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5 6 7	Tannin-based resins for 3D printing of porous carbon architectures
8	Pauline Blyweert †, Vincent Nicolas †, Jan Macutkevic‡,
9	Vanessa Fierro †, Alain Celzard †,*
10	
11	† Université de Lorraine, CNRS, IJL, F-88000 Epinal, France
12	‡ Faculty of Physics, Vilnius University, Sauletekio 9, LT-10222 Vilnius, Lithuania
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^{*}Corresponding author: Pr. Alain Celzard (<u>alain.celzard@univ-lorraine.fr</u>), Université de Lorraine

17 KEYWORDS

18 Tannin, Porous carbon, Stereolithography, Additive manufacturing

19

20 ABSTRACT

21 Carbon structures, unlike metals or polymers, can hardly be directly implemented in additive 22 manufacturing, and their sustainable nature is extremely limited. In this work, we demonstrate 23 that porous carbon structures can be obtained by laser stereolithography of acrylate-tannin mixed 24 resins, followed by a pyrolysis step. By adjusting the initial acrylate / tannin ratio in the resin, the 25 density and volume shrinkage of the resultant carbon, and consequently the mechanical properties 26 of the corresponding carbon-based architectures, could be modified. The moderate electrical conductivity of the structures, reaching about 7 S.cm⁻¹, and their broadband absorption in 27 28 microwaves open up possibilities for electrochemical or electromagnetic applications. Thus, the 29 possibility of obtaining complex free-standing structures with imperceptible warpage, low 30 volume shrinkage, and adjustable density offers an opportunity to develop 3D-printed carbon 31 materials with a significant proportion of bio-based precursors, which can be easily adapted for a 32 large number of applications.

33

34 INTRODUCTION

35 Porous carbon materials, thanks to their excellent chemical stability and unique mechanical, 36 electrochemical, thermal and electrical properties ^{1,2}, are used in many industrial processes such 37 as heat management ³, gas storage ⁴, electrochemical energy storage and conversion ^{5,6}, among 38 others. Given their central role, it would make sense to find a fast, inexpensive and accurate way to produce tailored complex systems, not only to broaden the application range of these materials,
but also to give them a more ecological character. This versatility can be achieved with 3D
printing methods ⁷, and especially by vat photopolymerisation, one of the most accurate
technologies ^{8,9}. Nevertheless, since pure carbon cannot be melted, sintered or polymerised, this
particular element is not directly applicable in additive manufacturing. Thus, most 3D-printed
carbon materials are obtained by printing a carbon precursor, followed by a pyrolysis step.

45 Existing studies have reported the printing of porous carbon through stereolithography 46 technology using commercial or custom resins, followed by a pyrolysis step. These approaches 47 have allowed for the realisation of highly porous carbon materials with tailored pores in a reduced number of steps ^{10–13}. Nevertheless, in most cases, the printed resin had a low carbon 48 yield 14,15 and underwent significant volume shrinkage (50 – 90%). Besides, the photocurable or 49 thermosets resins ¹⁶ used in those approaches are all of petrochemical origin. To make 3D-printed 50 51 porous carbons more sustainable, it is therefore interesting to use bio-based resins as carbon 52 precursors. Unlike previous studies in which bio-based carbon precursors such as lignin, cellulose or monosaccharides have been used in Direct Ink Writing technology^{17–19}, in this work we 53 54 present an approach to print highly porous carbons with low shrinkage and warpage from bio-55 based photoresin using laser stereolithography.

This approach consists of adding a wood extract, condensed (mimosa) tannin, to the photocurable resin to build complex structures and then converting the latter into carbon through a pyrolysis step. Indeed, mimosa tannin is an abundant phenolic bioresource with a carbon yield on pyrolysis of about 47%, which has been widely studied over the years to synthesise a wide range of carbon structures ²⁰. Herein, we show for the first time the possibility of printing a tannin-based resin to ensure high carbon yield with low shrinkage upon carbonisation, and we discuss some properties of these new materials based on rheological, structural and physicalcharacterisation.

64

65 EXPERIMENTAL SECTION

66 Materials

67 Hexanediol diacrylate (HDDA, SR238), pentaerythritol tetraacrylate (PETA, SR295) and

68 acrylate oligomer (CN154 CG) were kindly supplied by Sartomer (Arkema Group, France).

69 Bis(2,4,6-trimethylbenzoyl)phenylphosphine oxide (BAPO) was supplied by Lambson (Arkema

70 Group, England). Mimosa (Acacia Mearnsii) tannin extract, a light-brown powder containing 80-

71 82% phenolic material based on flavonoid units, commercially available under the name Fintan

72 OP, was kindly provided by SilvaChimica (St Michele Mondovi, Italy) and used as a bio-based

carbon precursor. All products were used as received without further purification.

74 Preparation of tannin-acrylate resins, 3D printing and pyrolysis

A BAPO photoinitiator (content 0.30 wt.%) and the acrylates monomers (SR238, SR295 and CN154 CG), added in the proportions shown in Table S1, were mixed thoroughly at 50 °C for 20 min to obtain a homogeneous resin. Then, a fixed amount of tannin was added and mixed by hand to the acrylate resin to obtain tannin-acrylate resin with a tannin content of 0, 10, 20, 25 or 30 wt.%. The corresponding formulations and printed structures were named T-0, T-10, T-20, T-25 and T-30, respectively, and the derived carbons were called CT-0, CT-10, CT-20, CT-25 and CT-30, respectively.

82 A commercial desktop SLA ($\lambda = 405$ nm, DWS J28, DWS, Italy) was employed to print 3D 83 structures from tannin-acrylate resins. After the printing procedure, the structures were removed from the build platform and rinsed twice in an isopropanol bath for 10 min to remove uncured resin. Afterwards, the products were post-cured in a UV oven ($\lambda = 405$ nm) for 20 min at room temperature.

The printed meshes were finally converted into carbon materials under pure nitrogen (flow rate 100 mL·min⁻¹) in a tubular furnace with a heating rate of $1.5 \, {}^{\circ}\text{C}\cdot\text{min}^{-1}$ to 400 ${}^{\circ}\text{C}$, then a 20 min step at 400 ${}^{\circ}\text{C}$, followed by a second ramp of 2.5 ${}^{\circ}\text{C}$ min⁻¹ up to the final temperature of 900 ${}^{\circ}\text{C}$, which was maintained for 1 hour (Figure 1). This particular temperature program was chosen to stabilise the architectures during thermal degradation.

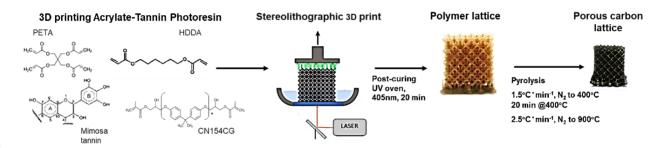




Figure 1. Schematic illustration of the fabrication strategy for 3D-printed tannin-based carbonmaterials.

95 Characterisation

96 Materials characterisation

- 97 A Brookfield DVII+ PRO viscometer with cone-plate geometry (spindle CPE-51, gap 0.013
- 98 mm) was used to measure the viscosity of the resins at room temperature.
- 99 The volume-based particle size distribution of the tannin powder was measured using a laser
- 100 diffraction particle size analyser (Mastersizer 3000, Malvern), equipped with an automated dry
- 101 powder dispersion unit (Aero S, Malvern). The powder was dried before measurements.

102	Thermogravimetric analysis (TGA) coupled with Differential Scanning Calorimetry (DSC)
103	was performed using an STA 449 F3 Jupiter [®] apparatus (Netzsch) and provided access to the
104	carbon yield (ratio of the weight of carbon after pyrolysis to the weight of the initial resin) and
105	glass transition temperature of the different resins, respectively. Cured resin samples (~ 10mg)
106	were tested from 30 to 800 °C with a heating rate of 10 °Cmin ⁻¹ under argon flow. The argon
107	flow was then replaced by air to deduce the ash content.
108	Elemental analysis (EA) of the resultant carbon materials was performed in an Elementar
109	Vario EL Cube analyser.
110	Structure, porosity, and textural properties of carbon materials
111	The linear shrinkage was calculated using Eq. (1), where the dimension can be height,
112	thickness or length of the printed architectures. The data reported are the average values of at
113	least 5 measurements on 3 samples obtained from each carbon type.
114	$Shrinkage (\%) = \frac{Dimension \ before \ pyrolysis - Dimension \ after \ prolysis}{Dimension \ before \ pyrolysis} \times 100 $ (1)
115	Scanning Electron Microscopy (SEM) and Energy-dispersive X-ray spectroscopy (EDX) were
116	performed with a Zeiss GeminiSEM 360 microscope. Transmission Electron Microscopy (TEM)
117	images were obtained using a JEM ARM 200F Cold FEG TEM/STEM operating at 200 kV,
118	using spherical aberration (Cs) probe and image correctors. The samples were prepared by
119	dispersing carbon powder in ethanol, after which a single drop of the suspension was deposited
120	on a carbon-coated copper grid (200 mesh).
121	X-ray diffraction (XRD) was carried out using a Panalytical X'Pert Pro diffractometer in a
122	Bragg-Brentano configuration in reflection with a Cu anode.

Raman spectra were recorded with a Horiba XploRa Raman apparatus equipped with a 50×
long-range objective (wavelengths of 532 and 638 nm, filtered at 10% of their maximum energy).
The spectra obtained at a wavelength of 638 nm were then fitted by the method described by
Mallet-Ladeira et al. ²¹ using a double Lorentzian for the D band and an asymmetric mixed
Gaussian-Lorentzian profile for the G band.

128 The bulk density ($\rho_{b,}$ g cm⁻³) was calculated from the measured mass and volume of three 129 samples of each carbon type obtained. The skeletal density ($\rho_{s,,}$ g cm⁻³) was determined by helium 130 pycnometry using an Accupyc II 1340 (Micromeritics) automatic apparatus, averaging 30 131 measurements. Sample were ground and dried before each measurement. The total porosity, Φ 132 (dimensionless), was then calculated according to Eq. (2):

133
$$\Phi = 1 - \frac{\rho_b}{\rho_s} \tag{2}$$

N₂ and H₂ adsorption experiments at -196°C were carried out with a Micromeritics 3-flex automatic adsorption apparatus. Samples (~ 1 g) were degassed under vacuum before analysis for at least 48h at 200 °C. The Brunauer-Emmett-Teller ^{22,23} (BET) method was applied to N₂ isotherms to obtain the BET area (A_{BET}). Pore size distribution (PDS) and total surface area (S_{NLDFT}) were calculated from the combination of N₂ and H₂ adsorption isotherms using a twodimensional non-local density functional model (2D-NLDFT) available in SAIEUS software (Micromeritics ²⁴).

141 Mercury intrusion was performed using a Micromeritics AutoPore IV 9500 up to a pressure of142 414 MPa.

143 Physical properties of carbon materials

144 The electrical conductivity σ (S cm⁻¹) was measured by the four-point method using a Keithley 145 6430 source-measure unit on 3D-printed and then pyrolysed 2 mm × 2 mm × 15 mm 146 parallelepipeds. The electrical conductivity was calculated according to:

147
$$\sigma = \frac{2I}{(|V^+| + |V^-|)} \frac{L}{s}$$
(3)

where V^{+} and V are the voltage values (V) measured with the current I (A) flowing in one 148 149 direction and then in the other, so as to correct the measurements for a slight thermoelectric effect 150 that is always present. S (cm²) is the cross-sectional area and L (cm) is the distance between the 151 two opposite sides on which the electrical contacts were placed. The electrical conductivity of the 152 conductive skeleton, σ_s , was calculated with Eq. (4), which is a good approximation for rigid open-cell foams with node effects ^{25,26}. This equation was preferred to others corresponding to 153 154 closed porosity, because polymeric precursors produce a significant amount of gases during 155 pyrolysis and hence a totally open porosity, as already proved by the mercury intrusion 156 experiments (see below):

157
$$\sigma_{s} = \frac{3\sigma}{\left[(1-\phi)+2(1-\phi)^{\frac{3}{2}}\right]}$$
(4)

Microwave measurements in the 24-39 GHz (Ka-band) frequency range were carried out with an ELMIKA R2-408R scalar network analyser. The samples were placed directly in a 7.2×3.4 mm cross-section waveguide, perpendicular to the wave propagation, and the modules of the amplitude of transmitted (S_{21}) and reflected (S_{11}) signals were measured. From these data, the (dimensionless) coefficients of reflection (R), transmission (T) and absorption (A) were calculated according to Eq. (5):

164
$$R = S_{11}^{2}$$

165 $T = S_{21}^{2}$
166 $A = 1 - S_{11}^{2} - S_{21}^{2}$
(5)

Mechanical properties of the different 3D printed carbon were investigated by quasi-static compression using an Instron 5944 universal testing machine equipped with a 2 kN load cell. The compression tests were carried out in triplicate at a constant load rate of 2 mm·min⁻¹.

170 More details of sample preparation and methods are given in the Supplementary Information.

171 RESULTS AND DISCUSSION

172 Rheological properties and thermal behaviour of the resins

In stereolithography, the viscosity of the resin is a crucial parameter. In most cases, a viscosity of less than 3 Pa·s (at the printing temperature and a shear of about 100 s^{-1}) is desired to promote recoating and self-levelling of the resin ^{27,28}. The viscosity of uncured acrylate-tannin resins with a tannin content between 10 and 30 wt.% was determined as a function of shear rate in a region of 10 to 200 s⁻¹ (Figure 2a). At higher tannin contents, the resultant resins were in the form of a paste that cannot be used in a stereolithography apparatus.

179 The unfilled resin (here referred to as T-0) showed Newtonian behaviour with a viscosity of about 1.0 Pa \cdot s at a shear rate of 100 s⁻¹. It turns out that Mimosa tannin is not soluble in the 180 181 acrylate resin, and is therefore dispersed in it in the form of solid particles with an average 182 diameter of 54.5 µm, according to measurements made with the Mastersizer (Figure S1). It is 183 well established that, due to particle interactions, the viscosity of suspensions increases with the volume concentration of solids ²⁹. Thus, as expected, the viscosity increased with the tannin 184 185 content (Figure 2b). For all filled resins, shear-thinning behaviour in the low-shear rate region (below 10 s⁻¹) was observed, with viscosity decreasing to a steady state (tending towards 1.3 Pa \cdot s) 186

187 at high shear rates. This observation, typical of highly particle-loaded suspensions, can be
188 explained by the variable interparticle and particle-fluid interactions under shear ^{30–34}. As the
189 shear stress increases, immobilised acrylate chains are indeed released, and the friction between
190 tannin particles decreases. Therefore, from a macroscopic point of view, the viscosity decreases
191 as seen in Figure 2a. Due to the initial viscosity of the acrylate resin, the suspension of tannin
192 particle remains stable during the printing processes (and even for weeks at room temperature),
193 which limits the heterogeneity of the final structure.

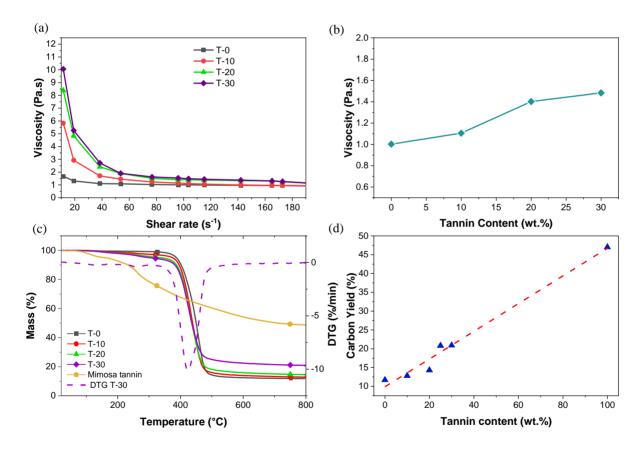


Figure 2. (a) Viscosity versus shear rate of acrylate resins containing different weight fractions of tannin (from 0 wt.% in T-0 to 30 wt.% in T-30). (b) Viscosity of acrylate resin at a shear rate of 100 s^{-1} as a function of tannin content. (c) TGA (mass loss) curves, under argon and up to 800 °C, of the same acrylate resins as in (a); that of pure tannin is given for comparison, and the DTG

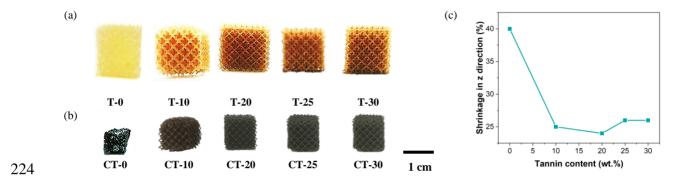
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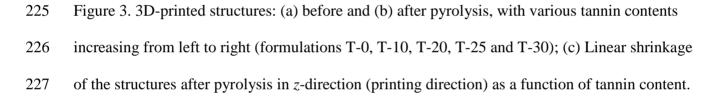
curve is provided as a guide for sample T-30 only, as the others DTG curves are similar. (d)
Carbon yield at pyrolysis as a function of initial tannin content; the value corresponding to a resin
with 25 wt.% tannin was also added.

202 The temperature behaviour of the acrylate-tannin resin was evaluated by TGA and the 203 corresponding differential curves (DTG) (Figure 2c). The UV-cured acrylate resin without tannin 204 started to lose weight at about 320 °C and degraded to char by 680 °C with a mass loss of about 205 85%. This degradation corresponds to the depolymerisation of acrylates at high temperature as 206 well as to the random radical scission of the free linear chains, leading to volatile chain fragments, 207 including ethylene and methane, the quantitative yield of which depends on the stability of the radicals formed but also on side-chain reactions (side group elimination or rearrangement)^{35–37}. 208 In the presence of tannin, three regions of mass loss were clearly identified ³⁸. The first region is 209 210 located between 30 and 170 °C, with a DTG peak around 132 °C, and corresponds to the 211 evaporation of absorbed water and the formation of small volatile compounds (CO, CO₂) released 212 at the beginning of tannin degradation. The second region between 170 and 310 °C, with a DTG 213 peak around 260 °C, is due to the degradation of tannin and acrylate resin. Finally, the 214 decomposition of the rigid segments of the resin and the formation of carbon material occur 215 between 310 and 680 °C. As expected, the addition of tannin significantly increased the final 216 carbon yield at 800 °C, from 11.7 to 20.8 % for the samples with 0 wt.% and 30 wt.% of tannin, 217 respectively. The carbon yield increased linearly with the weight fraction of tannin, as shown in 218 Figure 2d, where the carbon yield of pure tannin (47 %) has been plotted. In addition, for all 219 samples, the ash content remained below 0.5 wt.%.

220 Morphology and composition of printed carbon structures

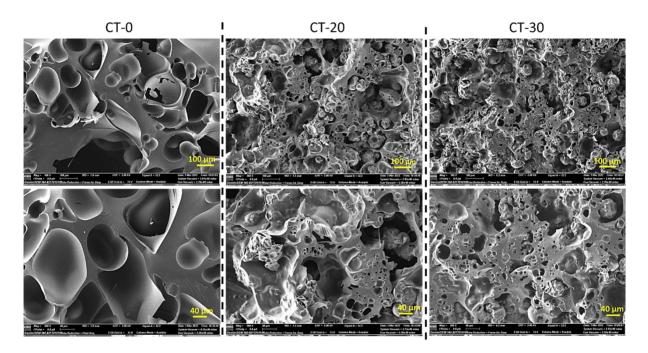
Tetrakaidecahedra meshes (Figure 3a, wall size of 200 µm) and solid cubes were printed from
30 µm-thick layers. After pyrolysis, the colour of the samples changed from light brown to bright
black for the acrylate resin without tannin to matte black for the CT-30 structure (Figure 3b).





228 The addition of tannin to the formulation was found to have an effect on the structure 229 shrinkage as evidenced in Table S2. The samples with tannin, regardless of their architecture, 230 tended to shrink linearly by 25% in each direction, while the tannin-free structure shrank by 40% 231 in the z-direction and by 30% in x-y directions with warpage. The shrinkage phenomenon results 232 from the evolution of volatile matters during carbonisation. During the pyrolysis process, the 233 acrylate polymer matrix is progressively heated to temperatures above its glass transition 234 temperature (found to be 126 °C by DSC measurement for the pristine resin, and reduced to 115 °C for the T-25 resin due to polymer/particles interactions ^{39,40}, Figure S2) and enters a more 235 236 flexible state subject to deformation. As the tannin degrades significantly less than the resin and 237 loses much less weight, even at much higher temperatures (Figure 2c), it stabilises the structure 238 during this plastic state and thus reduces warpage in different directions.

SEM images of carbon structures CT-0 and CT-30 are presented in Figure S3 and Figure 4.
Rather smooth surfaces with a few craters allowing gases formed during the pyrolysis step to
escape were observed on all structures (Figure S3). In contrast, SEM images after fracturing
(Figure 4) revealed extensive internal porosity with many large oval bubbles in the CT-0
structure, while narrower pores can be seen on the CT-30 carbon architecture. EDX results (not
shown) evidenced similar carbon and oxygen composition for each sample.



245

Figure 4. SEM images of a fractured area of CT-0, CT-20 and CT-30 samples of 3D-printed
structures. Scale bars: 100 μm (top row), and 40 μm (bottom row).

Table S3 presents the bulk chemical composition of the different 3D-printed carbons. The nitrogen (N) content is very low and no sulphur (S) was detected, while oxygen (O) is present in significant amounts in the carbon structures. These measurements are in agreement with the composition of other carbons based on Mimosa tannin ^{20,41}, as no doping treatment was carried out on these materials. Given the usual uncertainties on this kind of analysis, it is difficult to consider that the samples are really different in terms of composition, and we can assume that all of them have C, H, N, S and O contents of about 95, 1.3, 0.3, 0 and 3.5 wt.%, respectively. The differences between the samples do not show any clear trend, and can be attributed to inhomogeneities, despite the repetitions and calculated averages.

257 Structural properties of printed carbon structures

TEM observations were performed to investigate the nanostructure of the carbon materials (Figure S4). At high magnification, TEM evidenced a highly disordered structure in all samples, which was expected from a phenolic precursor in the absence of graphitisation catalyst ²⁰. Highresolution images and the presence of a well-defined halo seen by Fast Fourier Transform (FFT) also revealed the existence of very small lamellar domains dispersed in a highly disordered matrix.

264 Figure 5a shows the Raman spectra obtained at a wavelength of 532 nm, *i.e.*, with a green 265 laser, which corresponds to the vast majority of carbon spectra reported in the literature. They are completely typical of highly disordered carbons, as expected from the precursors used ²⁰. Indeed, 266 at first order, the D-band, whose maximum is around 1350 cm⁻¹, is very broad with a shoulder 267 around 1170 cm⁻¹, and the G-band, centered around 1600 cm⁻¹, is just a little less broad. The 268 269 intensity ratio of the D/G bands is very close to 1, and at second order, the profile between 2200 and 2700 cm⁻¹ is not structured. The spectra are very similar from one formulation to the other, 270 271 especially when tannin is present. The carbon obtained with the formulation without tannin 272 differs very slightly from the others by its somewhat shallower valley between the D and G bands. 273 Moving from sample CT-0 to sample CT-30, the D/G intensity ratio increases very slightly, from 1.01 to 1.04. Considering the very disordered and non-graphitizable character of these carbons, 274

this evolution corresponds to a slight increase in the size of the coherent domains, measured through the $L_{a,Raman}$ parameter.

277	Figure 5b shows the first-order spectra, now obtained with the 638 nm laser, <i>i.e.</i> , with a red
278	laser, at the wavelength of which equation (S2) applies. As expected, the D band has moved to
279	lower Raman shifts, around 1320 cm ⁻¹ , while the G band has remained fixed. The D/G intensity
280	ratios also increased, but the trend remained the same, with values increasing from 1.13 to 1.16
281	for samples CT-0 and CT-30, respectively. Figure 5d shows the fit of these two extreme samples
282	by the method described in ²¹ . The fits provided the half-widths at half maximum of the G band,
283	which decreased from 59.5 cm ⁻¹ for CT-0 to 55.2 cm ⁻¹ for CT-30, with which the values of $L_{a,}$
284	R_{Raman} could be calculated from Eq (S2). The corresponding values and their uncertainties are
285	given in Figure 5c. On average, $L_{a, Raman}$ thus increases from 1.6 to 2.5 nm, and the three
286	formulations containing tannin gives values very close to each other. Such small values definitely
287	confirm that these carbons are highly disordered and not graphitizable.

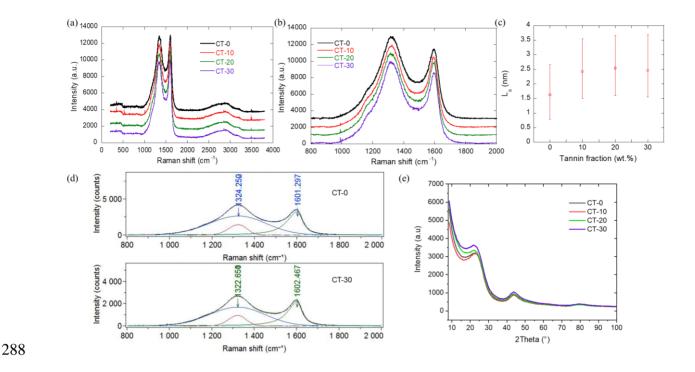


Figure 5. Raman spectra and XRD patterns of 3D-printed carbon structures for various initial acrylate / tannin ratios: (a) Full spectra obtained at a wavelength of 532 nm; (b) First-order spectra obtained at a wavelength of 638 nm. In (a) and (b), the spectra have been shifted with respect to each other along the intensity axis (whose units are arbitrary) to better distinguish them; (c) Changes of $L_{a, Raman}$ with the tannin content in the formulation; (d) Fits of the spectra shown in (b) according to the method detailed in ²¹; (e) XRD patterns.

The XRD patterns are shown in Figure 5e. As expected, the XRD profile was nearly the same for all samples, with the 002 reflection appearing as a broad band with a maximum at Bragg angle 2θ of 20-24°. The presence of such a broad reflection confirmed the existence of crystalline carbon nano-domains in a highly disordered carbon matrix, as expected from the observation in the high-magnification TEM images. Bands 10λ ($2\theta = 44^\circ$) and 11λ ($2\theta = 80^\circ$) appeared as broad and low-intensity reflections, which is indeed typical of disordered carbon. This is further confirmed by the corresponding values of $L_{a, XRD}$, calculated from Eq. (S1), which are extremely low, around 2 nm for all samples. These values are in agreement with those calculated above by
Raman spectrometry. Finally, the high intensities observed at low angles indicate the presence of
microporosity.

305 Textural properties of printed carbon structures

306 The bulk density of the 3D-printed structures after carbonisation depends on the acrylate / 307 tannin ratio (Table S4). The carbonaceous material obtained from the pure acrylate resin was denser (0.442 g cm⁻³) than its counterparts derived from tannin-acrylate mixed resins, with bulk 308 density values ranging from 0.305 to 0.408 g cm⁻³ when the tannin content in the formulation 309 310 increased from 10 to 30 wt.%. This difference can be explained by the significant linear 311 shrinkage observed with sample CT-0 (40 % in the z-direction), compared to the other samples 312 with an average shrinkage of 25 % (in the z-direction) (see again Table S2). In contrast, the 313 carbon skeletal densities remained the same, within the standard deviation, for tannin content between 0 and 30%, with an average value of 1.80 g cm⁻³. The high total porosity, ranging from 314 315 76 to 83 %, is mainly due to the depolymerisation of acrylates and the gas released during 316 transformation of tannin upon pyrolysis.

The influence of tannin content on the pore structure after pyrolysis, was studied. Figure 6 shows the N₂ and H₂ adsorption-desorption isotherms, from which the textural properties of CT-20, CT-25 and CT-30 were calculated. The outgassing of samples CT-0 and CT-10 was too long, which indicates extremely narrow porosity preventing any correct adsorption measurement. We therefore report results for only 3 of the 5 carbon structures, which are given in Table S4. Figure 6a shows type Ia nitrogen adsorption-desorption isotherms, characteristic of microporous materials with predominantly narrow (width less than ~ 1 nm) micropores, following the IUPAC

324	classification ²³ . All three isotherms exhibited a low-pressure hysteresis that it is due to slow
325	desorption in these very narrow pores. The amount of nitrogen and hydrogen adsorbed increased
326	with the amount of tannin added, although the largest increase was observed for CT-25 when
327	compared with CT-20, and slight differences were found between CT-25 and CT-30. Therefore,
328	the textural properties increased with the amount of tannin added, and thus A_{BET} varied from 265
329	to 375 m ² ·g ⁻¹ while S_{NLDFT} varied from 850 to 910 m ² ·g ⁻¹ (Figure 6c). These large differences
330	between A_{BET} and S_{NLDFT} are due: first, to the fact that the BET method underestimates the surface
331	area when a significant portion of the pores are smaller than 0.7 nm in diameter; secondly, to the
332	fact S_{NLDFT} was calculated using N ₂ and H ₂ adsorption isotherms together, and the diameter of the
333	H_2 molecule is about 0.3 nm while that of N_2 is about 0.35 nm. Therefore, H_2 has access to
334	porous surface that N_2 cannot penetrate. The corresponding isotherms are presented in Figure 6b.
335	Figure 6d shows that the three materials had very similar PSDs, the main difference being a
336	slight increase in pores around 0.7 nm when increasing the percentage of tannin from 20 to 25%.
337	Figure 6e shows the corresponding cumulative pore volumes. The total pore volumes, V_T , were
338	relatively low, between 0.20 and 0.23 cm ³ g ⁻¹ (see Figure 6f), and mostly composed of
339	micropores. Only CT-30 had a few mesopores (pore diameter between 2 and 50 nm) but in
340	almost negligible amount, 0.01 cm ³ g ⁻¹ (thus not shown in Table S4). Most of the micropores, 75%
341	for CT-20 and 65% for CT-30, were in the ultramicropore range, and their average diameter was
342	0.58 ± 0.2 nm. The latter means that only H ₂ molecules are able to form a bilayer inside pores
343	with a diameter of about 0.6 nm, and thus provides insight into the large difference between A_{BET}
344	and S_{NLDFT} . The fraction of micropores where the H ₂ bilayer cannot be formed, corresponding to
345	diameters smaller than 0.5 nm, is still remarkable and varies from 65 to 59% for CT-20 and CT-
346	30, respectively.

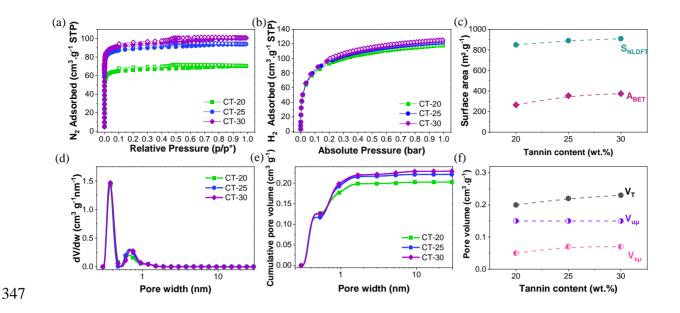


Figure 6. (a) Nitrogen; and (b) hydrogen adsorption (full symbols) – desorption (empty symbols)
isotherms; (c) BET area and NLDFT surface area of the printed carbons as a function of initial
tannin content; (d) pore size distributions; (e) cumulative pore volumes calculated by the 2DNLDFT model; and (f) pore volumes as a function of initial tannin content.

352 Raw and differential mercury intrusion curves plotted as a function of pressure and pore 353 diameter, respectively, are presented in Figure S5. Analysis of meso/macropores structure by 354 mercury porosimetry confirmed that most of the pore volume can be attributed to large pores with 355 a distribution centred on a diameter of about 10 µm, as observed in the SEM images. All pore 356 size distribution (PSDs) were very broad and rather multimodal, with pore diameters covering 357 three orders of magnitude. The addition of tannin to the initial resin was found to shift the median 358 pore diameter (calculated with Eq. S3) to lower values as the density increased, from 14 µm for 359 CT-0 to 5 μ m for CT-30.

360 The average pore diameters were significantly higher than those reported for other 3D-printed361 carbon monoliths prepared by stereolithography in the presence of porogens, followed by

362 oxidation and pyrolysis steps, which exhibited typical pore diameters (d_{90}) between 30 and 70 nm (but not including pores larger than 10 µm in diameter)¹². However, they were similar to those of 363 pyrolytic carbon electrodes derived from SLA technology with pore diameters of about 10 μm 11 364 365 and to those of tannin-derived monoliths prepared by emulsion templating, with values between 1 366 and 10 μ m^{42,43}. This is an important result because, unlike the latter materials, no particular 367 synthesis strategy was used here to obtain pores of these dimensions in such large numbers. This 368 result also indicates that this porosity is both wide and connected, and that it is therefore easy to 369 develop, by classical activation processes, the micro- and mesoporosity observed above by gas 370 adsorption. It should therefore be possible to obtain monolithic, architecturally structured 371 activated carbons, useful for example as highly permeable, rigid and easy-to-handle catalyst 372 supports, unlike granular or powdered carbons.

373 **Physical properties of printed carbon structures**

374 The electrical conductivity of a carbon structure may vary greatly as a function of texture, 375 which in turn is a function of precursor and heat treatment temperature, and also somewhat as a 376 function of the relative density of the material. Since the skeletal density and porosity of the 377 different 3D-printed carbons are similar (see again Table S4), their electrical conductivity was 378 assumed to remain of the same order of magnitude. Indeed, Table S5 shows values between 6.05 and 7.1 S cm⁻¹ for the different formulations. Similarly, the solid electrical conductivity, *i.e.*, the 379 380 electrical conductivity excluding the sample porosity, calculated using Eq. (4), was found to be 381 between 43 and 59 S \cdot cm⁻¹. This narrow range of values is reasonable considering the 382 approximations made to obtain Eq. (4) on the one hand, and is moreover typical of glassy carbon on the other hand, for which values of 11 S \cdot cm^{-1 44} to 70 S \cdot cm^{-1 45} have been reported. 383

These relatively modest conductivity values are of great interest, as 3D-printed carbonaceous architectures with low to medium conductivity struts have recently been shown to produce broadband photonic crystals in the GHz range, while higher conductivities combined with the same architectures result in higher absorption coefficients, but in an incomparably narrower frequency range ⁴⁶. To prove this, printed CT-25 parallelepipeds with final dimensions of 7.1 × 3.3×2 mm were analysed in the Ka band.

In the whole range of frequencies studied, the EM response is strongly dominated by the reflection, whose coefficient varies typically between 0.7 and 0.8. As a result, the absorption coefficient varies between roughly 0.2 and 0.3, and the transmission is almost zero (Figure 7). Analysis of the data collected yielded the shielding effectiveness, $SE = -20 \log(S_{21})$, and it can be concluded that the studied material provide a shielding level in the range of 20 to 22 dB. Although it is not very high, it is indeed a broadband absorption, which can be interesting to build photonic crystals and thus brings new perspectives to these porous biosourced carbons.

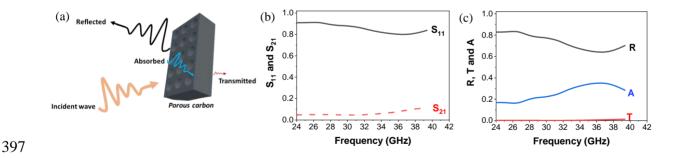


Figure 7. Electromagnetic properties of the CT-25 structure: (a) schematic of the EM response; (b)
frequency dependence of the *S* parameters ; (c) corresponding behaviour of the reflection,

400 transmission and absorption coefficients.

401 The strain-stress compression curves of 3D-printed carbon cubic samples are given in Figure
402 S6a-S6e. The compressive strength was defined as the highest stress at the first peak, the slope of

403 which was used to determine the Young's modulus. Following this first peak, the stress decreased 404 sharply, indicating a brittle failure of the material. The repetition of this phenomenon throughout 405 the compression of the samples led to serrated stress-strain characteristics explained by the 406 propagation of cracks in the carbon matrix. This behaviour then gave way to a progressive 407 densification of the material, as indicated by the increase in stress at strains above 35 %. 408 As expected, due to its denser structure obtained after pyrolysis (see again Table S4), the 409 sample printed with tannin-free acrylate resin showed relatively high compressive strength at 410 break (2.9 MPa) and high Young's modulus (156 MPa). These values decreased significantly to 411 1.4 MPa and 83 MPa, respectively, with the addition of 10 wt.% of tannin in the resin. This can 412 be explained by the decrease in density of the structures due to the low volume shrinkage 413 observed in the presence of tannin (Table S2). Then, due to this density gain, the mechanical 414 properties of the 3D-printed carbons increased with the increase of tannin content from 10 wt.% 415 to 30 wt.% to reach values close to those of the sample without tannin (CT-0), see Figure S6f and 416 Table S6. Because of the reduction in porosity with tannin loading from 10 to 30 wt.%, the 417 materials became more brittle, as indicated by the reduction in strain at break with values 418 decreasing from 7.0 to 2.8% (Table S6). In this printed system, tannin, being a bio-based carbon 419 precursor, helped to maintain the structure after pyrolysis and increase the material yield, but 420 only really appeared as a reinforcement at high loading. Under these conditions, at 30 wt.% of 421 tannin, mechanical properties are about equivalent to those of the carbonised pristine acrylate 422 resin, but with a significantly better yield and with a non-negligible share of renewable and 423 ecological material. Generally speaking, these mechanical properties are higher than those obtained for carbon manufactured by stereolithography with a similar density 1^{12} . 424

425

426 CONCLUSION

427 The present paper demonstrated for the first time an effective strategy to develop 3D-printed 428 carbon materials from bio-based resins using stereolithography (SLA). A new resin formulation 429 based on an acrylate-tannin suspension was developed to meet the viscosity requirements for 430 SLA printing. Indeed, despite its simple preparation and interesting mechanical properties, the 431 shrinkage and warpage of the carbon structure derived from the pure acrylate resin remained 432 significant and the carbon yield upon pyrolysis was low. Thus, the addition of tannin to the 433 formulation allowed porous structures to be printed with an encouraging carbon yield close to 434 20 %. At a loading above 20 wt.%, tannin significantly reduced the warpage and volume 435 shrinkage of the resulting carbon structures by about 35 %.

436 Increasing amounts of tannin were found to increase the bulk density of the structures while 437 the skeletal density remained the same, leading to a high total porosity of about 80%. However, 438 the non-graphitizable and very disordered carbon structures exhibited low surface area and wide 439 pore size distribution, demonstrating the presence of a limited amount of very narrow pores 440 together with very large pores with a median diameter ranging from 5 to 10 μ m. These 441 characteristics demonstrate the possibility of preparing architectured activated carbons, for 442 instance for catalytic applications, as the pre-existing porosity can be further developed by 443 conventional activation methods.

Theses carbon architectures also showed correct electrical conductivities in the range of 6-7 S cm^{-1} , varying with the textural properties, with a solid electrical conductivity between 43 and 59 S·cm⁻¹, quite typical of glassy carbon. These properties open the way to the preparation of more sustainable 3D-printed devices for electrochemical applications, in which glassy carbon is widely used, but also for electromagnetic applications, for which moderate conductivities might be preferable if broadband absorption is required. Finally, the new 3D-printed carbon structures
presented here exhibited high mechanical properties with elastic-brittle behaviour. Nevertheless,
in this system, tannin as a carbon precursor really appeared as a reinforcement only at high
loading (> 25 wt.%).

The possibility of obtaining complex architectures with a non-negligible share of renewable and bio-based material, with good carbon yield and little shrinkage to tailor their properties, makes our approach promising for a broad range of applications. Thus, more detailed studies of such 3D carbon architectures will be carried out in the near future.

457 ASSOCIATED CONTENT

458 Supporting information: Additional experimental details and results including data and figures
459 related to elemental analysis, porous texture, mechanical properties, particle size distribution of
460 mimosa tannin, and SEM and TEM images of two different 3D-printed carbons.

461 AUTHOR INFORMATION

462 **Corresponding Author**

^{*} Corresponding author. Tel: + 33 372 74 96 14. E-mail address: <u>alain.celzard@univ-lorraine.fr</u>
(A. Celzard)

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626 GRAPHICAL ABSTRACT

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- 629 1 For Table of contents Use only
- 630
- 631 SYNOPSIS
- 632 Highly porous, conductive and mechanically resistant 3D-printed bio-based carbon were
- 633 produced from tannin precursors through stereolithography.
- 634