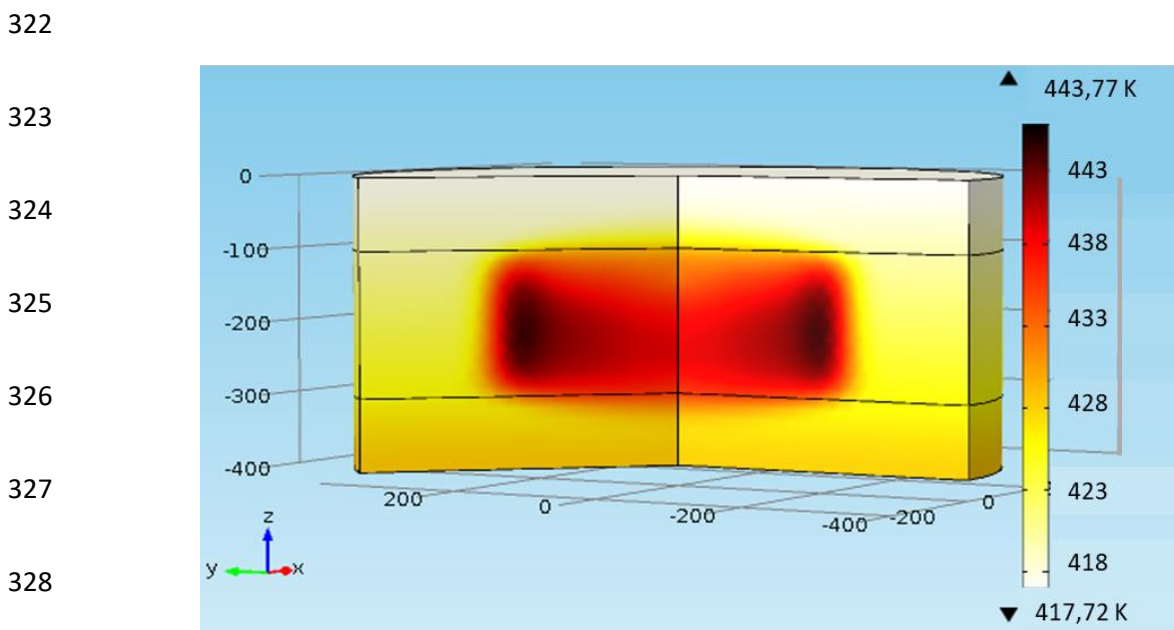
304

305 3.1. Simulations at constant enthalpy of reaction

306 The temperature and *n*-octane concentration in the entire reservoir have been computed up to about
307 10 years. For an example, Figure 3 shows the results after 1500 days (about 4 years) along a cut-line.
308 Near the well, in the first 20 m, almost all *n*-octane is consumed and the consumption of *n*-octane has
309 begun in the entire reservoir. The temperature increases in less than half of the reservoir's radius, but
310 the increase is less than 10 K. The maximum temperature is reached after 2500 days (almost 7 years –
311 end of injection) and it is equal to 443 K. The temperature profile is represented in Figure 4. The
312 maximum temperature is approximately located at 200 m of the injection well (radial position) and
313 half height. The temperature of the cap and base rocks have also increased by thermal conduction.
314 After 2500 days, the pressure reaches 100 bar and so no more gas is injected. It should be noted that
315 the auto-ignition temperature of *n*-octane is 493 K (De Soete and Feugier, 1976), which means that
316 there is apparently no auto-ignition in the reservoir in the studied conditions.



319 **Figure 3.** Temperature (a) and *n*-octane conversion (b) as a function of the distance from injection well,
 320 after 1500 days (cross-sections) – constant enthalpy of reaction (end of injection: 2500 days – $C_{p, rocks} =$
 321 $933 \text{ J.kg}^{-1}.\text{K}^{-1}$).

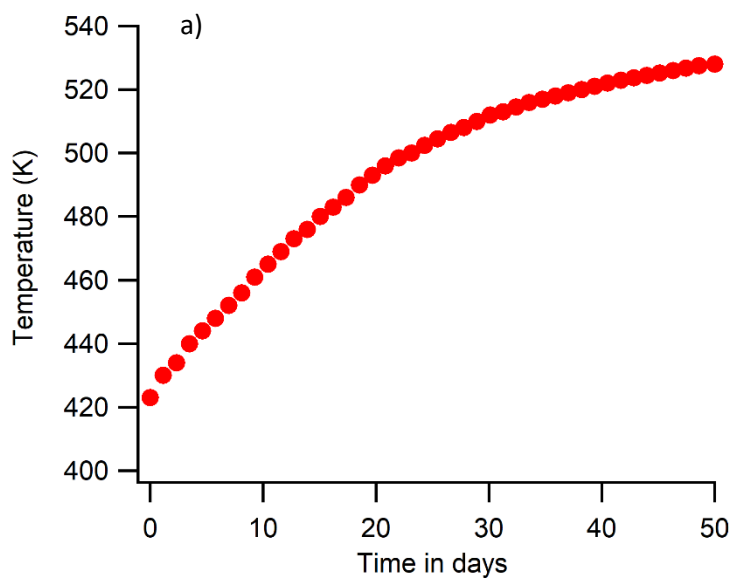


329 **Figure 4.** Temperature profile in the reservoir (between -100 and -300 m), in the cap rock (between 0
 330 and -100 m) and in the base rock (between -300 and -400 m), after 2500 days (end of injection – $C_{p, rocks}$
 331 $= 933 \text{ J.kg}^{-1}.\text{K}^{-1}$).

332 *3.2. Simulations with the heat release estimated from CHEMKIN results*

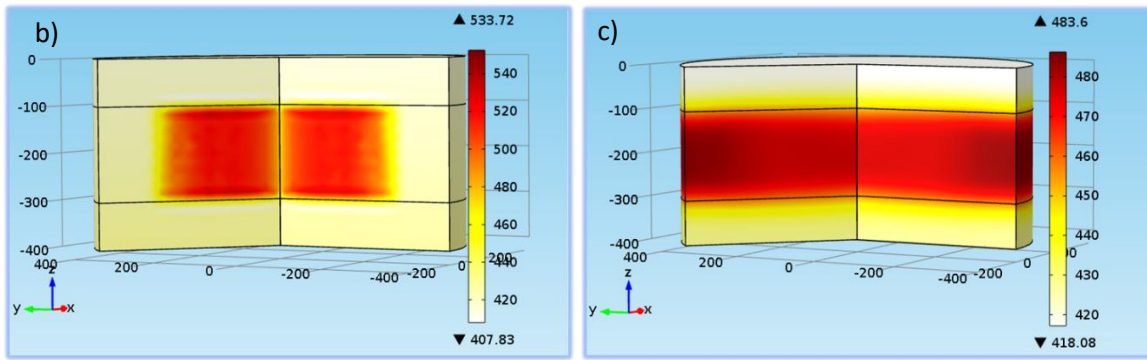
333 In this subsection, the reservoir temperature has been computed using the volumetric power release
334 estimated from CHEMKIN as a function of remaining *n*-octane concentration. The temperature
335 increases much faster than in the previous computation (Figures 5a and b) and reaches 533 K in the
336 center of the reservoir after only 60 days, which is above the auto-ignition temperature. Figure 6 shows
337 the temperature evolution as a function of time and radial position in a cross section and so the
338 reaction wave can be observed. The heat transfer by the rocks is much slower than by the fluid, so the
339 cap and base rocks warm up very slowly. According to Figure 7a, *n*-octane is reduced to half of the
340 initial value of the reservoir after 50 days. Figure 7b shows the progression of oxygen into the reservoir.
341 After the total consumption of *n*-octane, the dissipation of heat in the base and cap rocks can be
342 observed (Kihm et al., 2019) and the temperature slowly decreases and is almost uniform in the
343 reservoir and equal to 480 K after 500 days (Figure 5c).

344



345

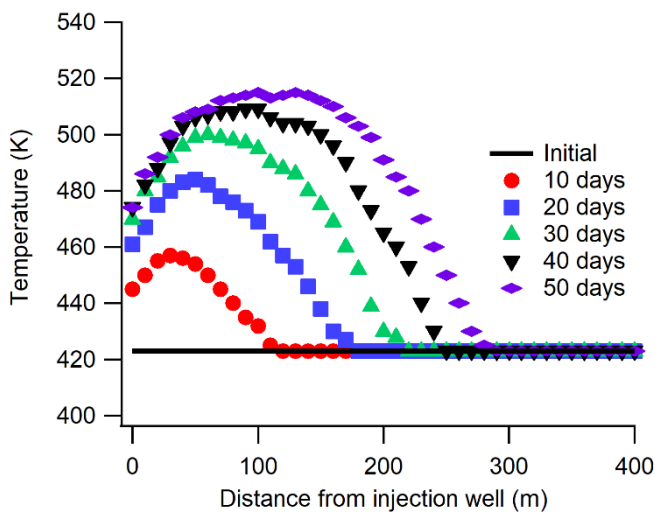
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347

348 **Figure 5.** a) Average temperature in the reservoir as a function of time until 50 days; b) Temperature
 349 profile in the reservoir after 50 days; c) Temperature profile in the reservoir after 500 days (end of
 350 injection: 300 days– $C_{p,rocks} = 933 \text{ J.kg}^{-1}.\text{K}^{-1}$).

351

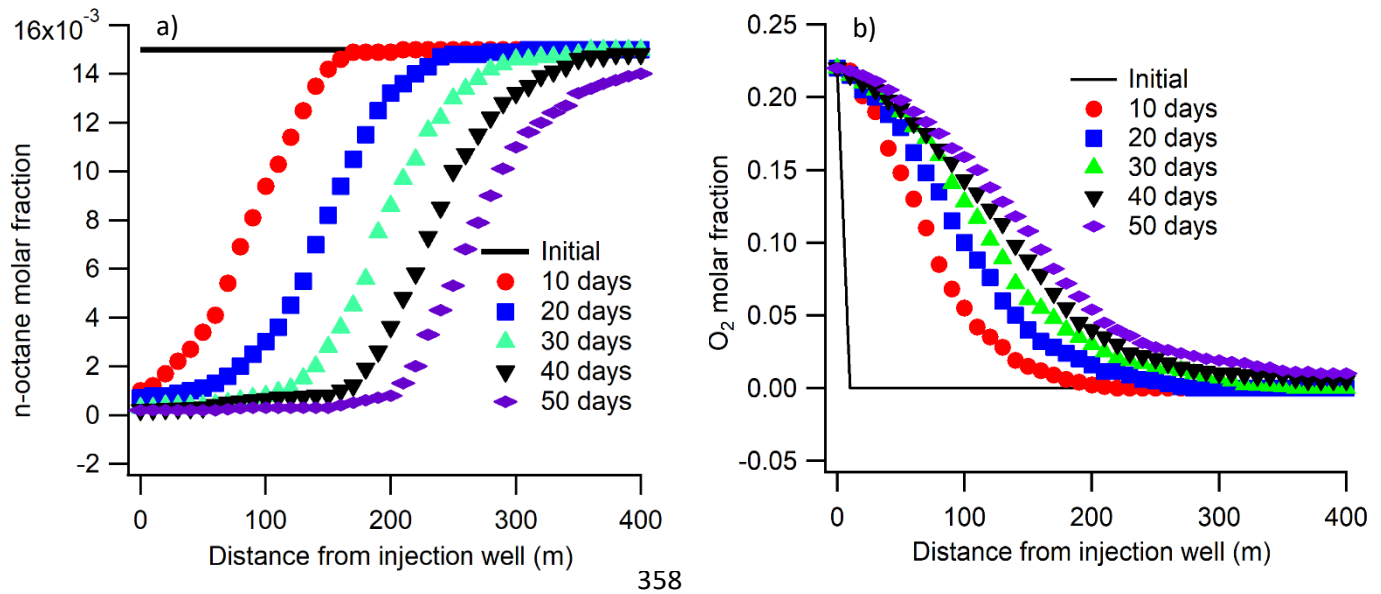


352

353

354 **Figure 6.** Temperature as a function of the distance from injection well (cross-sections after 10, 20, 30,
 355 40 and 50 days) – heat release from CHEMKIN results (end of injection: 300 days– $C_{p,rocks} = 933 \text{ J.kg}^{-1}.\text{K}^{-1}$
 356 1).

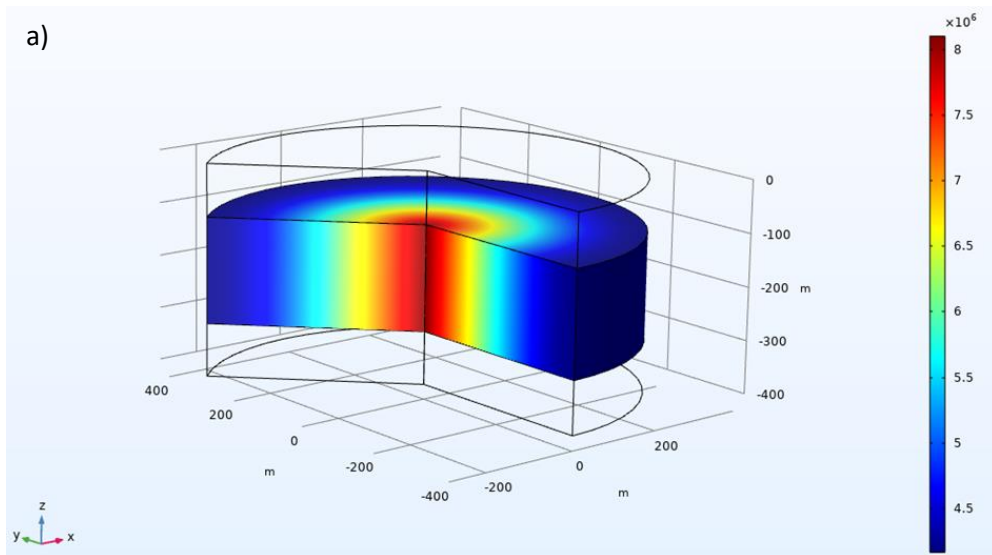
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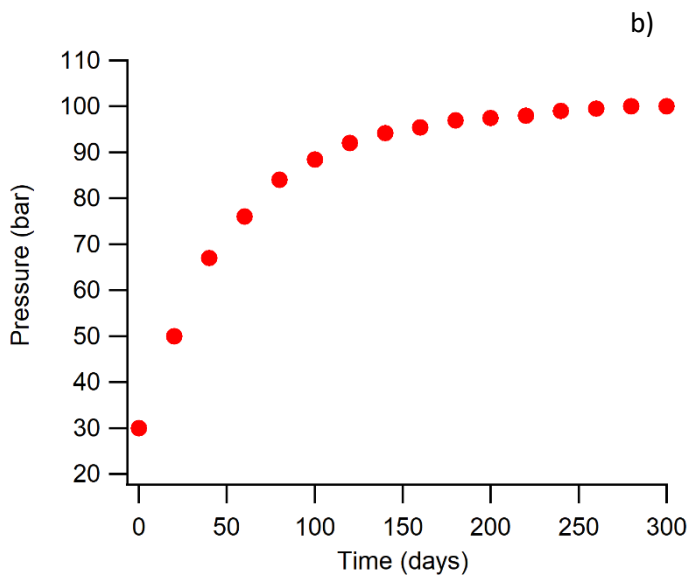
359 **Figure 7.** Molar fractions (gas phase) of the reactants (a, n -octane and b, O_2) as a function of the
 360 distance from injection well (cross-sections after 10, 20, 30, 40 and 50 days - end of injection: 300
 361 days - $C_{p, rocks} = 933 \text{ J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$).

362 Figure 8a shows the pressure profile after 50 days. Pressure reaches 80 bar near the injection well. The
 363 average pressure reaches 100 bar after almost one year (300 days - Figure 8b), which corresponds to
 364 the end of the gas injection into the well in the simulations. Indeed, the temperature increase is greater
 365 and faster with the heat release estimated from CHEMKIN results than with a constant enthalpy of
 366 reaction, and so the pressure limit is reached sooner than by the first simulations. In the meantime,
 367 the oxygen concentration slowly stabilizes.

368



370



372 **Figure 8.** a) Pressure (bar) profile in the reservoir 50 days after the beginning of the injection; b)
 373 Average pressure in the entire reservoir as a function of time (end of injection: 300 days– $C_{p_rocks} = 933$
 374 $J.kg^{-1}.K^{-1}$).

375

376 According to this study, which considers the heat released in a more realistic way than only the
 377 enthalpy of total oxidation, the gas injection into the reservoir would strongly affect the hydrocarbon
 378 composition and the temperature over a few months after the beginning of the injection.

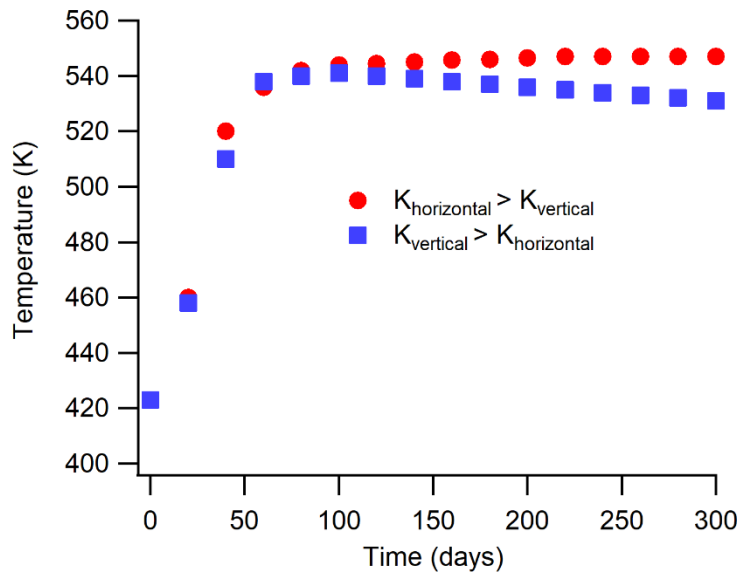
379 Nevertheless, the reactivity of *n*-octane against oxidation was estimated in homogeneous phase. The
380 reactivity in a porous medium could be lowered by the recombination of the radicals onto the walls
381 (Garo et al. 1984). The heat release and the reaction rate would be lower than estimated here.
382 Moreover, the temperature increase depends on the competition between the kinetics of oxidation
383 and the thermal properties of the rocks, which is studied in the next sub-section.

384

385 *3.3. Influence of thermal properties anisotropy on heat dissipation into the reservoir*

386 The evolution of temperature in the reservoir is governed by the competition between the
387 exothermicity of reaction, the dynamics of the gas injected and the thermal properties of the reservoir,
388 the cap and the base rocks. In order to assess the influence of anisotropy in thermal conductivity on
389 heat dissipation, first, a horizontal thermal conductivity 10 times higher than the vertical one was set
390 in the simulations (Figure 9). The opposite situation was also considered (Figure 9). The maximum
391 temperature is reached after 80 days and is about 540 K in both cases. No major difference is observed,
392 but the slope after the maximum temperature point in the case of a higher vertical conductivity than
393 the horizontal one. It could be explained by the higher vertical conductivity, which leads to more
394 thermal diffusion into the cap and base rocks. Nevertheless, the thermal conductivity of the rocks has
395 only a small influence in comparison with the dynamics of the gas injected and the heat released by
396 the oxidation.

397

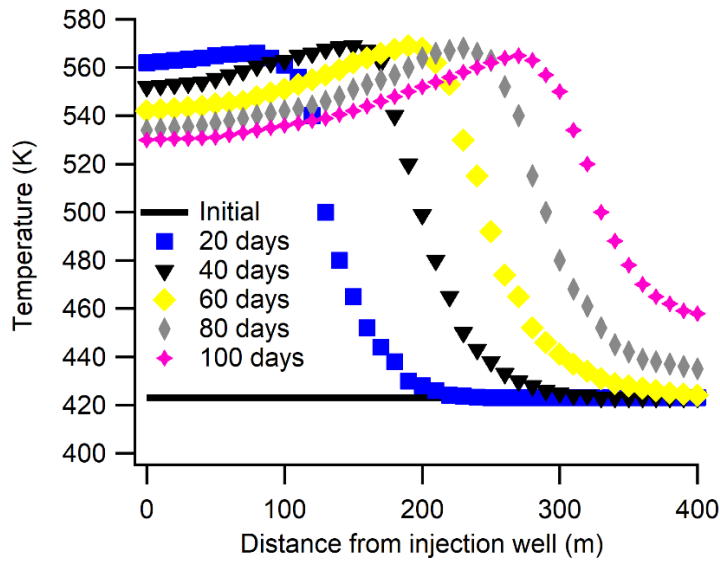


398

399 **Figure 9.** Temperature as a function of time in case of anisotropy (red circle markers: horizontal
 400 thermal conductivity 10 times higher than the vertical one; blue square markers: opposite situation -
 401 end of injection: 300 days— $C_{p \text{ rocks}} = 933 \text{ J.kg}^{-1}.\text{K}^{-1}$).

402

403 Secondly, the influence of the heat capacity of the reservoir rock was considered. The heat capacity of
 404 usual reservoir rocks ranges from 500 to 1500 $\text{J.kg}^{-1}.\text{K}^{-1}$, depending on mineral composition. These two
 405 limit cases were successively investigated. First, the heat capacity of the reservoir, cap and base rocks
 406 was set to 500 $\text{J.kg}^{-1}.\text{K}^{-1}$. In this case, the rocks warm up very quickly (Figure 10). The reservoir
 407 temperature reaches its maximum (about 570 K) after 40 days, and then the temperature slightly
 408 decreases. After 100 days, the temperature in about three quarters of the reservoir is greater than 530
 409 K.



410

411 **Figure 10.** Temperature as a function of time and distance from the injection well (heat capacity of the
 412 reservoir, cap and base rocks set to $500 \text{ J.kg}^{-1}.\text{K}^{-1}$ - end of injection: 300 days).

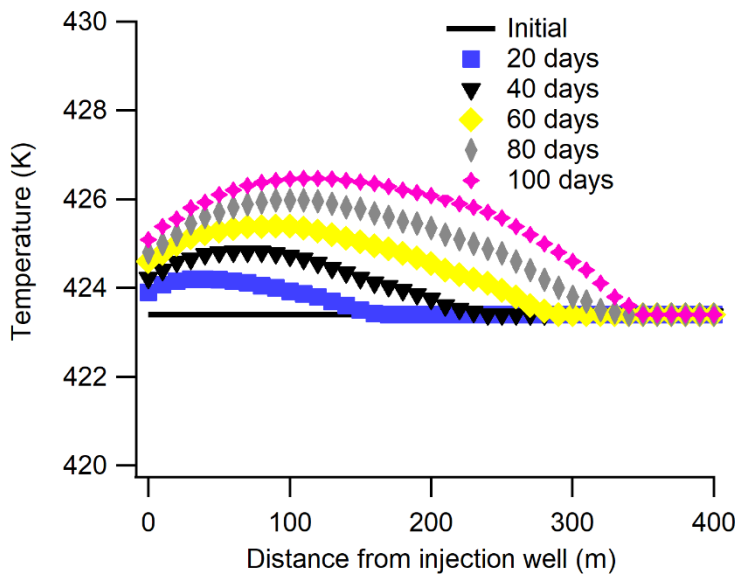
413

414 Secondly, the heat capacity was set to $1500 \text{ J.kg}^{-1}.\text{K}^{-1}$ (Figure 11). In this situation, since the rocks absorb
 415 most of the energy released by the exothermicity of the reactions, the temperature only increases by
 416 3 K after 100 days, which is completely negligible. Figure S2 compares the evolution of temperature
 417 computed for the three heat capacities tested in order to better highlight their influence. As expected,
 418 the consumption of the reactants is all the faster than the heat capacity is lower and the temperature
 419 higher, since the kinetics is temperature dependent (Figure 12). Figures S3 and S4 compare the
 420 evolution of respectively *n*-octane and O_2 molar fractions computed for the three heat capacities. The
 421 difference between the results obtained with 500 and $933 \text{ J.kg}^{-1}.\text{K}^{-1}$ is more pronounced than between
 422 933 and $1500 \text{ J.kg}^{-1}.\text{K}^{-1}$ since the increase of temperature with $500 \text{ J.kg}^{-1}.\text{K}^{-1}$ is very important. It appears
 423 that there is no major difference between 933 and $1500 \text{ J.kg}^{-1}.\text{K}^{-1}$ except at the very beginning of the
 424 injection. In the case of O_2 , since it is progressively injected into the reservoir, the concentration
 425 increases faster when the heat capacity is high. As O_2 is less consumed at low temperature, which
 426 corresponds to a high heat capacity, it penetrates the reservoir faster.

427

428 Contrary to the thermal conductivity, the heat capacity has a major influence on temperature in the
429 reservoir. In order to choose the location of the storage site, which ensures the maximum thermal
430 stability of the reservoir, it is particularly important to pay attention to the heat capacity of the
431 reservoir, base and cap rocks and to choose rocks with the highest heat capacity. Furthermore, for a
432 given depleted reservoir, it is important to adapt to the thermal characteristics of the reservoir and to
433 the remaining hydrocarbons, the composition of the injected gas, its flow rate, and the temperature
434 and the pressure in the well, in order to limit the risk.

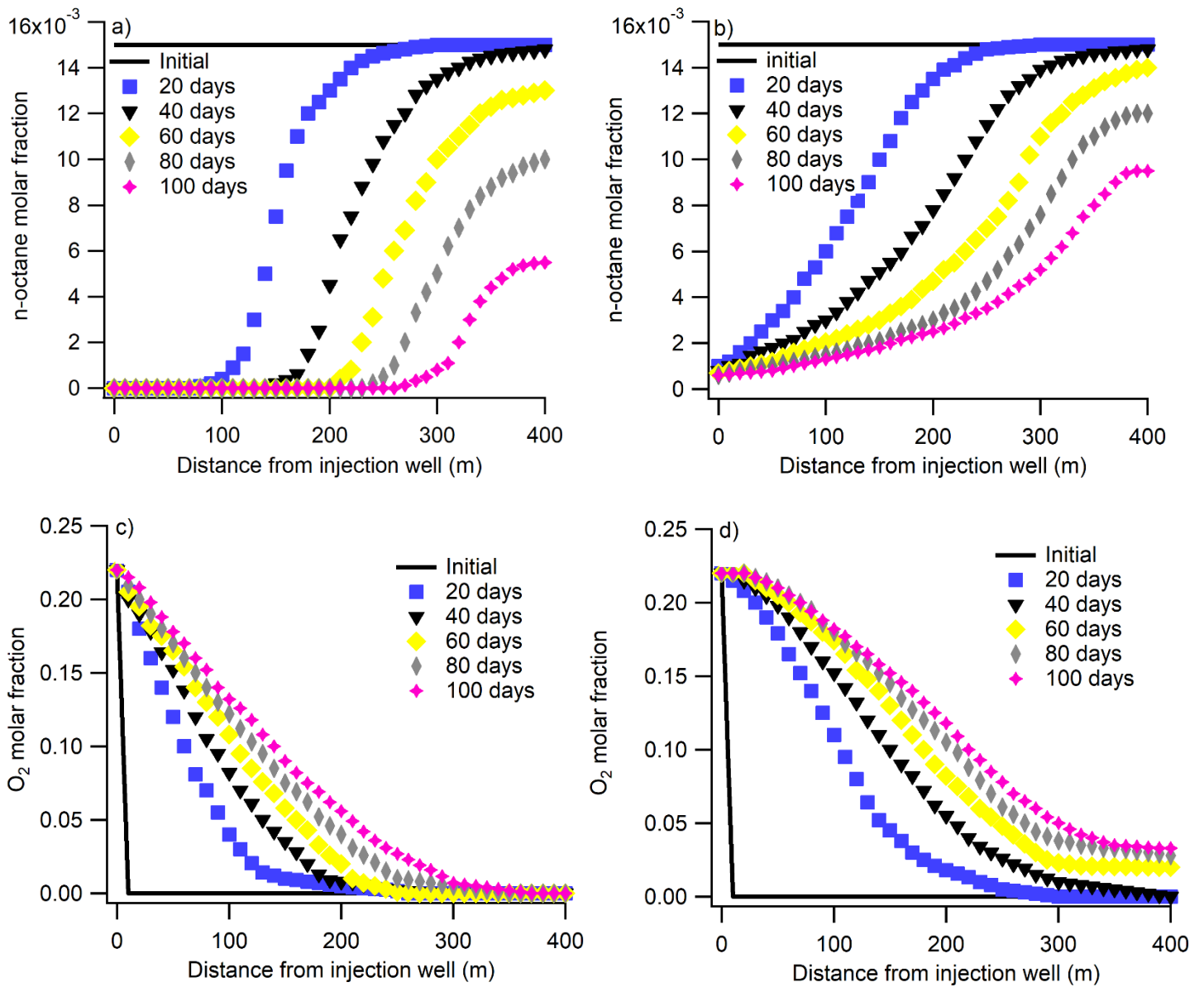
435



436

437 **Figure 11.** Temperature as a function of time and distance from the injection well (heat capacity of the
438 reservoir, cap and base rocks set to $1500 \text{ J.kg}^{-1}.\text{K}^{-1}$ - end of injection: 300 days).

439



440

441 **Figure 12.** Reactant molar fractions as a function of time and distance from the injection well: a) *n*-
 442 octane (reactant present in the reservoir) and heat capacity of the reservoir, cap and base rocks set to
 443 $500 \text{ J.kg}^{-1}.\text{K}^{-1}$; b) *n*-octane and heat capacity set to $1\ 500 \text{ J.kg}^{-1}.\text{K}^{-1}$; c) O_2 (injected into the reservoir) and
 444 heat capacity set to $500 \text{ J.kg}^{-1}.\text{K}^{-1}$; d) O_2 (injected into the reservoir) and heat capacity set to $1\ 500 \text{ J.kg}^{-1}.\text{K}^{-1}$ (end of injection: 300 days).

446

447

448 **4. Conclusions**

449

450 The paper aimed at studying the influence of oxygen contained in the flue gas, for an example captured
451 after oxy-combustion, on the oxidation kinetics of the residual fuel in the depleted reservoir and its
452 influence on some of the reservoir parameters or properties. More precisely, the paper focused on the
453 influence of the oxidation exothermicity and the thermal properties of the rocks on the pressure and
454 the temperature evolution, and on the kinetics of the reaction inside a depleted oil reservoir.

455 Despite some assumptions and a quite high initial concentration of O₂, it clearly appears that the
456 system is not inert with respect to the hydrocarbons reactivity. The residual oil is probably consumed
457 in less than one year in relatively small reservoirs, leading to the temperature increase throughout the
458 reservoir. After the end of the reaction, the temperature decreases slowly by heat dissipation in the
459 cap and base rocks.

460 The extent of increase in the reservoir temperature is clearly a function of the thermal properties of
461 the rocks. It has been demonstrated that the thermal conductivity of the rocks and their potential
462 anisotropy have only a small influence on the temperature evolution. Nevertheless, the temperature
463 evolution is highly dependent on the heat capacity and quite contrasting behaviors can be observed
464 depending on low or high heat capacities of the rocks. A high heat capacity of the rock leads to a
465 negligible increase in the temperature of the reservoir. Therefore, in order to limit the temperature
466 increase and the kinetics of the reaction, it appears important to select reservoirs whose rocks as well
467 as the rocks of the cap and the base have a high heat capacity (greater than 1000 J.kg⁻¹.K⁻¹), especially
468 when O₂ concentration in the injected gas is high (at least greater than 3%). In this case, the
469 temperature increase seems negligible and the oil oxidation is very slow, which appears to be an
470 essential condition for a safe injection operation. Consequently, for a depleted reservoir with rocks of
471 low heat capacity, it is important to decrease the O₂ concentration in the injected gas, its flow rate,

472 and the temperature and the pressure in the well, in order to limit the risk of significantly rapid rise in
473 the reservoir temperature and its associated thermal instability issues.

474

475

476

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479

480

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